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FACULTAD DE CIENCIAS FÍSICAS

DEPARTAMENTO ÓPTICA



TESIS DOCTORAL

**Optical and structural properties of the crystalline lens:
Accommodation and aging**

**Propiedades ópticas y estructurales del cristalino:
acomodación y envejecimiento**

MEMORIA PARA OPTAR AL GRADO DE DOCTORA

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Introduction

1.1 Motivation

This thesis addresses the influence of shape and gradient refractive index to the optical properties of various mammal lenses, in relation to age or accommodation. With age, the human lens undergoes various physical, biometrical and optical changes. Physical and biometrical changes are well documented in the literature. Optical changes of the lens have been most often assessed from geometrical and biometrical measurements of cornea and lens, and not from the optical properties from the lens itself. The optics of the crystalline lens, especially the gradient refractive index (GRIN) distribution are still under discussion. Although lens shape and GRIN have been measured with various methods, the information on it is usually obtained separately. In this work, we measured the lens shape and the focal length of crystalline lenses and combined this information with the reconstructed GRIN of the same lens to draw conclusion about the lens' optics and the influence of the GRIN as a parameter of optical lens property. The understanding of the structural changes of the crystalline lens with aging and accommodation is important to gain insights into the mechanisms of aging of the eye, and in particular presbyopia development, and its potential treatment. For a deeper understanding of lens structure and mechanics of accommodation on a microscopic level, we used a custom built confocal microscope to image microscopic structures of the crystalline lens. The

development of instruments to analyze the accommodating lens on a microscopic level is important for the understanding of the lens mechanics. This is especially interesting when thinking of future developments of accommodative intraocular lenses.

1.2 The crystalline lens with age and accommodation

The crystalline lens, as part of the anterior segment of the eye, is located behind the cornea and the iris. It is composed of a clear membrane (the lens capsule), a subcapsular lens epithelium (only anteriorly), and lens fibers, which emerge from the epithelium cells. The lens, together with the cornea, transmits and refracts light to form an image of the

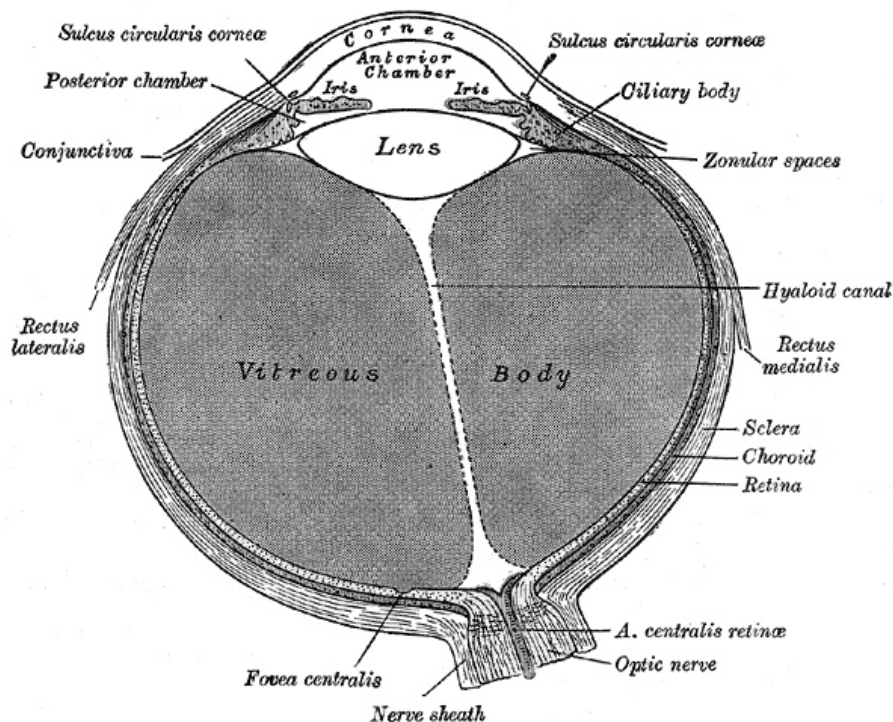


Figure 1.1: The crystalline lens is located behind the cornea and the iris. Reproduced from Gray's Anatomy[240].

world on the retina. The elastic young crystalline lens changes shape via constriction and/or relaxation of the ciliary muscle to which the lens zonules are attached, resulting in a change of the optical power of the eye which provides a focused image of objects both near and far. With aging, the human lens becomes thicker, relatively steeper, and

looses the ability to accommodate [10, 58, 72], a condition known as presbyopia. The refractive index of the lens is not homogeneous. It is distributed over the lens gradually, with a high refractive index in the lens center, and a lower refractive index towards the lens surface. This gradient refractive index (GRIN) changes with age[39, 48, 89, 104].

1.2.1 Change of ex vivo crystalline lens shape and optics with aging

Measurements on lens shape and optical properties have been done, both on in vitro and in vivo lenses. An important difference between in vivo and in vitro is, that in vitro lenses are usually cut off from any zonular tension and, therefore, remain in their most accommodated state. In the last decades, there have been numerous reports of in vitro measurements[32, 73, 96, 145, 183, 195, 218]. The first studies typically used frozen cadaver lenses, cut into thin sections, to study the lens shape[96, 183]. In the last years, laser ray tracing, linked to photographic or digitized recording were used. Pierscionek and Augusteyn[195] for example measured lens dimensions using laser ray tracing. The entrance beam was split into several component rays and the ray paths were recorded photographically. The lens shape was then determined from the photographic negatives. The results suggested that growth in sagittal width is predominantly in the anterior part of the lens. Glasser and Campbell[73] studied the biometric, optical and physical properties of in vitro human lenses. Focal length and spherical aberration were measured using a similar scanning laser apparatus. The radius of curvature and lens were obtained by digitizing the lens profiles.

Manns et al.[145] measured the shape of the anterior and posterior surface of human cadaver lenses in situ using a corneal topography system. Rosen et al.[218] and Borja et al.[32] determined dimensions, curvatures and asphericity of excised human lenses using shadow photography.

In general, with the in vitro lens maximally accommodated, the anterior and posterior lens surface radii of curvature tend to increase with age, at least up to past the presbyopia onset[145]. At the same time, the lens thickness increases with age from around 3.5 mm to 4.5 mm. However, between the age of 40 and 60 years, when the lens suffers from presbyopia, the radii of curvature decreases again.

Measurements of other geometrical aspects of the lens surfaces, such as their asphericity, are scarce. From above mentioned authors for example, only Manns et al. and Parker and Howcraft included asphericity measurements in their studies. It was suggested by Perez-Escudero et al.[203] that lens radius of curvature and asphericity should not be interpreted separately.

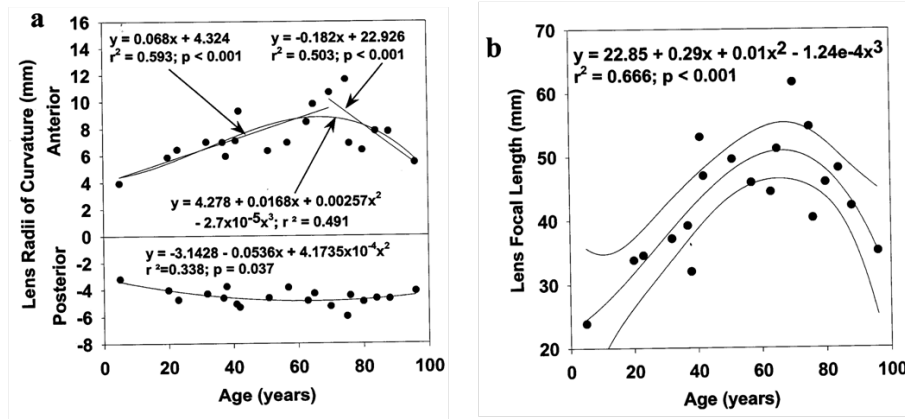


Figure 1.2: (A) Lens radius of curvature with age. While the posterior radius of curvature is almost constant, the anterior radius of curvature can be fitted by two linear regression lines below and above 65 years of age and with a third order polynomial. (B) Linear increase in lens focal length with age from five to roughly 65 years of age and thereafter a decrease in lens focal length with age for isolated lenses. Reproduced from Glasser et al.[73]

Several authors measured lens power with increasing age[18, 32, 33, 72]. In their study on presbyopia and optical changes with age, which included the use of a stretching device to simulate accommodation, Glasser and Campbell[72] found that the focal length (inverse of lens power) of the unstretched lenses increased linearly with age (measurements were done with a scanning laser apparatus). Borja et al. measured the lens power on humans[32] and non-human primates[33] with a modified commercial lensmeter or an optical system based on the Scheiner principle. In human lenses, the total refractive power showed a biphasic age dependency (decreasing between ages 6 and 58.1 from around 45 D to 20 D, and increasing between ages 58.1 and 82). For non-human primates, the authors found that the total isolated lens refractive power decreased with age. A recent study by Augusteyn et al.[18] found that the lens power decreased with age from 40 to 50 D at birth to 20 D after age 50.

Due to the age related geometrical changes in the lens, the optical properties of the lens change as well. Spherical aberration were measured on in vitro animal lenses, including fish lenses[120, 122, 121, 229], amphibious lenses[120, 229], bird lenses[74, 120, 229, 230], mammal lenses[2, 47, 227, 229], and non-human primate lenses[213]. In many of the mentioned literature, a split beam laser technique was used to study the amount of spherical aberration. The parallel beams were directed through the excised crystalline lens and the spherical aberration was measured by photographing the variation in back vertex distance for the laser beams. In most species, negative spherical aberration was

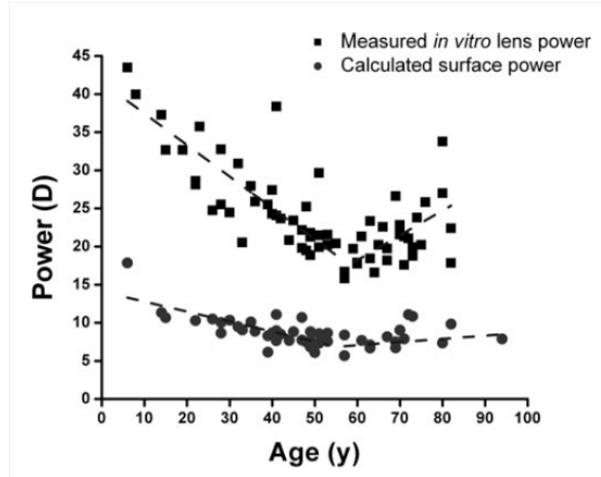


Figure 1.3: In vitro lens refractive power and calculated surface refractive power. The total refractive power showed a biphasic age dependency. Reproduced from Borja et al.[32].

found in the crystalline lens, for example for pigs[259], rats[40, 229], fish[122], duck[229], dog[229], frog[229], squirrel[227], chicken[74]. Aberration changes with age in in vitro human lenses were studied for example by Sivak et al.[226] and Glasser et al.[72, 73]. In both cases, the above mentioned scanning laser technique was used. Sivak found positive and negative spherical aberration in five human lenses between 16 and 78 years. Glasser found that spherical aberration was negative for young lenses and positive for older lenses in a study of 27 humans between 10 and 87 years[72], and another study of 19 pairs of human lenses, between 5 and 96 years of age[73].

1.2.2 Change of ex vivo crystalline lens shape and optics with simulated accommodation

In vitro, changes in lens shape and spherical aberration with accommodation can only be measured using a stretching device which mimics the radial accommodative forces on the lens (see section 2.1.4). Stretching experiments were recently done in human lenses[18, 72, 147], and monkey lenses[147, 142, 213]. Obviously, stretching does not reproduce natural accommodation exactly, mainly because of the absence of vitreous and intraocular pressure and the differences in the direction of the forces involved in accommodation[213], but the good correspondence between the geometrical and optical changes in the lens as a function of accommodation measured in vivo and in vitro has demonstrated in the recent years that stretching is an adequate model and that the results are in general consistent with the Helmholtz theory of accommodation[18,

of the accommodative amplitude from 12 to 14 D in the young lens to around 0 D, already by the age of 40-50, leaving the mature eye at a lens power of around 20 D. This age-dependency in lens power-decrease with stretching was also found for monkey lenses[147].

There are only a few studies which measure spherical aberration with simulated accommodation [72, 213]. Usually a laser scanning technique is applied to define the amount of spherical aberration. Roorda et al.[213] used a least squares procedure to fit a Zernike polynomial function, which was obtained from local slopes of the wavefront which emerged from the measured trajectories for an array of narrow refracted laser beams scanned through the lens[213]. Glasser et al. also used a scanning laser technique to calculate the spherical aberration on human lenses. He scanned the lens, and recorded beam height from the optical axis at the entrance beam, and the distance from the principal point of the exit beam to its intersection with the optical axis. A fourth order polynomial was fitted to these coordinates. The spherical aberration was then calculated from the difference in dioptric distance of peripheral refracted rays incident at the lens periphery and paraxial refracted rays incident near the optical axis of the lens. For both, monkey lenses and human lenses it was found that spherical aberration became less negative with stretching (that is, more negative with accommodation).

1.2.3 Change of in vivo Crystalline lens shape and optics with aging

The first attempts to measure the shape of the anterior crystalline lens surface in vivo used slit lamp photography on iridectomized eyes[35, 36, 139, 140]. Lowe et al.[139, 140] measured only the anterior curvature of eyes, while Brown et al.[36] measured via Scheimpflug photography, on average in younger eyes than Lowe, and included the posterior radius of curvature.

The first Scheimpflug imaging system used for investigation of the crystalline lens in vivo was developed by Brown himself[35]. The Scheimpflug configuration allows to image the anterior segment of the eye with large depth of focus, but a geometrical and optical distortions are introduced, since the magnification is not constant, and each surface of the anterior segment is seen through the previous surface. In order to obtain reliable information from those images, a distortion correction must be introduced. If corrected, for example with the approximation of a constant refractive index within the lens, lens shape and its change with age and accommodation can be obtained[35, 36, 59, 116, 115, 214, 216]. In Brown's study on one hundred healthy, emmetropic subjects between the age of 3 and 82 years, he demonstrated that the lens becomes more convex with

age and that the anterior radius of curvature decreases around 6.5 mm between the age of 20 and 80. The posterior radius of curvature decreased less (around 1.5 mm) in comparison. Koretz et al.[115] used Scheimpflug imaging to measure lens radii of curvature in unaccommodated lenses in subjects between 18 and 70 years. Anterior and posterior lens radius of curvature decreased with increasing age, in qualitative but not quantitative agreement with the earlier observations of Brown. Dubbelman et al.[57] measured the unaccommodated in vivo lens with Scheimpflug images on 102 subjects between the age of 16 and 65 years. Correction for the refraction distortion produced by the cornea were applied. The anterior and posterior radius of curvature decreased with age, however, posteriorly, the radius of curvature decreased slightly. The conic constant of both surfaces were not age-dependent. In another study, Dubbelman et al.[59] used Scheimpflug imaging to show that, with age, the increase in thickness of the cortex is approximately 7 times greater than that of the nucleus. The increase in thickness of the anterior cortex was found to be 1.5 times greater than that of the posterior cortex. He also found, that it was only a specific zone in the lens cortex that is responsible for the increase in lens thickness with age. As mentioned before for in vitro lenses, also for in

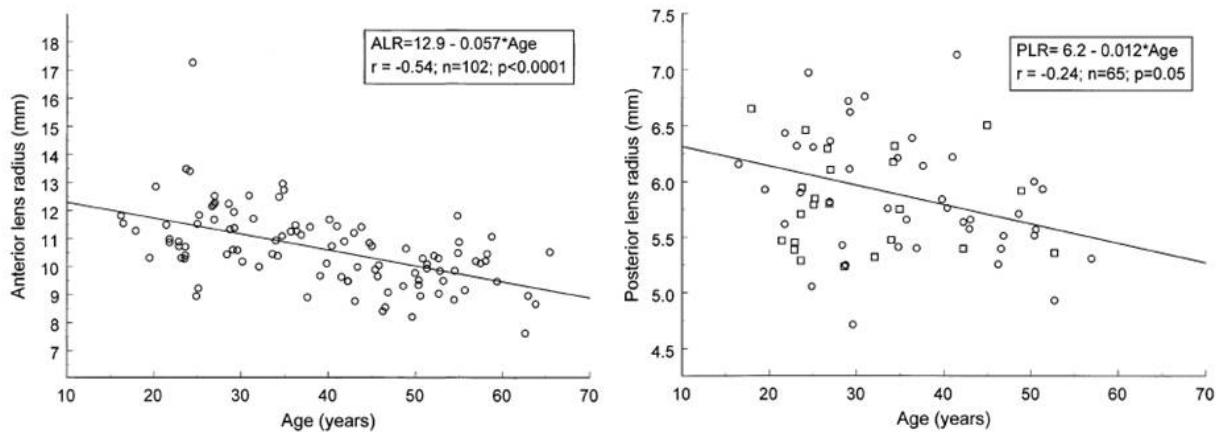


Figure 1.5: Anterior and posterior radius of curvature with age, measured by Dubbelman et al. on 202 subjects in vivo (unaccommodated). Reproduced from Dubbelman et al.[57].

vivo lenses, the optical properties changes with age. In particular, the overall spherical aberration of the eye shifts toward more positive values[6, 72, 157]. This change in lens spherical aberration leads to age-related loss of the corneal/lens spherical aberration compensation and a decrease of the optical performance [8, 22, 81, 232].

Assessment of the optical aberrations of the lens in vivo has generally involved measurements of the total aberrations of the eye using wavefront sensing and subtraction of the corneal aberrations estimated from corneal topography. Alternative approaches

have involved measurements in immersion, which canceled the aberrations of the anterior cornea. In both cases, these indirect measurements include, apart from the crystalline lens aberrations, contributions from the posterior cornea as well as the effect of the beam convergence onto the lens. By using the method of subtracting the corneal aberrations, Jenkins et al.[103] suggested that the lens' spherical aberration was approximately zero. El Hage et al.[82] measured the contribution of the crystalline lens to the spherical aberration of the eye, and concluded that the lens, showing negative spherical aberration, plays an important compensatory role, reducing the positive spherical aberration of the cornea. The results of Millodot et al.[159], who eliminated the corneal power by using swimming goggles filled with saline water, showed that the aberration of the lens does not systematically neutralize that of the cornea. Tomlinson et al.[252] reported in vivo data for 20 subjects between 20 and 56 years, and found that the spherical aberration of the crystalline lens was generally negative and in the range of 0 to -0.81 D. Smith et al.[232] supported Tomlinson's findings with 26 tested subjects (between 20 and 72 years). He found an age effect, with older eyes having a larger aberration for the whole eye, and concluded that this change is most likely due to the age related changes in the crystalline lens with the spherical aberration of the lens becoming less negative with age. In a study by Artal et al.[9], the aberrations of the internal surfaces were obtained

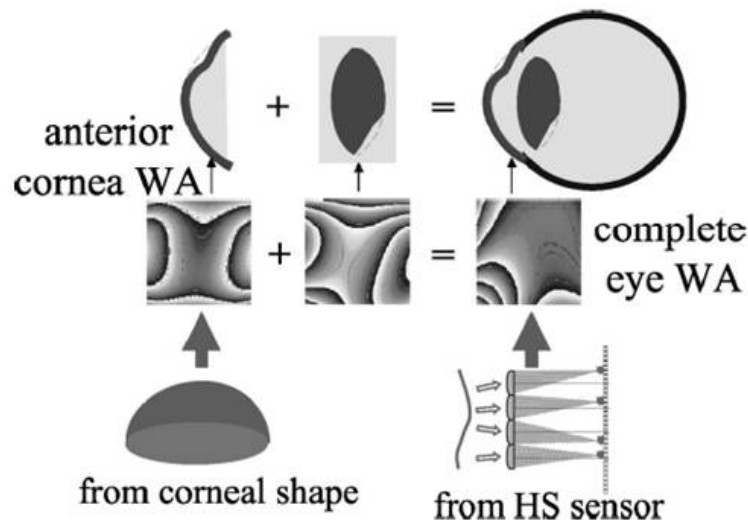


Figure 1.6: Schematic representation of the methods by Artal et al. The aberrations of the internal ocular media were obtained by direct subtraction of the anterior corneal wave aberration from the ocular wave aberration. Reproduced from Artal et al.[9].

by direct subtraction of the ocular and corneal wave-front data. The data was obtained using a Hartmann-Shack sensor and a corneal topography system. He measured a difference between total ocular aberration and corneal aberration in younger eyes, with total

ocular aberration lower than corneal aberrations, and found that the opposite occurred in older eyes. These results were supported by Barbero et al.[22].

1.2.4 Change of in vivo Crystalline lens shape and optics with accommodation

Several works have studied the lens shape, power, tilt and decentration in vivo using Purkinje imaging[12, 164, 187, 214, 215, 217, 238, 245, 257, 270]. Purkinje imaging was described by Purkinje in 1832, and since then it has been widely used to obtain the above mentioned characteristics of the crystalline lens. Purkinje images show the reflections of the structure of the eye. They are categorized from first Purkinje image (PI) to fourth Purkinje image (PIV). PI is the reflection from the anterior surface of the cornea, PII from the posterior surface of the cornea, PIII is the reflection from the anterior surface of the lens, and PIV is the reflection from the posterior surface of the lens. Information on the shape of the crystalline lens are obtained using PIII and PIV reflections. Early systems were based on photography[257, 270], and some versions were used to study the correlation between the refractive error of the crystalline lens and its geometrical properties[238]. Purkinje measurements as a function of accommodation

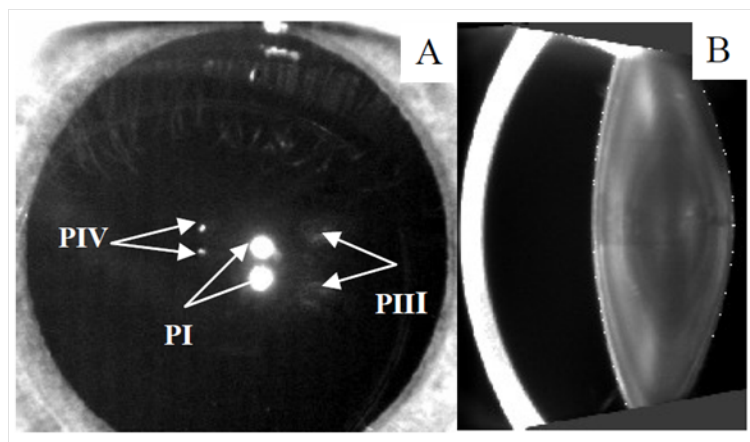


Figure 1.7: A typical example of a Purkinje image showing double PI, PIII and PIV (A) and an example of a corrected Scheimpflug image (B), for the same unaccommodated eye. Reproduced from Rosales et al.[214].

were done for example by Garner et al.[70] and Rosales et al.[217]. Rosales calculated the values of the relaxed lens radius of curvature of iridectomized rhesus monkey eyes to (anterior) 11.11 ± 1.58 mm and (posterior) -6.64 ± 0.62 mm. With accommodation these values decreased at a rate of (anterior) 0.48 ± 0.14 mm/D and (posterior) $0.17 \pm$

0.03 mm/D for anterior and posterior lens surfaces, respectively. Tilt and decentration did not change significantly with accommodation except for tilt around the horizontal axis. In his studies on age and accommodation, Dubbelman et al. took Scheimpflug images of both, unaccommodated and accommodated eyes of 102 subjects (16-65 y)[59]. He concluded that, with accommodation, the nucleus becomes thicker, but the thickness of the cortex remains constant. In a later study, including 65 subjects between 16 and 51 years, he showed a decrease in anterior and posterior radius of curvature during accommodation[58]. He found that the increase in lens thickness with accommodation is higher than the decrease in the anterior chamber depth, and, therefore, suggested that the posterior lens surface moves backwards with accommodation. A comparison study between Scheimpflug imaging and Purkinje imaging was done by Rosales et al.[214]. Measurements of the lens radius of curvature were performed using both instruments. Comparisons were also made as a function of accommodation (0 to 7 D) in a subset of 11 eyes. It was concluded that both techniques provide comparable lens radii and similar changes with accommodation. Purkinje imaging tended to overestimate the posterior lens radius, while pupil size limited the acquisition of posterior lens data with the Scheimpflug camera.

The change in lens shape during accommodation, results in a change in aberrations of the eye. There are several reports on the change in optical quality of the eye with accommodation in vivo in humans[10, 42, 68, 87, 141] and non-human primates[261, 260]. Atchison et al.[10] measured monochromatic ocular aberrations of 15 human eyes as a function of accommodation using an advanced Howland Aberroscope Technique. The instrument included a fundus camera for photography, a stimulus system for accommodation, and a method for correction for grid projection at the entrance pupil of the eye. With increase in accommodation, a trend towards negative spherical aberration was found for about half of the subjects tested. These results were also confirmed by He et al.[87], who measured wave-front aberration for eight undilated subjects using a spatially resolved refractometer and accommodation stimuli from 0 to -6 D. Besides, He et al. found a systematic shift towards more negative or less positive aberration in all eyes. On average, ocular spherical aberration crossed zero for about a 2-D stimulus. Dubbelman[58] indirectly supported these results with his earlier mentioned study on lens shape with accommodation. He found that during accommodation, the anterior lens surface becomes more hyperbolic, which is consistent with the mentioned shift in spherical aberration. Cheng et al.[42] used a Hartmann-Shack wavefront sensor to measure wave aberrations in crystalline lenses, using accommodative demands of 0, 3, and 6 D. It was found that spherical aberration showed the greatest change with accommodation, changing to more negative values with accommodation, while coma and astigmatism showed minor changes. Gamba et al.[68] also used a Hartman-Shack sensor to measure

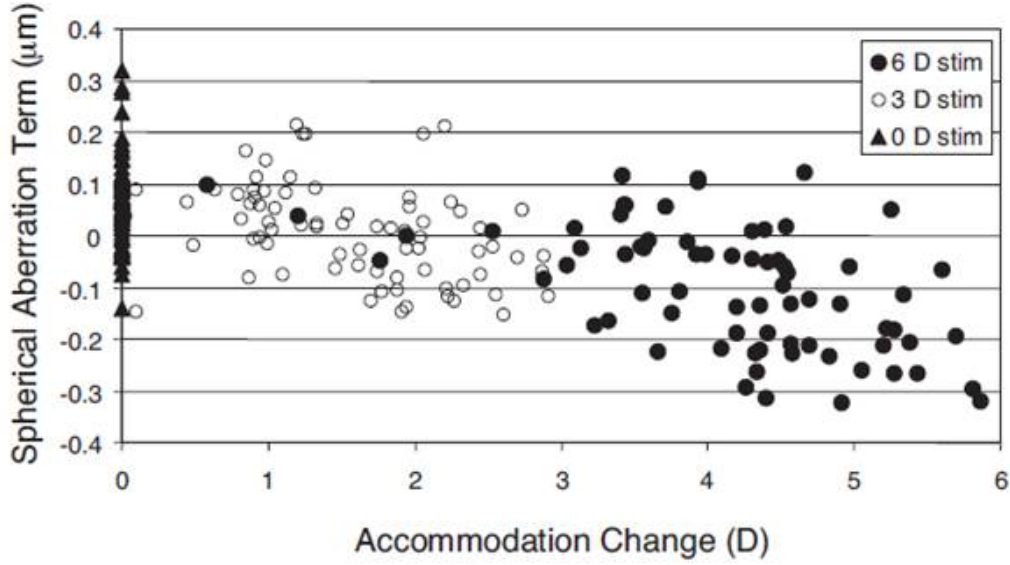


Figure 1.8: The coefficient for spherical aberration as a function of the change in accommodative response 76 young subjects under different accommodation stimulus. Reproduced from Cheng et al.[42].

the accommodative response (to a stimulus moving from 0 to 6 D in increments of 1 D) and aberration changes on 5 subjects for natural, corrected and induced optical aberrations. In this study it was demonstrated that aberrations influence accommodative lag and fluctuations of accommodation. Spherical aberration was the term that varied the most during accommodation (shift towards more negative values). Other aberration terms, like coma and astigmatism did not change systematically with accommodation. Their slope could be compared to above mentioned literature[42, 87].

In a study with iridectomized rhesus monkey eyes by Vilupuru et al.[261], accommodation was centrally stimulated to various accommodative demands (0 D to ~ 11 D). The spherical aberration became more negative with accommodation, while vertical coma increased. It was also found that the change in lens power with accommodation was much more pronounced in the lens center than in the lens periphery.

1.2.5 Lens Astigmatism

Early attempts to measure the astigmatism of the internal ocular surfaces involved the use of videokeratography, A-scan ultrasonography, and autorefractometry along with multi-meridional phakometric measurements of Purkinje images[214]. As recognized by the authors[26], the ophthalmo-phakometric method was prone to considerable accumu-

lated experimental errors. However, the work reports some interesting findings, such as predominance of inverse astigmatism (defined as the one where the steeper meridian is in the horizontal axis) in the posterior cornea and in the posterior lens surface, while the astigmatism of the anterior surface of the lens was direct (defined as the one where the steeper meridian is in the vertical axis). Compensatory effects of corneal astigmatism by the crystalline lens[7, 8, 110, 109] by evaluation of corneal and total astigmatism have been reported, as well as compensatory effects of the anterior corneal astigmatism by the posterior corneal surface[55, 60, 61, 180, 205, 209]. In addition, in his PhD thesis[63], Elawad measured ocular component contributions to residual astigmatism in human eyes, and found, whilst the astigmatic contributions of the posterior corneal and lens surfaces were found to be predominantly inverse, that direct astigmatism came from the anterior lens surface.

1.2.6 Lens shape and OCT imaging

Optical Coherence Tomography (OCT) allows a non-invasive insight into the human body, therefore, it has become an important tool in medicine. 3D OCT imaging creates a set of sliced images which can be reconstructed into a 3D image. OCT is based on light interference. The coherent light is separated into a reference beam and the beam which focuses on the sample. The light is reflected by the sample and returns to the reference beam where they interfere[54]. The use of light allows a high resolution and does not require any contact with the tissue. Since ocular tissue is transparent, it is very suitable for OCT imaging. In comparison to aforementioned imaging modalities, OCT has some strong advantages when imaging the eye, because of the high resolution images, the imaging range, and the direct view it gives into the eye. Usually the tissue is scanned both, in depth and laterally, by means of a time-coherent light source. If the depth scan is attained by moving the reference mirror, the OCT is time-domain. If the interference signals are obtained by splitting up from a light source with a wider bandwidth, the OCT is spectral-domain. OCT imaging was used to image the crystalline lens ex vivo and in vitro. Uhlhorn et al.[254] used a custom built time-domain OCT (2-D) to develop a method to measure the axial thickness and the average refractive index of the 40 ex vivo lens from a single recorded image. He used the optical thickness of the lens and the optical displacement of the posterior window of the sample chamber to calculate thickness and group refractive index. Borja et al.[34] quantified the posterior surface distortion in 2-D OCT images of in vitro crystalline lenses. The lens was flipped over in between the measurement to obtain the undistorted and distorted (due to refraction at the anterior lens and the index gradient) shape of the posterior lens surface. Radius of curvature and asphericity could then be compared for the distorted

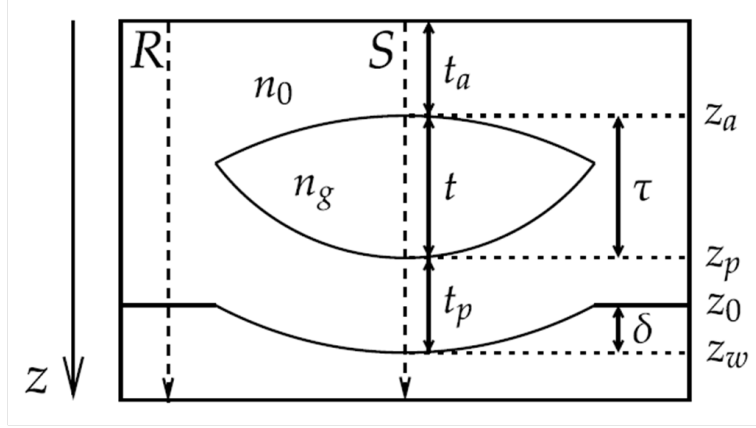


Figure 1.9: The crystalline lens with group refractive index n_g , and geometric thickness t , is immersed in a solution with group index n_0 . The optical thickness of the lens s , and the optical displacement d , of the apparent location of the window surface is measured from the OCT image data. The lens thickness t , and group refractive index n_g , are then calculated. Reproduced from Uhlhorn et al.[254]

and the undistorted image. It was concluded that the error in the radius of curvature is within the measurement reproducibility, which can be corrected assuming a uniform refractive index. Asphericity values on the other hand require a correction algorithm that takes into account the gradient refractive index. Siedlecki et al.[225] then estimated the contribution of the gradient index structure to the optical distortion of the posterior surface, using Borja's data. He found that the gradient refractive index played a significant role in the distortion of the posterior surface of the lens, viewed by OCT. To reconstruct the gradient refractive index, de Castro et al.[47] used a custom-developed high resolution spectral domain OCT system to obtain 3-D images of an ex-vivo porcine crystalline lens. He also used the method of flipping the lens to obtain distorted and undistorted images of the different lens surfaces.

In collaboration with the Bascom Palmer group, de Castro et al.[48] used the developed method of OCT imaging and reconstructing the gradient refractive index on human lenses of various ages to conclude about their age-dependency (see also next section). Maceo et al. used OCT to measure the lens shape in a stretching system. She combined the lens shape measurements with measurements of the calculated back vertex power of the stretched lenses, to determine the contribution of the gradient refractive index to the change in lens power.

Ortiz et al.[175] imaged the eyes of three subjects to characterize the human lens in vivo and in 3-D. All images were corrected for fan distortion and optical distortion (i.e. scanning and refraction distortion, respectively), and the lens surfaces were fitted by

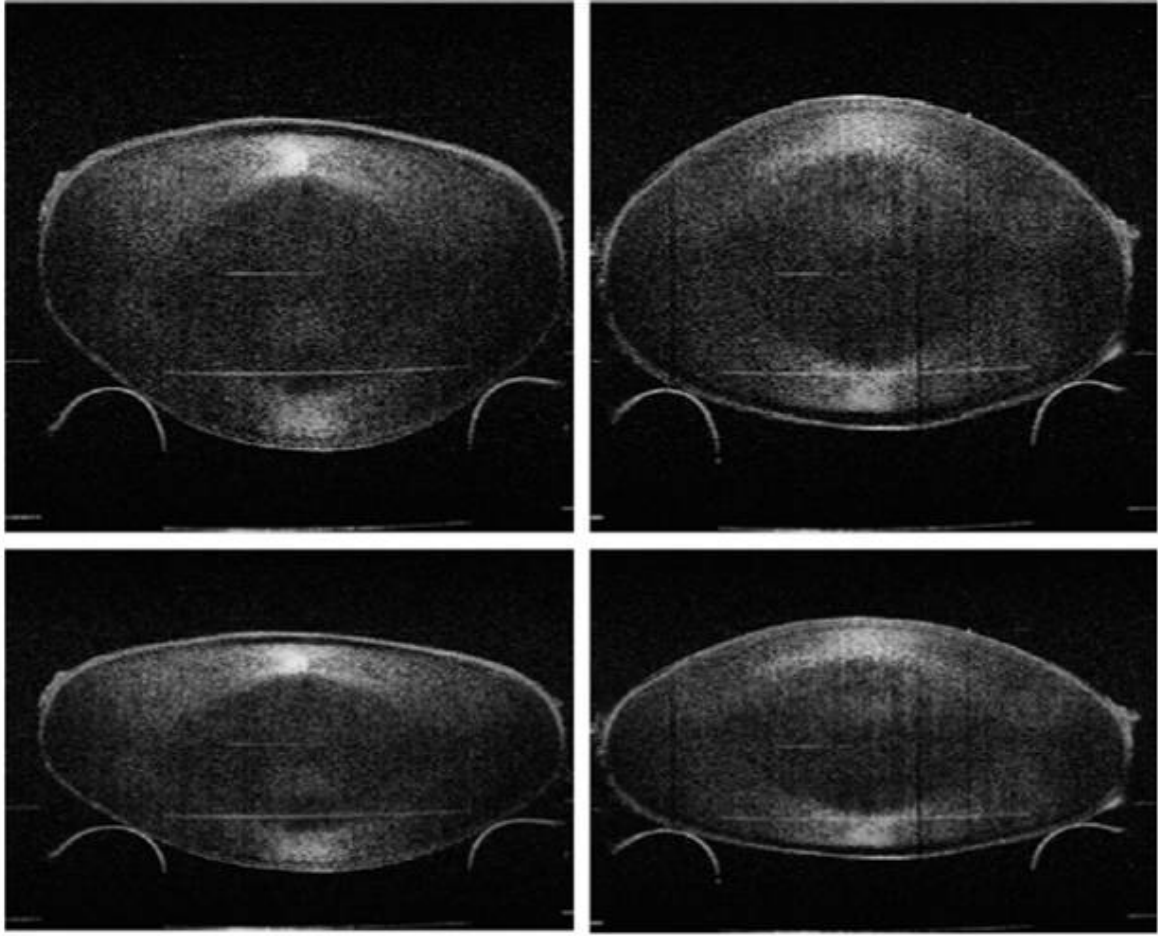


Figure 1.10: Raw (top) and rescaled (bottom) OCT images of a 49 year old human crystalline. Left: Anterior-up OCT image; Right: Posterior-up OCT image. Reproduced from Borja et al.[34].

by biconicoids and Zernike polynomials. The calculated anterior lens radii of curvature ranged from 10.27 to 14.14 mm, and the posterior lens radii of curvature from 6.12 to 7.54 mm. Ortiz concluded that the correction of optical distortion is critical and that the retrieved lens radii of curvature from OCT agree with phakometric data using Scheimpflug and Purkinje image along one meridian.

As an alternative to Scheimpflug and Purkinje imaging, OCT was also used to measure tilt and decentration of in vivo crystalline lenses and IOLs after cataract surgery. Ortiz et al.[174] analyzed in vivo anterior segment biometry before and after cataract surgery, using a 3-D spectral domain OCT. He retrieved biometric data of the complete anterior segment with a high degree of accuracy. Sun et al.[244] estimated tilt and decentration

of IOLs after cataract surgery, using the same OCT as Ortiz et al. Sun et al. used a Purkinje-based methodology from anterior segment en face OCT imaging. En face OCT images as Purkinje-like images that correspond to the specular reflections from the corneal and IOL surfaces. It was shown that average IOL tilt and decentration from Purkinje were 3.30 ± 4.68 deg and 0.16 ± 0.16 mm, respectively, and differed on average by 0.5 deg and 0.09 mm, respectively, from direct measurements on distortion-corrected OCT images.

1.2.7 Lens topography

Given its high accessibility, the corneal topography has been characterized much more excessively than the lens topography. Corneal irregularities in the elevation map and astigmatism are common. Interestingly, the posterior corneal surface appears to compensate part of the errors of the anterior cornea. In particular, the posterior cornea has been reported to compensate around 3.5% of the coma of the anterior surface[56]. There are also numerous reports of the compensation of the corneal astigmatism by the astigmatism of the crystalline lens using corneal and refractive parameters[110, 119, 136, 182]. However, there is little information on possible compensations within the crystalline lens itself, i.e. between anterior and posterior surface or between its surfaces and its GRIN. Manns et al.[145] presented an experimental technique to measure the topography of the lenses of Eye-Bank eyes in situ, using the PAR Corneal Topography System. He presented typical lens topography maps and height data, but the focus of the work was to introduce the measurement method as an effective way to obtain the shape of the anterior and posterior surface. Above mentioned study of Ortiz et al.[175] on in vivo human lens topography found that the surface lens astigmatism was significant, and that the three subjects in the study showed perpendicular orientation in the astigmatism of the anterior versus the posterior surface. The horizontal astigmatism for the anterior surface ranged from -11 to 1 μm , and the vertical astigmatism on the posterior lens surface ranged from 6 to 10 μm . The lens elevation maps showed minimal amounts of trefoil, coma or other high order aberrations. It was concluded that the fact that the lens surface topographies are relative smooth suggests a contribution of the internal lens structure to the reported lens optical irregularities.

1.3 Gradient Refractive Index

It is well known that the optical properties of the lens are not only dependent on its shape, but also on its refractive index [1, 14, 70, 236, 246]. Because of the gradient

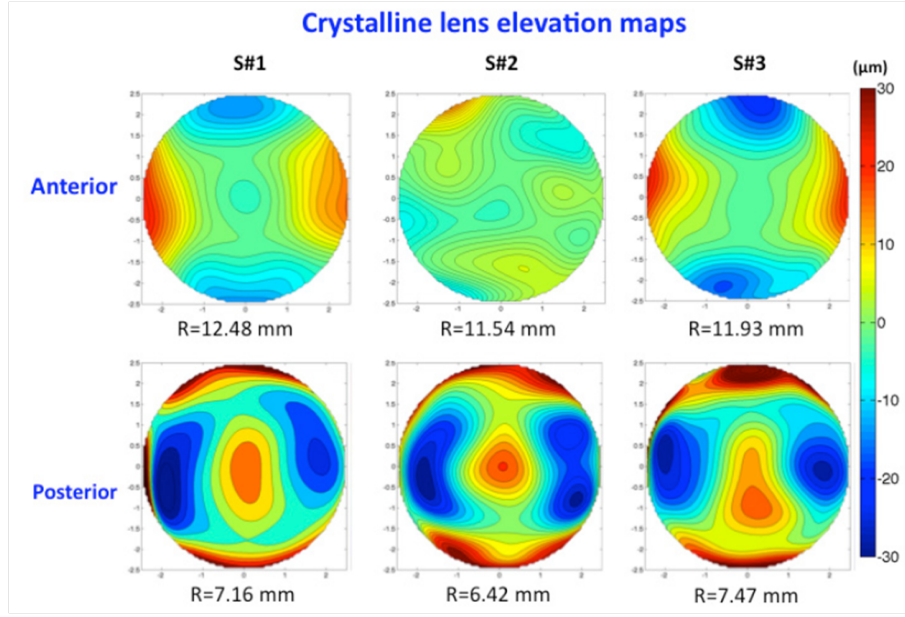


Figure 1.11: Quantitative anterior (top) and posterior (bottom) crystalline lens elevation maps in 3 eyes in vivo, after full distortion correction. Maps are Zernike fits to the elevation maps, relative to the best fitting sphere. Data are for 5-mm pupils. Reproduced from Ortiz et al.[175]

refractive index of the crystalline lens (GRIN), refraction does not take place only at the lens surface, but also within the lens medium. The gradient distribution of the refractive index is due to a non-uniform distribution of protein concentrations within the lens and varies across species[16, 198, 236].

1.3.1 Measurement methods

The gradient refractive index has been measured in several species, for example fish[100, 258], rat[39], cat[99, 165], rabbit[166, 190], porcine[256, 47, 196], and human[167, 194, 104, 107, 48]), using both, direct destructive[190][194] and indirect non-destructive methods (Ray Tracing[39, 256, 71, 196]; MRI[104, 107, 71], and OCT[258, 47, 48]). In the following, we give a brief overview over the techniques.

Destructive Methods

Destructive techniques to measure the lens gradient refractive index were done as early as 1880 by Matthiessen et al.[156]. In this study, the GRIN was measured using Abbe refractometry, a method which was also employed by Huggert et al.[98]. Nakao et

al. used Schlieren interferometry to measure the GRIN of rabbit[166] and human[167] lenses. Therefore, the lenses were frozen and thin layers were cut off from the lens to study Fraunhofer diffraction patterns. This method would show the different sections of the GRIN. Other destructive methods to measure the GRIN include Pulfrich refractometry, where lenses were cut and the refractive index was measured in the sample surface[181, 99] and Reflectometry[190]. In Reflectometry, a fiber optic sensor is used to estimate the lens GRIN. The method is based on the fact that the amount of light reflected in the interface between different media depends on the difference of the refractive index. Pierscionek used this method to analyze the GRIN of different species, like sheep, rabbit[190], bovine[191], and human lenses[194, 192].

Non Destructive Methods

Ray Tracing A report on the first non destructive method to measure the GRIN was written by Campbell et al. in 1984[39], based on the work by Chu[43] and Barrel and Pask[25]. Campbell measured the refractive index relative to the value at the lens surface. After passing through the lens, angular deviations of laser beams are measured and an inversion procedure is applied to the data leading to the refractive index. The lens itself has to be immersed into a medium that matches the surface refractive index of the lens. Ray tracing was also used by Garner et al. [71] with similar results. Pierscionek et al. used ray tracing to measure the GRIN in different species ([196, 197, 188]. For porcine lenses, she used two different wavelength and found the wavelength dependency to be small.

Magnet Resonance Imaging Moffat and Pope[160] reported a linear relationship between the refractive index and nuclear spin transverse relaxation rates. This knowledge permitted measurements of the GRIN using MRI[162]. Subsequently, Garner et al. used MRI technique to measure the GRIN and compared it to laser ray tracing results[71]. The comparison showed good agreement in the cortical zone, but MRI measurements were not possible in the lens core, due to the lack of free water. Jones et al.[104] measured the GRIN of 20 human lenses between the age of 7 and 82, using MRI. They found the refractive index profiles to become flatter in the central region and steeper in periphery. Kasthurirangan et al.[107] used MRI to map the GRIN in vivo, and found a high refractive index plateau region and a decline of the refractive index towards the periphery. In the youngest age group the peripheral decline was the least steep. The axial thickness of the central plateau increased with age, but not with accommodation. In general it was found that the refractive index of the plateau region in the lens center did not change neither with age nor accommodation, but that the size of the plateau

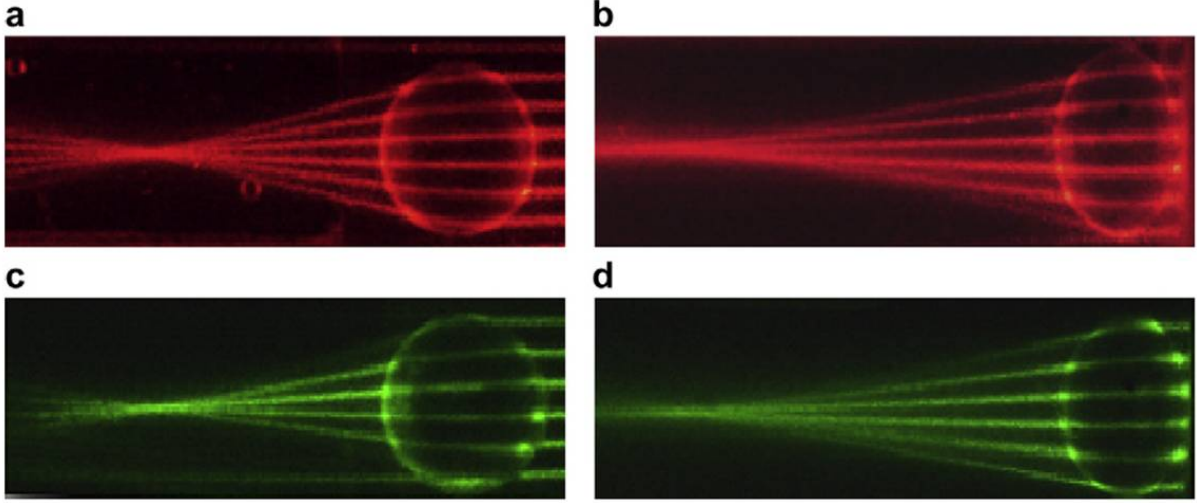


Figure 1.12: Ray tracing method to measure the GRIN (reproduced from Pierscione et al[198]). The rays paths through a porcine lens in equatorial (Figure a and c) and sagittal (Figure b and d) planes for two wavelengths (633 nm and 532 nm).

increases with age.

Optical Coherence Tomography In recent years, OCT has been used to characterize the GRIN. First experiments were done by Verma et al. on zebrafish[258]. Verma compared the OPD (optical path difference) from the OCT measurements with computed ray paths. The computed ray paths were based on the assumption of a polynomial GRIN. First OCT GRIN measurements with human donor lenses were done by Uhlhorn et al.[254]. 40 human donor lenses between 6-82 years were imaged by means of OCT, and thickness and average refractive index were measured. The measured average group refractive index was then converted to the average phase refractive index. The average refractive index was in agreement with mentioned MRI measurement (1.408 ± 0.005). Recently, de Castro et al.[47] reconstructed the GRIN of a in vitro porcine lens from 3D OCT images. He used the OPD extracted from the OCT data and the lens shape data as input parameters for a 4-variable GRIN model. The reconstruction algorithm was based on a genetic algorithm that searched for the parameters that best fit the distorted posterior surface of the lens. Subsequently, de Castro et al.[48] used 2D OCT to investigate the age-dependent variation of the GRIN profile in human crystalline lenses, using a similar reconstruction algorithm as before. Age-dependency of the GRIN was consistent with previous data using MRI. The formation of a refractive index plateau with age was also supported.

1.3.2 GRIN Models

The first time the GRIN was incorporated into an eye model was in 1924 with the Gullstrand Eye model[88]. Since then, many models were proposed to describe and calculate the GRIN. In 1971, Pomerantzeff et al.[200] built a wide angle optical eye model and simulated the GRIN distribution as a shell like structure with a high number of shells. Each shell would have a different refractive index, curvature and thickness. As input data focal length and spherical aberration were used. Almost 20 years later, Smith et al.[235] used a mathematical model based on ellipsoidal iso-indical contours that are concentric with the lens surfaces to describe the lens GRIN. Smith's model was studied by Atchison et al.[11] who described mathematically how to transform a GRIN distribution into a shell model. Within the shells, the refractive index was thought to be constant.

Al-Ahdali[3] proposed a model of the human eye, including a GRIN lens with 300 layers. He described the GRIN through the layers as exponentially increasing refractive index from the lens core to its periphery. Pérez et al.[201] used a continuous asymmetric bi-elliptical model to determine the GRIN and studied paraxial propagation of light through the lens. This work was continued in a later studies[202, 66]. Navarro et

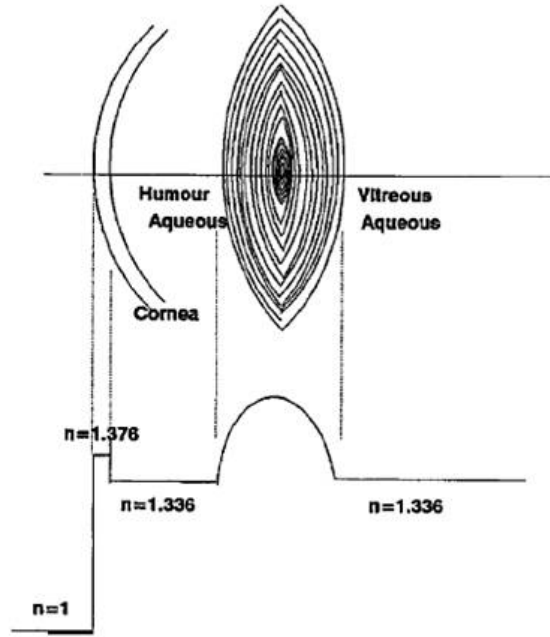


Figure 1.13: Schematic representation of the variation of the index of refraction distribution laminated lens model. Reproduced from Al-Ahdali et al[3].

al.[170] proposed a parametric GRIN model of with conicoid surfaces. The GRIN was

modeled using the power equation by Smith et al.[233]. The model was able to adapt to individual distributions and lens changes with age and accommodation. In recent years, Campbell[38] introduced a nested shell model, based on anatomy and growth of the lens throughout life. In this model, cell layers are added to an embryonic lens model. Every shell has a unique, uniform index of refraction, slightly different from its neighboring shells.

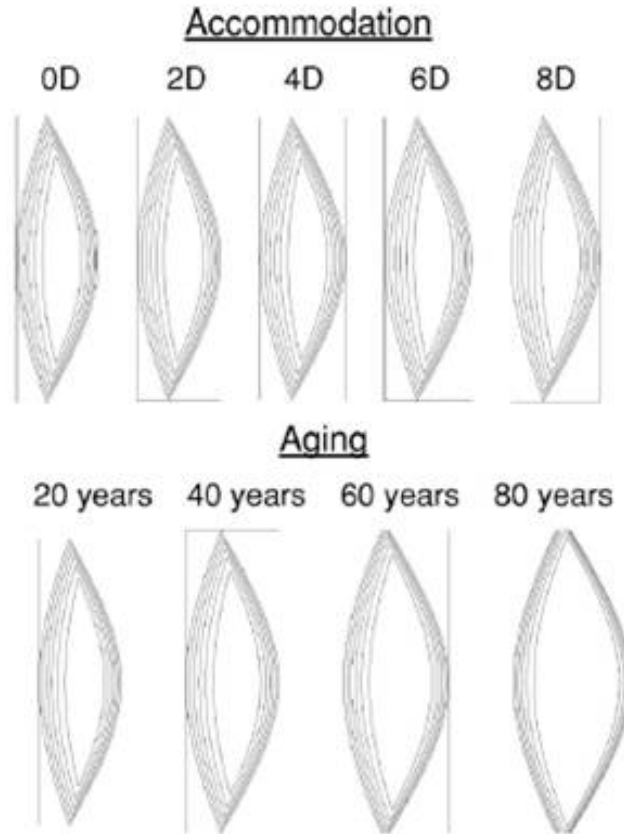


Figure 1.14: Changes of the isoindicial surfaces of the lens model by Navarro et al. with accommodation and age. Reproduction from Navarro et al.[169]

A model by Manns et al.[146] proposed a GRIN model with a power equation which described the GRIN from the center to the periphery of the lens, using conics (similar to Smith et al.[233]). This model is the basis of the one used by de Castro et al.. It was further developed for this thesis and is described in detail in chapter 2 of this thesis. Finally, Bahrami et al.[20] proposed a GRIN model where the GRIN distribution was described by a power law. The external geometry of the model was defined by a conic equation.

1.3.3 Equivalent Refractive Index

The equivalent refractive index is defined as a homogeneous lens which has the same shape and the same lens power as the corresponding GRIN lens. The expression is usually used to quantify the GRIN contribution to the optical properties of the lens. Studies for both, in vivo and in vitro lenses have shown that the equivalent index, as does the contribution of the GRIN to the lens power, decreases with age[32, 33, 57].

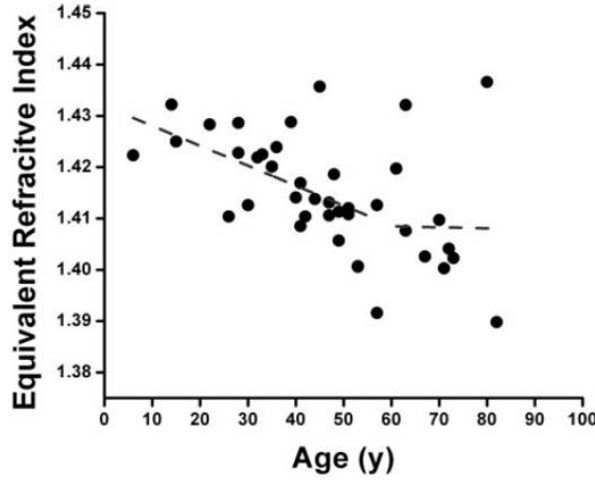


Figure 1.15: Age dependence the equivalent refractive index of in vitro human lenses. Reproduced from Borja et al.[32]

1.3.4 GRIN change with age

Several studies reveal changes of GRIN in human crystalline lenses with age, with a progressive flattening of the GRIN profile in older lenses[17, 104, 160, 89, 107, 48], likely associated to age-related changes of the lens spherical aberration with age ([32, 146, 73, 104]. Hemenger et al.[89] investigated the GRIN change in two age groups (19-31 years and 49 to 61 years) and evaluated a continuum of index values between a fixed surface and nucleus index, which were dependent on only one parameter. This GRIN parameter was found to be well separated between the age groups with a flatter GRIN near the lens center for the older age group. In already mentioned MRI studies, Augusteyn et al.[15], Moffat et al.[161] and Jones et al.[104] confirmed this behavior with age. The progressive formation of a refractive index plateau was also found by de Castro et al. using OCT imaging on human lenses with ages ranging from 6 to 72 years.

1.3.5 GRIN change with accommodation

The redistribution of the GRIN in the accommodating lens has remained relatively unexplored. Garner and Smith[70] used Purkinje-based phakometry data, in combination with a one-variable bi-elliptical GRIN model to predict, amongst others, the change of GRIN power with accommodation. The GRIN power increased from 9.70 D to 13.74 D for the maximum accommodation stimulus of 8 D. It was concluded that a substantial part of the increase in power with accommodation results from the GRIN distribution within the lens. Maceo et al.[142] investigated the GRIN contribution on stretched primate lenses. To obtain the GRIN contribution to the change in lens power, the power of the lens surfaces was subtracted from the measured lens power. The contribution of the surfaces and the GRIN increased linearly with accommodation. It was found that the GRIN contributed around 65% to the total power change during accommodation. When expressed in percent of the total power change, the relative contribution of the GRIN was constant with accommodation and age. Prior theoretical studies on the potential impact of the GRIN on accommodation point to an important contribution of the GRIN in the value[20, 50] and change of spherical aberration with accommodation[237].

1.4 Microscopic Anatomy of the human lens

1.4.1 The Lens Capsule

The lens capsule is the thickest base membrane in the human body, and appears dense and homogeneous under light microscopy, appears very dense and homogeneous. At the ultrastructure, the lens capsule consists in most part of type IV collagen combined with about 10 % glycosaminoglycan[53, 199]. The type IV collagen plays an important role in the formation of a 3D molecular network[251, 123]. The thickness of the capsule varies throughout the lens perimeter and age. The lens capsule is the thickest at the anterior midperiphery. Depending on age, it varies between 13.5 μm to 16 μm . Pre-equatorially and on the posterior pole, there are no age changes and the capsule is significantly thinner (7 μm and 3.5 μm , mean value, respectively). The average thickness at the posterior periphery decreases from 9 to 4 μm [24]. The growth of the lens capsule is of practical relevance, to adjust to the increasing volume of the crystalline lens.

1.4.2 The Lens Epithelium

The inner anterior lens capsule is in immediate contact with the basal surface of the lens epithelium. Indeed, it is the lens epithelium that synthesizes the anterior lens capsule[272, 208]. The lens epithelium consists of a monolayer of non-specialized, predominantly cuboidal cells (except the cells at the equatorial zone which are more cylindrical). The epithelium layer surfaces the anterior pole up to the lens equator. On the other side of the lens epithelium, the cells make contact with the elongated lens fiber cells. There are three different cell regions of lens epithelium, which are in non-dividing phase (central epithelium), in dividing phase (germinative) and differentiating phase (equatorial)[29]. Therefore, the size of the cells differ between $8\text{ }\mu\text{m}$ and $21\text{ }\mu\text{m}$ (in diameter) with most of the cells lying in a range from $9\text{ }\mu\text{m}$ to $17\text{ }\mu\text{m}$ [37]. The cell height has been reported to be about $5.5\text{ }\mu\text{m}$ to $8\text{ }\mu\text{m}$ [94]. With age, the cell size increases[186]. All lens epithelial cells have a large lens nucleus, together with granular cytoplasm, mitochondria, Golgi apparatus, ribosomes, rough endoplasmic reticulum (ER), and numerous small filaments[266, 94]. Unlike lens fibers, the lens epithelium has characteristics of a "normal" cell monolayer.

1.4.3 Lens Fibers and Lens Suture

Lens fibers form the main part of the lens. They are very long (up to 12 mm) and thin (typically around $7\text{ }\mu\text{m}$ wide, and $4.5\text{ }\mu\text{m}$ thick)[94]. Most fiber cells are hexagonal and stacked one upon another in well organized rows[44][206][125][231, 28, 4, 129, 130]. Their intercellular distance is only about 15-20 nm[27]. Interestingly, the average equatorial fiber width and fiber thickness does not vary much between different species. Therefore, the shape and the size of crystalline lenses are dependent on how the containing lens fibers are organized and how many fibers are part of the lens[131]. Lens fibers are divided into primary and secondary lens fibers. Primary lens fibers develop during embryonic development and become an inert structure of the lens. Secondary lens fibers are generated throughout life. They arise from the germinative zone, and elongate in direction of anterior and posterior lens surface. Finally, they lose contact with the capsule and the epithelium. They indeed unwrap the lens from both sides. This continuous addition of lens fibers throughout life leads to the changes in lens shape and thickness. As the lens grows larger, the lens fiber ends cannot reach the poles of the lens anymore, which results in a so called lens suture. Lens sutures vary between species. There are basically four distinct types of lens suture types, the umbilical, line, Y, and star, in order of increasing complexity[131]. Examples with umbilical and line suture are chicken and rabbits, respectively. Pigs and primate lenses are an example of Y sutures,

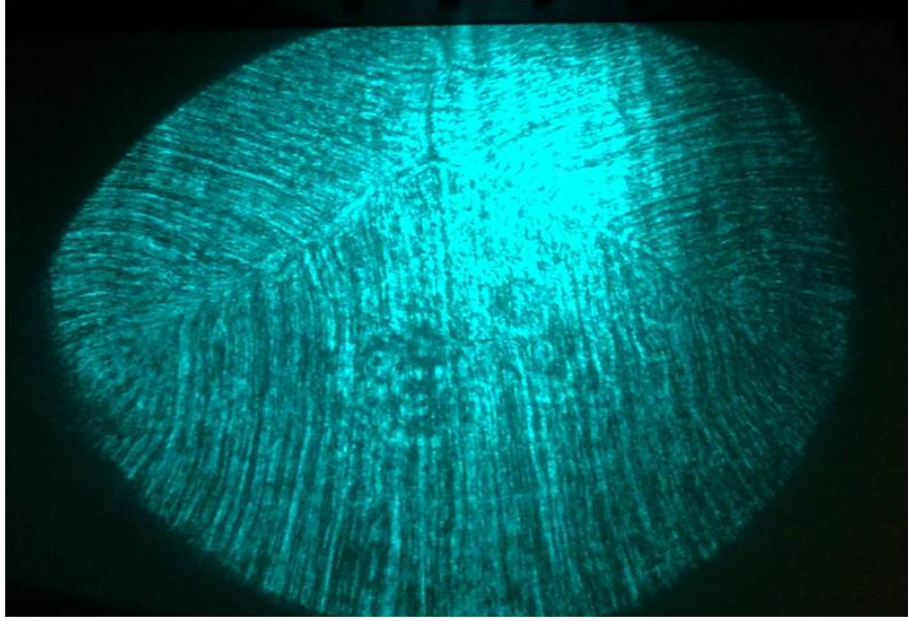


Figure 1.16: Lens fibers with Y-suture branch. This image was taken in our lab, shining 488nm light through a fresh porcine lens (Image courtesy: Maria Viñas Peña).

whereby primate lenses develop into a more complex star suture throughout life[131, 207]. This growing complexity is important because lens optical quality is directly related to lens suture type and a negative influence of sutures on lens optical quality increases with age and as a result of some ocular surgeries[127, 131, 228, 204]. The lens fibers

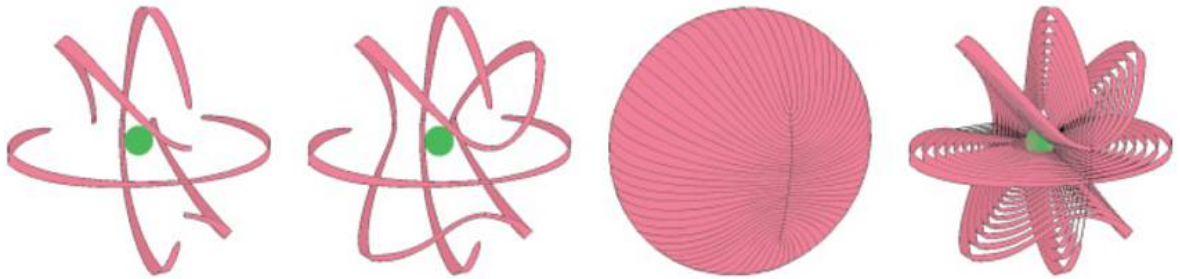


Figure 1.17: Fiber shape and arrangement in rabbit and amphibian lenses with straight and curved fibers. The typical line suture is formed. Reproduction of Kuszak et al.[131]

are linked together through various junctions, including for example ball-and-socket interlocking devices[52, 51, 128], tongue-and-groove interdigitations[132, 248, 207], and gap junctions[264, 45, 46, 76, 75, 126, 137]). The type of linking element depends on lens fiber age and where it is located within the crystalline lens. In general, more superficial

lens fibers are linked together via so called "ball-and-socket" interdigitations. Deeper fibers have additionally "tongue-and-groove" junctions.

1.4.4 Measurement of Lens Structures using Microscopy

Different microscopy imaging modalities have been used to characterize the capsule, the lens epithelium and lens fibers, including scanning electron microscopy, confocal microscopy, multi photon microscopy, or Brillouin microscopy.

Scanning electron microscopy (SEM) was used extensively to investigate the structure of lens fiber cells of different mammals, including rats[65, 83, 266], rabbits [84, 266, 267], porcine[97], monkeys[52, 132, 266, 268], and humans[44, 132, 155, 248, 253, 263]. SEM has also been used to investigate on the lens epithelium[83, 265], the lens capsule[83], cataract [51, 111, 133, 219, 253], the influence on calcium on the lens fibers[64] or on the electric coupling of lens epithelium with lens fibers[206].

Advantages of an SEM include its large depth of focus and its ability to resolve details on the surface of the sample. An obvious disadvantage for biological samples is the required preparation procedure. Samples have to be dried and heavily prepared (coated) since they must be electrically conductive. In contrast, confocal microscopy allows observation of the eye in an unstained, unfixed, living condition, confocal microscopy has been used. It provides high-resolution, high-contrast images, and the capability to section the sample without contact. The cornea has been a primary target for confocal microscopy, with commercial instruments used in the clinic for diagnostics on a regular basis [31, 101, 134, 135, 144, 152, 153, 171, 173, 184, 185, 271]. Applications of crystalline lens investigation, realized with confocal microscopy are scarce. Jeacocke et al.[102] and Xiao et al.[271] observed an isolated crystalline lens of a rabbit with a confocal microscopy. Xiao et al. could image lens fibers of a freshly excised lens about 200 μm below the lens surface. Masters et al.[150] used confocal microscopy to image the fine structure of an in-situ rabbit crystalline lens from an ex-vivo rabbit eye through the full thickness of the cornea and aqueous humour. Masters showed the feasibility of confocal microscopy to achieve high-contrast images of transparent objects across 1.7 mm (crossing cornea and aqueous humour) and conducted a comparison study between confocal microscopy and SEM to further investigate the potential of confocal microscopy for observations of pathology in the anterior segment of the eye [151]. In vitro confocal microscopy showed high resolution images of the epithelium and superficial lens fibers. The compared SEM images fully supported the confocal images, also the study showed that confocal microscopy has the potential to become a useful tool for imaging the anterior segment of the eye. The technique has been scarcely used after those first evidences.

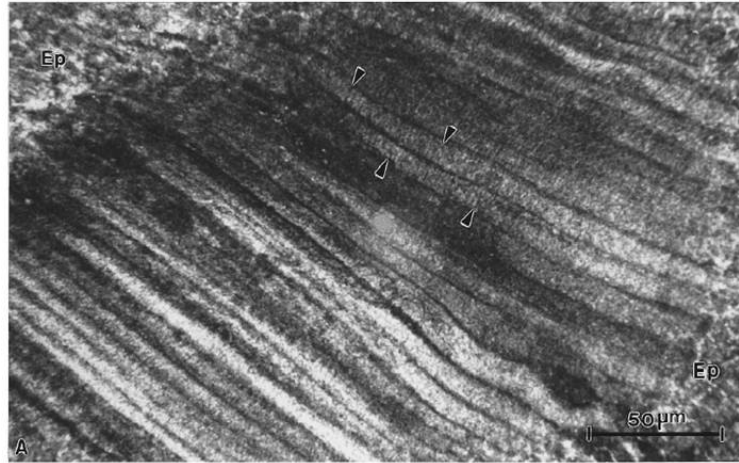


Figure 1.18: Confocal images of superficial lens fibres with epithelial cell border. Reproduction of Masters et al.[151]

More recent studies focus on somewhat specific questions or practical applications. Marcantonio et al.[148] used confocal microscopy to determine the level of cell coverage in human donor capsular bags with implanted intraocular lenses, and concluded that they contain a large population of cells, even up to 13 years after surgery.

Stachs et al.[239] analyzed and reported cutting effects of femtosecond laser cuts with confocal microscopy (Stachs1009). The authors analyzed the cuts inside of crystalline lens tissue and found altered tissue scattering properties with intact lens fibers to concrete fiber separation, depending on femtosecond laser pulse energies.

A study by Maddala et al.[143] also imaged lens fibers with a confocal microscope. The transparency of the lens depends on packing symmetry and membrane organization of the lens fibers. Concluding from his results, Maddala suggested a key role for the protein Periaxin in maturation, packing, and membrane organization of lens fiber cells. Reiss et al.[210] revealed a degenerative aging process of porcine lens tissue, using Brillouin microscopy and confocal microscopy.

An interesting experiment was done by Scarcelli et al.[222] using Brillouin microscopy. The authors characterized material acoustic properties and measured the longitudinal elastic moduli of in vivo mouse lenses over a period of two months. 3-D elasticity maps were obtained, and age-related stiffening of the lens nucleus over a period of two months. It was suggested that Brillouin microscopy has the potential to be a useful instrument for basic and animal research and clinical ophthalmology.

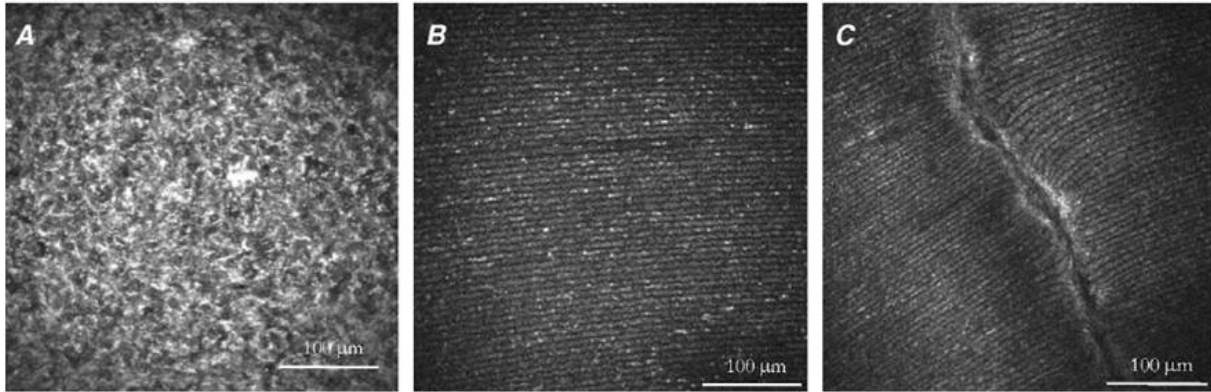


Figure 1.19: Lens capsule morphology (A), lens fibers (B), and surface suture (C) by Stachs et al., images with a confocal laser scanning microscope. Reproduced from Stachs et al[239].

1.5 Open Questions

- Although numerous studies address theoretically or experimentally the contribution of surface shape and index of refraction (generally based on an equivalent homogeneous refractive index) to the lens power, but to our knowledge there is no study that measures directly the contribution of lens shape and structure to the overall aberrations of the lens. The relative contribution of both, GRIN and lens shape, to the optical properties of the lens needs to be evaluated.
- It is well known that the spherical aberration of the young lens is negative and shifts towards positive values with age. Relative contribution of the GRIN to spherical aberration with age has never been studied experimentally.
- The role of the GRIN with accommodation is not clear. Some studies support the theory that the GRIN changes with accommodation. However, the change of GRIN with accommodation has never been measured experimentally.
- Indirect measurements of crystalline lens astigmatism suggest that its contribution to overall astigmatism can be relevant. However the relative contribution of surface astigmatism and GRIN to the crystalline lens astigmatism is not known
- There is little known about the coordination of anterior and posterior lens surface geometries and their topographical change with age
- The mechanisms of accommodation within the lens structures are still not fully understood at a microscopic level. It is beyond the scope of this thesis to investigate the fiber movements during accommodation, but current development in the laboratory using the instruments developed in this thesis address that goal.

1.6 Goal of this Thesis

The specific goals of this thesis are

- To measure the GRIN distribution in porcine, non human primate and human lenses
- To use available technologies in the Visual Optics and Biophotonics Lab (and collaborative institution Bascom Palmer Eye Research Institute) -Optical Coherence Tomography and crystalline lens stretcher- and technology newly developed for this thesis, to measure crystalline lens shape and gradient index ex vivo.
- To measure the change of the GRIN with age in donor human lenses, in 3-dimensions. To explore the relative contribution of lens shape and GRIN to the spherical aberration in unaccommodated lenses, as a function of aging and as a function of accommodation
- To explore the topography of the lens surfaces and investigate their individual contribution to the lens' optical properties like astigmatism
- To develop a custom built confocal microscope and design an experimental setup to image microscopic lens structures

1.7 Hypothesis

- The negative values of the lens spherical aberration are in parts a result from the gradient refractive index in the lens. We will test the relative contribution of the GRIN to the lens spherical aberration by evaluating both, GRIN and shape of several crystalline lenses of different species, and evaluate the influence of the GRIN by comparing the results to lenses with an equivalent homogeneous refractive index.
- The GRIN changes with accommodation. We will reconstruct the GRIN on lenses in different accommodation states, using a lens stretcher.
- The GRIN profile is age dependent. We will evaluate the GRIN dependency on human donor lenses with age and its influence on the lens optical properties with age
- The individual lens surface topographies (anterior surface and posterior surface) are to some degree correlated to each other
- It is possible to evaluate lens structures using confocal microscopy

1.8 Structure of this thesis

This thesis has been organized by chapters, which are based on published articles.

- Chapter 1: gives a state of the art and motivation of the thesis
- Chapter 2: introduces the instruments and methods used
- Chapter 3: presents a study on the influence of GRIN and shape on the lens properties of ten in vitro porcine lenses
- Chapter 4: presents the results of a collaboration with the Bascom Palmer Eye Institute (University of Miami, FL), on the change of GRIN and lens properties with accommodation
- Chapter 5: presents an extensive study on human donor lenses. The GRIN and lens shape was calculated and computational laser ray tracing was used to find out about the influence of lens properties on spherical aberration
- Chapter 6: is a continuation study of Chapter 5, in which the lens astigmatism is evaluated, using the calculated GRIN and lens shape
- Chapter 7: introduces a custom-built confocal microscope and its use to study structural properties of the crystalline lens.
- Chapter 8: includes a summary of all chapters in English and Spanish, and gives an overview of the activities done during the Ph.D. period 2010-14.

2

Methods and Materials

In this chapter the experimental and theoretical techniques are described. The imaging systems used for this work were custom-developed instruments, either at the Visual Optics and Biophotonics Lab, Instituto de Óptica, CSIC, Madrid (high speed spectral 3-D OCT, laser ray tracer, confocal microscope), or at the Ophthalmic Biophysics Center, Bascom Palmer Eye Institute, Miami, FL (time-domain 2-D OCT, Ex Vivo Accommodation Simulator II). The GRIN algorithm and the distortion correction algorithm were developed by Alberto de Castro and Sergio Ortiz, respectively, at the Visual Optics and Biophotonics Lab. The author of this thesis built the laser ray tracer and the confocal microscope, designed the experiments, and developed image processing and analysis routines (mainly in collaboration with Alberto de Castro, Jorge Lamela, Sergio Ortiz, and Bianca Maceo). Principal investigators of the laboratories were Susana Marcos, Fabrice Manns, and Jean-Marie Parel.

2.1 Experimental Techniques

2.1.1 Spectral domain 3D Optical Coherence Tomograph

The OCT system in our group was developed in collaboration with Nicolaus Copernicus University (Poland)[79]. The system (figure 2.1) is based on a fiber-optic Michelson

interferometer, an 840-nm SLD diode as illumination source (near Gaussian emission in 849nm, bandwidth 50nm, Superlum, Ireland, and a spectrometer consisting of a volume diffraction grating and a 12-bit line-scan CMOS camera (4096px) as detector. The system has an acquisition speed of 25000 A-Scans/s. The axial range is 7 mm (which corresponds to a pixel resolution of $3.42 \mu\text{m}$, and the axial resolution was calculated to $6.9 \mu\text{m}$. The imaging of the crystalline lens with the OCT results in results in different distortions, which has to be corrected to quantitatively analyze the captured images. Therefore, Ortiz et al.[177] have developed a distortion correction algorithm to correct for fan and optical distortion. Fan distortion (see figure 2.2) arises directly from the scanning architecture of the OCT (separation of scanning mirrors and aberrations of optical components), affecting the radius of curvature by about 3% and asphericity by about 125%). Optical distortion arises from viewing the internal optical surfaces through refracting preceding surfaces (for example viewing the posterior lens surface through the anterior lens surface and the lens GRIN). After distortion correction the radius of curvature and asphericity of an artificial lens could be measured with $> 90\%$ accuracy (compared to non contact profilometry)[178]. For all measurements, a special platform was used, so that the incoming laser beam would hit the horizontally positioned cuvette with the lens sample vertically (figure 2.3).

2.1.2 Time domain 2D Optical Coherence Tomograph

The OCT system is a custom built device developed at the Ophthalmic Biophysics Center at the Bascom Palmer Institute in Miami, Florida (Uhlhorn2008)(see figure 2.4). It consists of a superluminescent diode as light source (Gaussian emission centered in 825 nm, bandwidth 25 nm, SLD-38-HP, Superlum, Cork, Ireland). The axial resolution is $12 \mu\text{m}$, the maximum lateral length imaged is 20 mm. The scan length in tissue is approximately 7.5 mm. A flat scanning field is produced by a telecentric scanner as beam delivery. The depth of focus was around 10 mm.

2.1.3 Laser Ray Tracer

A Laser Ray Tracer (LRT) was custom developed in our group to measure the focal length of the lenses (see figure 2.5). Illumination came from an 849 nm SLD (Superlum, Cork, Ireland). The beam diameter at the plane of the lens was 600 m. A 2-mirror galvanometric scanning system, in combination with a 400 mm collimating lens was used to scan rings of light of 2 and 4 mm diameter onto the lens surface. A CMOS camera placed behind the cuvette captured a series of through-focus images around the focal plane of the system composed by the lens. The step size from image to image

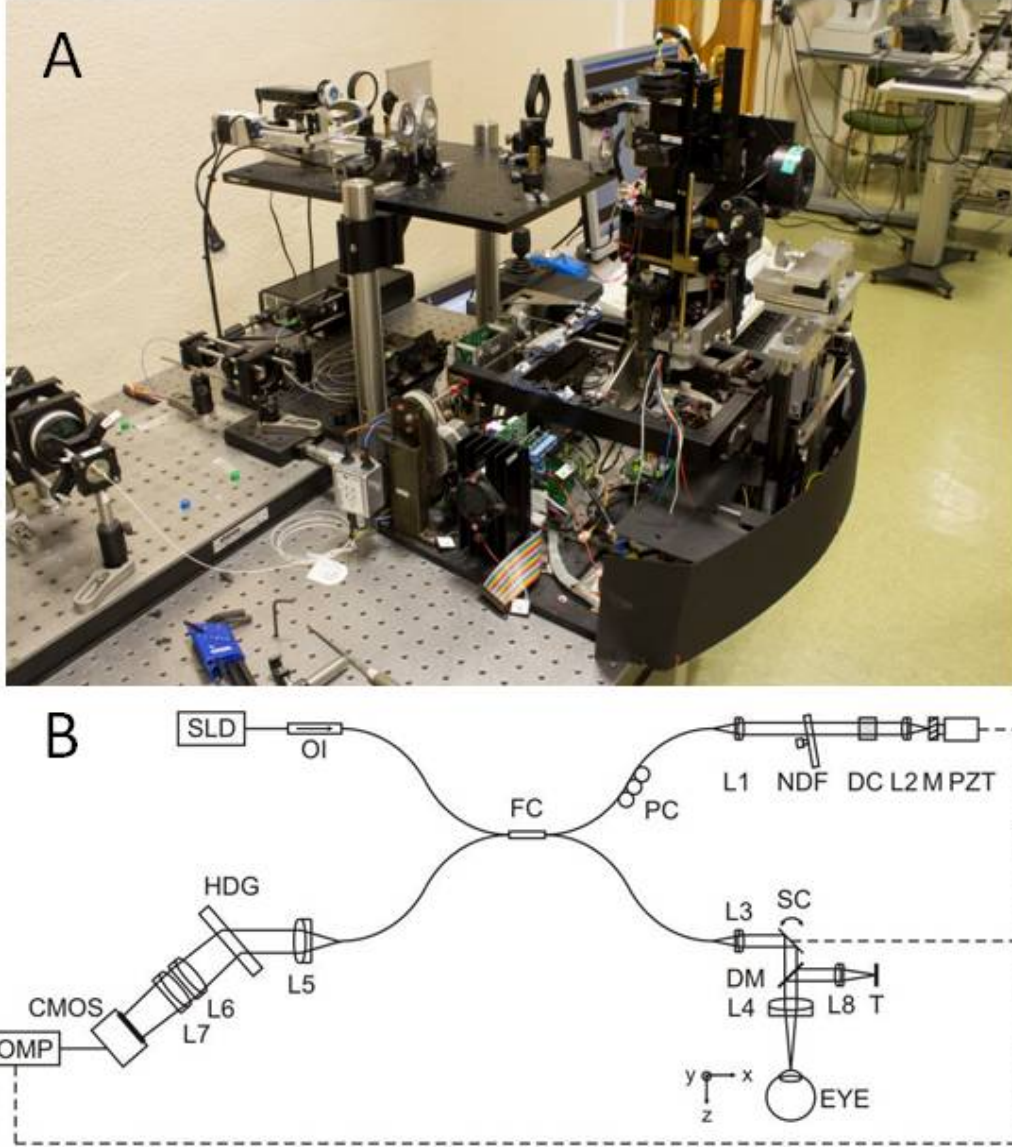


Figure 2.1: A Spectral domain 3D OCT system at the Institute of Optics "Daza de Valds"
 B Experimental SOCT set-up: SLD - superluminescent diode, OI - optical isolator, FC-80:20 fiber coupler, PC - polarization controller, NDF - neutral density filter, DC - dispersion compensator, L1-L8 - lenses, M - silver mirror, PZT - piezotranslator, SC - galvanometric scanners, DM - dichroic mirror, T - target, HDG - holographic volume diffraction grating, CMOS - linescan camera, COMP - computer. (Reproduced from Grulkowski et al. [79])

was adjustable. For our measurements, it was set to 0.5 mm. The lens was aligned such that the center of the rays captured with the CMOS camera would not change

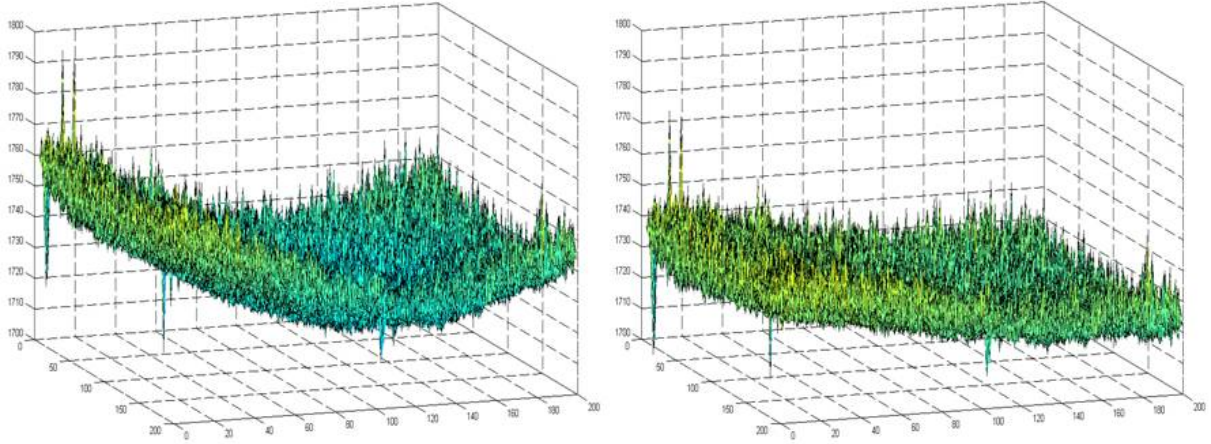


Figure 2.2: Visualization of fan distortion (Reproduced from Ortiz et al. [179])

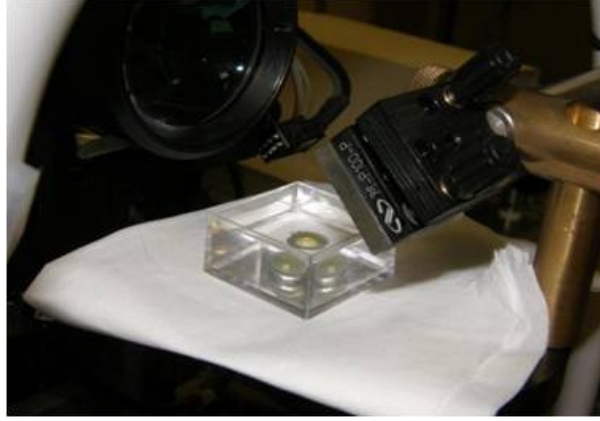


Figure 2.3: Human crystalline lens, mounted on the spectral-domain OCT at viobio group.

within the range of movement of the camera (35 mm). The system was calibrated with commercial artificial lenses (Edmund Optics Inc. Barrington, NJ USA). Comparison of the back focal length estimates from the experimental system with computational ray tracing programs (Zemax, Radiant Zemax, Redmond, WA, USA) allowed estimation of the back focal length of the lens in preservation media (see figure 6). The precision of the measurements was estimated to be 0.8 mm. The back focal length could be expressed in diopters as the inverse of the back focal length multiplied by the group refractive index of the solution $n = 1.345$ at 825 nm[34]. The captured images from the CMOS camera were analyzed with a custom developed algorithm developed in imagej to find the focal spot.

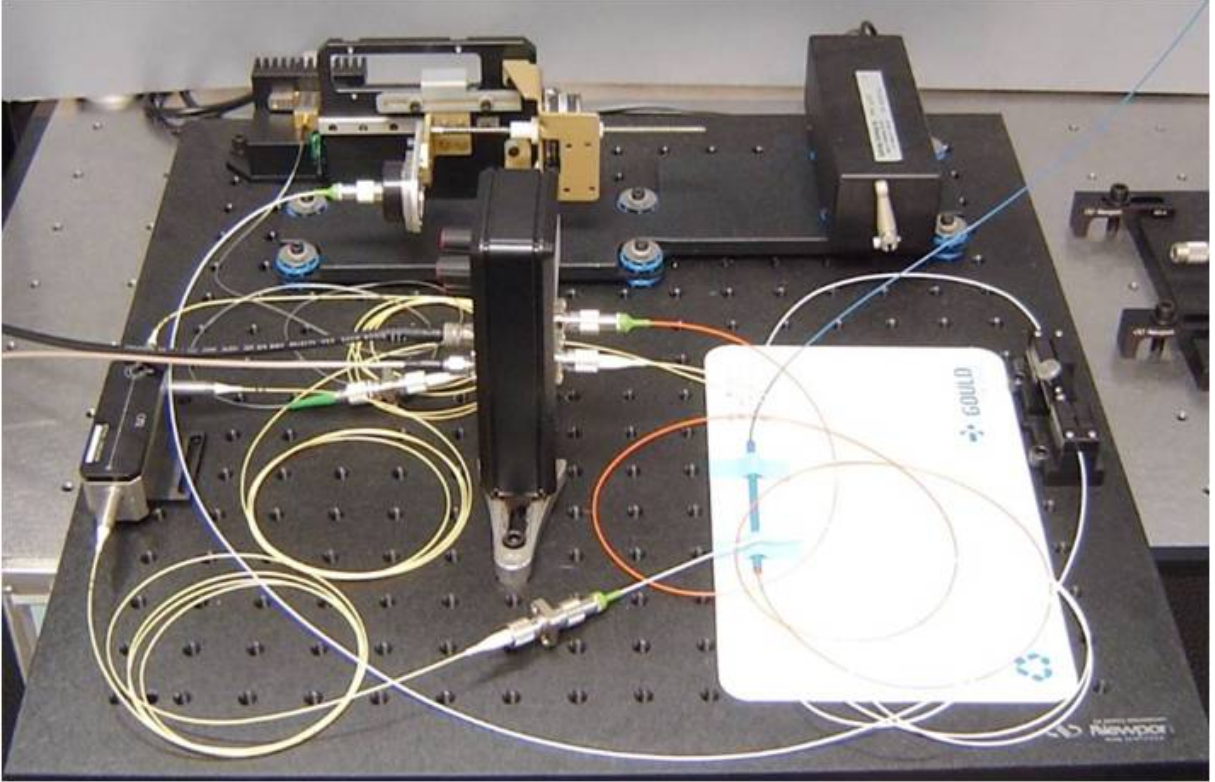


Figure 2.4: Time domain 2D OCT at the Ophthalmic Biophysics Center at the bascom Palmer Eye Institute, Miami, FL.

2.1.4 Ex Vivo Accommodation Simulator II (EVAS II)

The EVAS II reproduces dis-accommodation by simultaneous radial stretching of the eight scleral segments of the eye up to 2.5 mm radially. The following description of the stretcher is based on an article by Klaus Ehrmann, one of the designers of EVASII[62]. The stretcher consists of eight identical linear motion stages with integrated force transducer. Each stage is mounted on a motorized linear stage with five millimetres of travel (M110.1DG, Physik Instrumente, Germany). The eight stages are symmetrically arranged and connected to a computer. At the center of the instrument, a lens power measurement assembly can be inserted (0.25 inch CCD camera (GP-CX261V, Panasonic, Japan), vertically motorized). To control the stretching steps, and to receive the optical power readings, customized control software was developed. When using the stretcher, a scleral ring of the eye is glued symmetrically to eight shoes which are connected to the motion stages. The section of the scleral ring contains intact ciliary muscle, zonules and the crystalline. The posterior section of the globe, as well as the cornea and the iris, are removed.

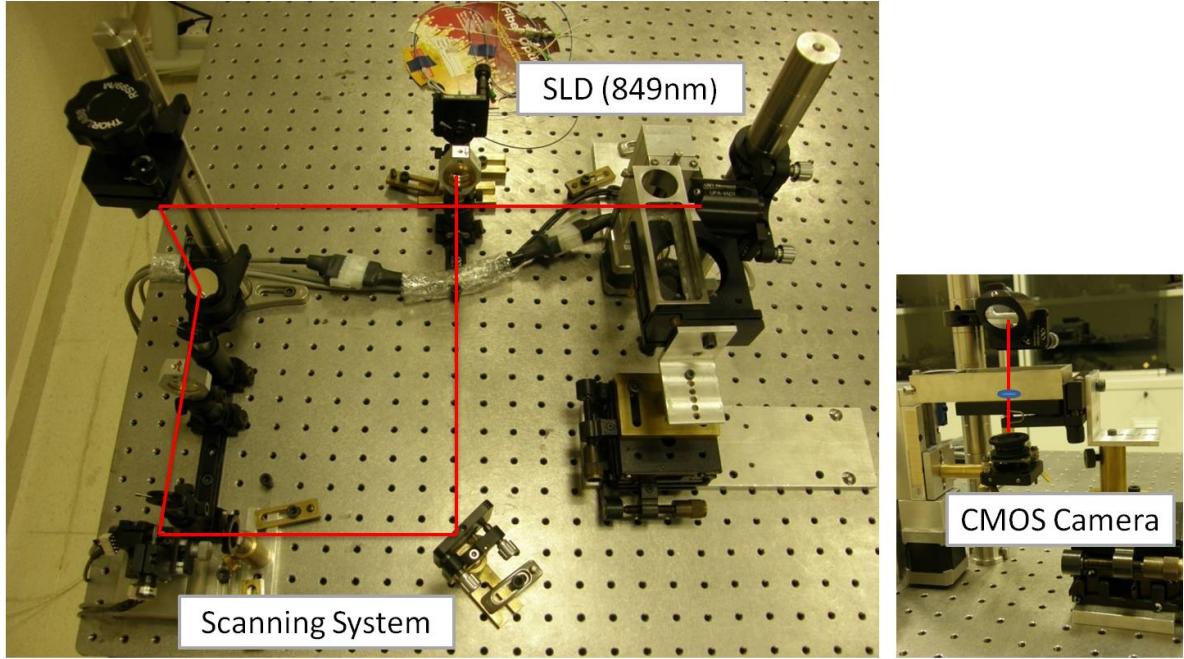


Figure 2.5: Custom built LRT. The larger photo shows the SLD light passing the scanning system and various mirrors so that a ring of light can be projected onto the lens surface (smaller photo). The CMOS camera captures images around the focal plane of the lens.

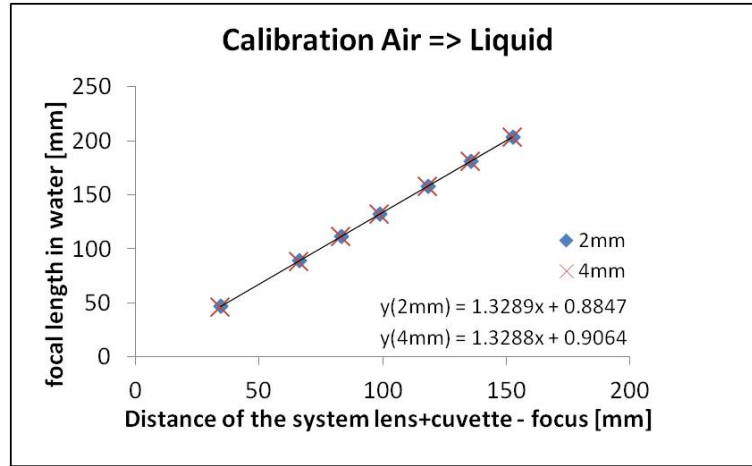


Figure 2.6: Dissected human eye, bonded to stretching shoes, shown in fully stretched position. (Reproduced from Ehrmann et al. [62])

2.1.5 Confocal Microscopy

The confocal microscope used was custom-developed in our group. It is part of a Multiphoton-Confocal combination-microscope. The set-up and layout is shown in fig-

ure 8. The confocal part of the microscope is equipped with a diode laser at 488 nm and a avalanche photodiode (APD) as detection channel. The laser source goes through a telescopic system to clean the beam. It passes through a beam splitter, reaches the scanner, and is then directed to a water dipping objective. The microscope objective was a Olympus LUMPIanFL 40x waterimmersion with a numerical aperture of 0.8 (Olympus LUMPIanFL 40x NA 0.8). The beam expander between the scanning system and the objective is important to expand the laser beam to the numerical aperture of the objective, and to transport the image plane to the APD. The microscope has a lateral resolution of $0.5\mu\text{m}$, an axial resolution of $3.5\mu\text{m}$, an axial range of approximately 1 mm. The acquisition rate can be up to 100.000 A-Scans/s. The pinhole is placed in front of the APD in a conjugate plane with the focal plane of the microscope. It acts like a spatial filter, thus only the signal coming from the focal point reaches to the detector, removing all signals from out of focus. There is a close relationship between the pinhole size and the axial resolution of the system, that means it is important to use a pinhole as small as possible, which implies that the beam must be focused as much as possible. As a consequence, it is not possible to illuminate a large area of the sample, but rather a very small, specific region. It is therefore important that the system includes a scanning method that allows to scan the sample point by point, to obtain an image with a considerable size that allows to obtain information about the sample and not only local information.

2.1.6 Eyes

porcine eyes

Cadaver porcine eyes (around 6 months of age) were obtained from a local slaughterhouse and transported at a temperature around 4°C (see figure 2.8). Experimental protocols had been approved by the Institutional Review Boards.

rabbit eyes

Rabbit eyes (age unknown) were obtained from a local slaughterhouse and transported at a temperature around 4°C . Experimental protocols had been approved by the Institutional Review Boards.

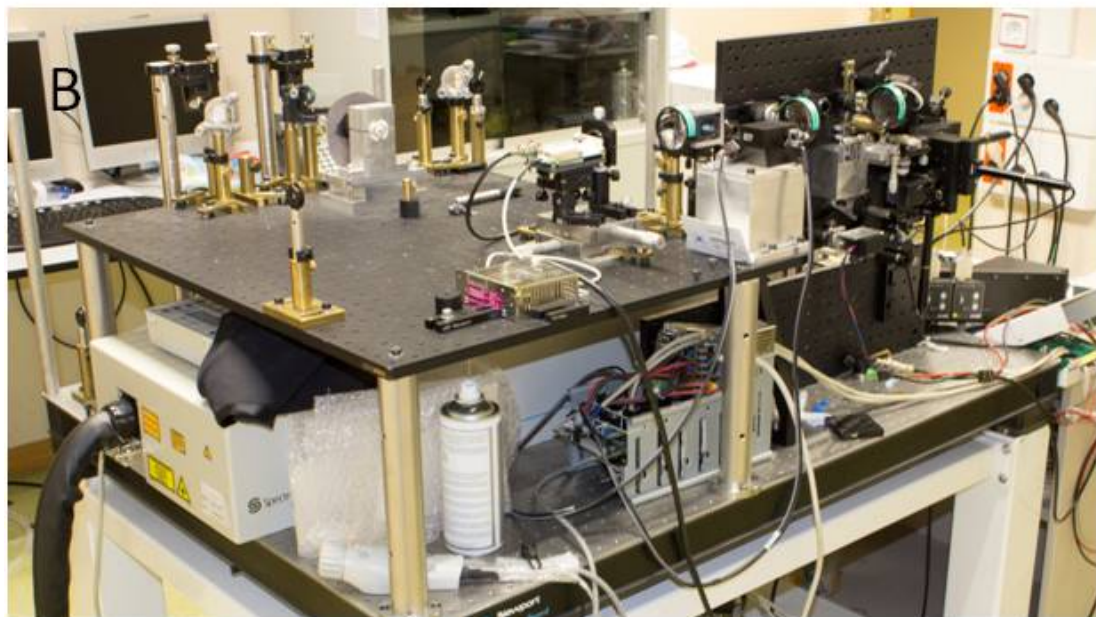
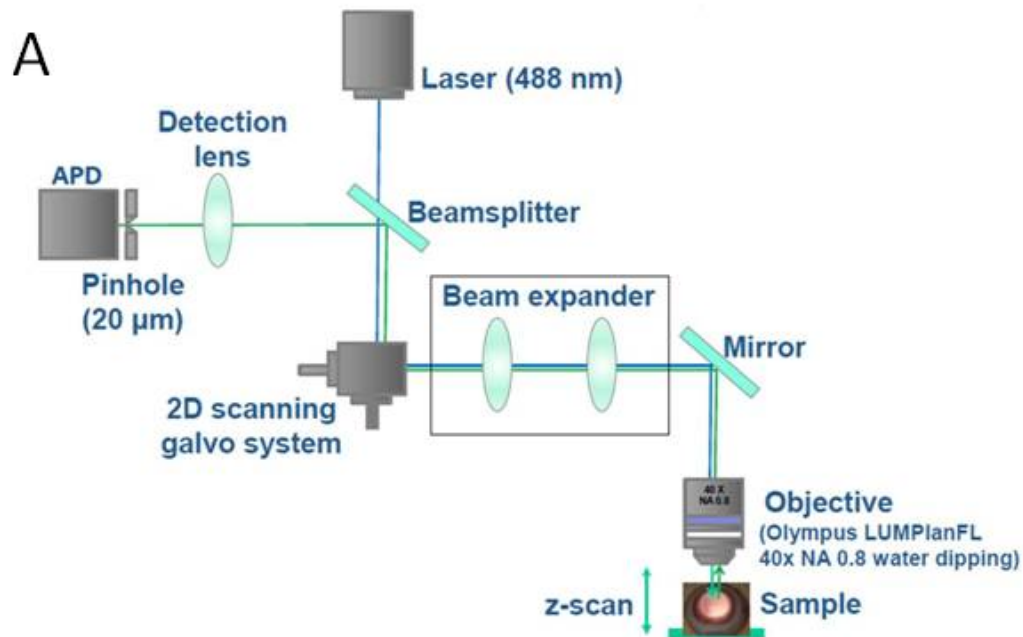


Figure 2.7: A Layout of the custom-developed confocal microscope, showing the incoming beam (dark blue) and the outgoing beam (light blue). B Confocal-Multiphoton Microscope at the Institute of Optics.

cynomolgus eyes

Experiments with cynomolgus lenses were executed at the Ophthalmic Biophysics Center at the Bascom Palmer Institute (Miami, FL). All experiments adhered to the Associ-



Figure 2.8: Albino rabbit eye, 2 hours post-mortem

ation for Research in Vision and Ophthalmology Statement for the Use of Animals in Ophthalmic and Visual Research. The eyes were obtained from the Division of Veterinary Resources at the University of Miami as part of a tissue-sharing protocol and were used in accordance with Institutional Animal Care and Use Guidelines. No animals were euthanized for the sole purpose of this study.

Human donor eyes

Human donor eyes were obtained from Transplant Service Foundation (TSF) Eye Bank. Methods for securing human tissue were in compliance with the Declaration of Helsinki. The handling and experimental protocols had been previously approved by the Institutional Review Boards of TSF and CSIC. Presence of cataract was considered an exclusion criterion in the study. During the experiments, all damaged and swollen lenses were also excluded (see figure 2.9). All eyes were shipped in sealed vials at 4°C, wrapped in gauze soaked in preservation medium (DMEM/F-12, HEPES).

2.1.7 Data Analysis

All OCT images were processed to obtain the shape of the lens surfaces and the surface of the cuvette. Custom-developed software was used to merge the images of the lens, obtained in two different planes, and to detect both surfaces of the lens, as well as the cuvette distorted by the lens and preservation medium (see figure 2.10). All images were corrected from fan distortion[179] and optical distortion (in most cases, the distortions due to the preservation medium were corrected to extract the geometry of the first surface of each image)[47]. All surfaces (lens and cuvette) were fitted by Zernike polynomials (up to 7th order) within a 6-mm pupil. For the purposes of this study, only symmetric

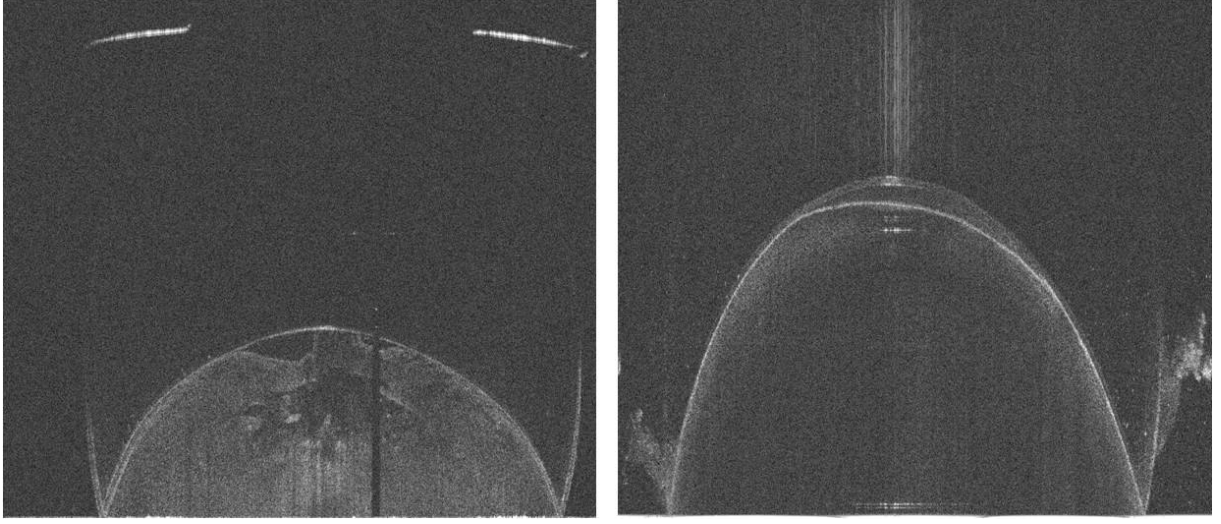


Figure 2.9: OCT image of a damaged lens (left) and a swollen lens with detached lens capsule (right), which were excluded from the study.

Zernike polynomials and astigmatism were used. Additionally, the radii of curvature and conic constant of the surfaces were estimated for 18 meridians, using the following equation

$$z = z_0 - \frac{(x - x_0)^2}{r + \sqrt{r^2 - k(x - x_0)^2}} \quad (2.1)$$

where z is the surface sag, x the radial position along the meridian, r the apical radius of curvature at the vertex, and k the conic constant. The lens thickness was calculated from the distortion of the image of the cuvette surface[254].

2.1.8 GRIN model

The refractive index in the lens was modeled so that it varies continuously from the nucleus to the surface in both, axial, and meridional directions. The center of the lens is assumed to be in the optical axis at a distance from the anterior surface vertex equal to 0.41 times the lens thickness [218]. The GRIN is described as a 4-variable model in polar coordinates:

$$n(\rho, \theta) = n_N - \Delta n \cdot \left(\frac{\rho}{\rho_S} \right)^{p(\theta)} \quad (2.2)$$

where n_N is the refractive index of the nucleus, Δn the difference between the refractive index of surface and nucleus, ρ_S is the distance between nucleus and surface, and $p(\theta)$ is the exponential decay in axial (p1) and meridional (p2) direction; p1 is constant across

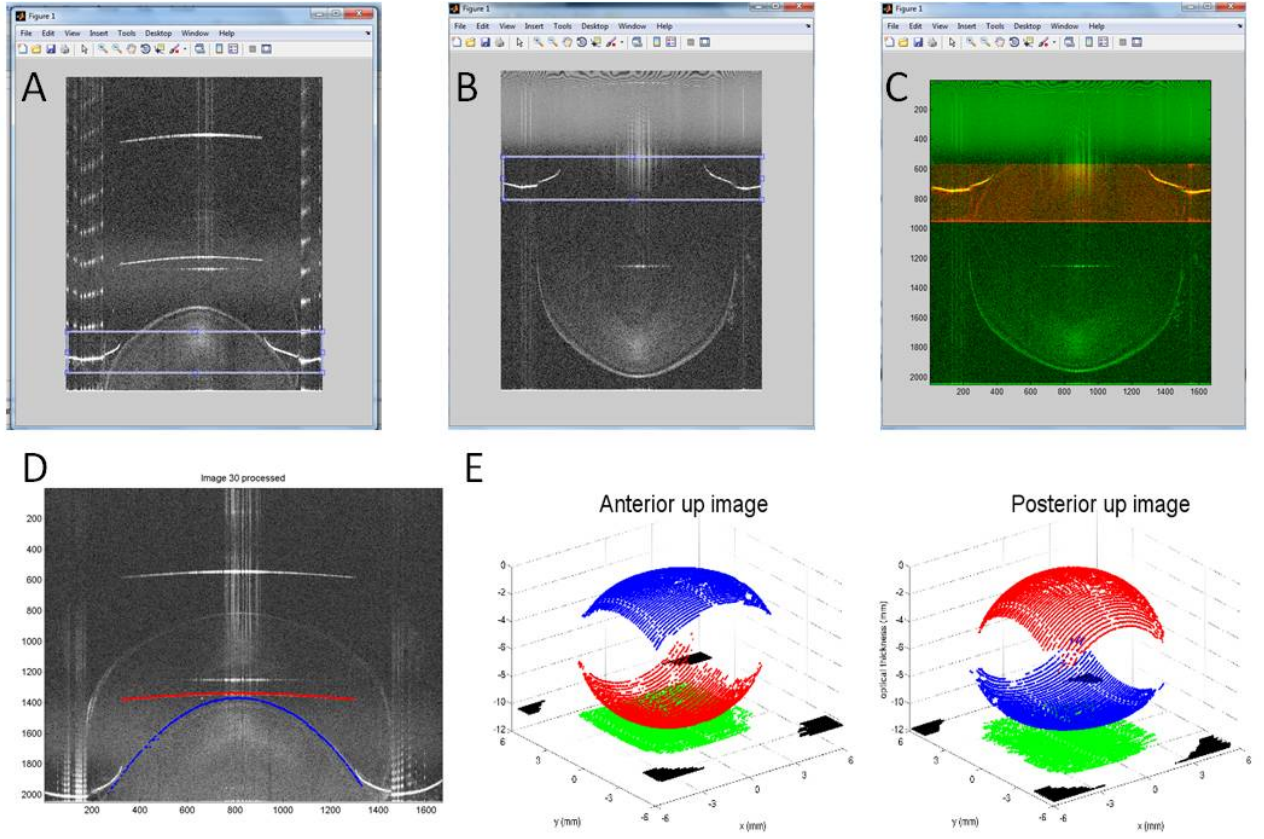


Figure 2.10: Image merging ($A+B=C$) and surface detection (D) to obtain 3D lens surfaces (E).

meridians, while p_2 can vary to account for differences between meridians [47].

2.1.9 Search algorithm

The GRIN distribution which best fits the experimental data is searched through minimization of a merit function which is defined by the sum of the root mean square of the differences between the simulated and the experimental OPD from the OCT images and the simulated and experimental focal lengths of the lenses. Since the search involves variables that are strongly coupled it is probable that a local algorithm would get stuck in local minima. Therefore, a genetic optimization algorithm⁴⁸ was implemented as a global search algorithm to prevent the solution falling in a local minimum [47].

2.1.10 Computational Ray Tracing and Zernike fitting

To investigate the relative effect of the lens geometry and GRIN distribution on the aberrations of the lens, a computational ray tracing analysis was performed at multiple meridians, based on the experimental lens parameters retrieved for each meridian of the lens. The custom developed ray tracing algorithm was programmed in MatLab (Math-Works, Natick MA). The basic components of the algorithm are the implementation of the Stavroudis formula to trace rays through conical surfaces [241], the Sharma algorithm for ray tracing in the GRIN media [223], and a procedure to calculate the impact of the ray in the posterior surface of the lens [242]. Aberrations were calculated for a 6 mm pupil diameter by fitting the wave front up to a 7th order Zernike polynomial expansion using the reconstructed 3D GRIN distribution or the equivalent refractive index. The equivalent refractive index, defined as the homogeneous index of a lens with the same external geometry and focal length as the crystalline lens, was calculated. Astigmatism and spherical aberration coefficient (obtained from the wavefront fitting) were compared across lenses, assuming a homogeneous equivalent refractive index or the estimated GRIN. The astigmatism was calculated with the Zernike polynomials using the power vector defined by Thibos et al. [249].

2.1.11 Power Vector Analysis

Lens GRIN astigmatism

The lens total astigmatism was calculated using the measured lens shape and the reconstructed GRIN by means of a computational ray tracing analysis[47]. The calculated wave aberrations were fit by Zernike polynomials and the magnitude of astigmatism C and angle α were calculated in [250] as:

$$C = -2\sqrt{J_{45}^2 + J_{180}^2} \quad (2.3)$$

$$\alpha = \frac{1}{2} \tan^{-1} \left(\frac{J_{45}}{J_{180}} \right), \quad (2.4)$$

where J_{180} is the power at axis $\alpha = 0^\circ = 180^\circ$ and J_{45} is the power at axis $\alpha = 45^\circ$, and

$$J_{180} = -\frac{2 \cdot \sqrt{6} \cdot Z_2^{-2}}{r^2} \quad (2.5)$$

$$J_{45} = -\frac{2 \cdot \sqrt{6} \cdot Z_2^2}{r^2} \quad (2.6)$$

where Z_2^{-2} and Z_2^2 are the corresponding astigmatism Zernike terms, and r the pupil radius (3 mm in this study).

Surface Lens Astigmatism

The lens surface astigmatism C_S was calculated as

$$C_S = (n_2 - n_1) \cdot \left(\frac{1}{R_x} - \frac{1}{R_y} \right) \quad (2.7)$$

where R_x is the astigmatic angle and R_y is the axis perpendicular to R_x . n_1 and n_2 are the refractive indices of the lens immersion medium and the lens cortex, respectively.

2.1.12 Software

Matlab

Matlab (The MathWorks, Natick, MA) was used for image processing, surface detection, GRIN reconstruction, computational ray tracing, and data analysis.

Amira

Amira (FEI Visualization Sciences Group, Bordeaux, France and the Zuse Institute Berlin (ZIB), Germany) is a software platform for 3D and 4D data visualization. For this thesis, it was used for image processing, image analysis, and visualization of data obtained with the confocal microscope.

Imagej

Imagej (National Institute of Health) is a Java-based image processing program. For this thesis it was used for Image processing and image alignment.

Zemax

Zemax (Radiant Zemax) is an optical design program. It is used for design and analysis of imaging and illumination systems.

3

Contribution of the GRIN and Shape to the Crystalline Lens Spherical Aberration and Astigmatism

This chapter is based on the paper by Birkenfeld et al. "Contribution of the gradient refractive index and shape to the crystalline lens spherical aberration and astigmatism" (Vision Research, 2013, 86, 27-34). The coauthors of this study were Alberto de Castro, Sergio Ortiz, Daniel Pascual, and Susana Marcos. The author of this thesis (i) designed the Laser Ray Tracer (in collaboration with Susana Marcos) (ii) built the Laser Ray Tracer (in collaboration with Sergio Ortiz and Daniel Pascual), (iii) modified the experimental setup (OCT) (iv) did the sample preparation (v) performed the experiments (in collaboration with Alberto de Castro) (vi) performed data processing (in collaboration with Alberto de Castro), (vii) analyzed the results (in collaboration with Alberto de Castro) (viii) prepared the manuscript (in collaboration with Alberto de Castro and Susana Marcos) . As a result of this work, we could perform the first systematic experimental study of the relative contribution of geometry and GRIN to the aberrations in a mammal lens. It was found that the presence of GRIN shifted the lens' spherical aberration towards negative values, and, in addition, produced a decrease in the total amount of lens astigmatism in most lenses.

3.1 Introduction

The optical properties of the crystalline lens are determined by its shape and refractive index distribution. However, to date, those properties have not been measured together in the same lens, and therefore their relative contributions to the lens' optical aberration are not fully understood. In this chapter, we measured the shape, the optical path difference, and the focal length of ten in vitro porcine lenses (age around 6 months) using Optical Coherence Tomography and a custom built laser ray tracer, both described in chapter 2. The 3D Gradient Refractive Index distribution was then reconstructed by means of the optimization method. The optimization method searched for the parameters of a 4-variable GRIN model that best fits the distorted posterior surface of the lens in 18 different meridians. Spherical aberration and astigmatism of the lenses were estimated using computational ray tracing, with the reconstructed GRIN lens and an equivalent homogeneous refractive index. Although there is extensive literature that studies theoretically [30][235] or experimentally [32][57][73][104][117][142][160][254] the contribution of surface shape and index of refraction (generally based on an equivalent homogeneous refractive index) to the lens power, to our knowledge there is no study that measures directly the contribution of lens shape and structure in the overall aberrations of the lens. The only systematic experimental analysis, which investigated the role of the GRIN on spherical aberration we are aware of is on fish lenses. Fish lenses are known to have very low amounts of spherical aberration, despite being nearly spherical in shape. Jagger[100] compared the expected spherical aberration of an artificial spherical lens with homogenous index of refraction with that of the fish lens, and hypothesized that the presence of a GRIN distribution in the lens balanced the lens surface spherical aberration. Kroeger et al. [122] demonstrated the relative contribution (and balance) of the lens surface and GRIN on spherical aberration in African teleost fish lenses. A more recent study used Optical Coherence Tomography in combination with an iterative optical path fitting to estimate the GRIN distribution in Zebrafish [258].

3.2 Methods

3.2.1 Lens Specimens and preparation

Ten cadaver porcine eyes (around 6 months of age) were obtained from a local slaughterhouse and transported at a temperature around 4°C. Experiments were performed within three to eight hours after enucleation. Immediately before the measurements, the eye globes were cut and the crystalline lens was carefully dissected and placed in a cuvette

filled with a preservation solution (DMEM/F-12, HEPES, no phenol red, Invitrogen, GIBCO) [19]. The overall duration of the experiment was about one hour per lens. Experimental protocols had been approved by the Institutional Review Boards.

3.2.2 Optical Coherence Tomography system

The lenses were imaged in 3D using a custom developed OCT system described in chapter2. All images were provided with fan distortion correction algorithms that arise from the scanning architecture of the instrument [177]. The lens axis was aligned with the OCT scan axis such that a specular reflection was seen from the surfaces of the lens. Care was taken that the lens did not appear tilted in the image preview, for neither 0 and 90 degrees. One image contains 1668 A-Scans and 70 B-Scans on a 12x12 mm lateral range. Each set of 3D images consisted of two acquisitions at two different planes of focus which were merged to produce a complete image of anterior surface, posterior surface, and the cuvette holding the lens, using customized Matlab programs[47]. The acquisition time for each OCT image was 4.5s

3.2.3 Laser Ray Tracing

Our custom developed LRT system, described in chapter2, was used to measure the focal length of the lenses for two different pupil diameters. The beam diameter at the plane of the lens was 600 μm . A 2-mirror galvanometric scanning system, in combination with a 400 mm collimating lens was used to scan rings of light of 2 and 4 mm diameter onto the lens surface. A CMOS camera placed behind the cuvette captured a series of through-focus images around the focal plane of the system composed by the lens immersed in preservation media and the cuvette. The step size from image to image was set to 0.5 mm. The lens was aligned such that the center of the rays captured with the CMOS camera would not change within the range of movement of the camera (35 mm). The precision of the measurements was estimated to be 0.8 mm. The back focal length could be expressed in diopters as the inverse of the back focal length multiplied by the group refractive index of the solution $n = 1.345$ at 825 nm [34]. To estimate spherical aberration from the experimental power measurements it was assumed that the focal length measured for a circle of light with a radius of $r=1$ mm corresponds to the paraxial focal length. The 4th order Zernike spherical aberration Z_4^0 was then obtained using

$$Z_4^0 = \frac{M r^2}{12 \sqrt{5}} \quad (3.1)$$

where M is the difference in power, experimentally determined between the paraxial (2 mm diameter) and the non paraxial ring (4 mm diameter), and r is the pupil radius, in our measurements 2 mm[247].

Experimental Protocols

All lenses were imaged in the cuvette immersed in preservation medium, first with the OCT system and then with the Laser Ray Tracing (LRT) apparatus. OCT 3D imaging was performed with the posterior surface lens facing the OCT beam (posterior up image). The lens was then flipped around a predetermined horizontal axis (the same for all lenses) to image it again, this time with the anterior surface facing the OCT beam (anterior up image). The procedure is similar to that described in prior publications [34][47]. After completing the OCT imaging, the lens in the cuvette was aligned in the LRT with the anterior surface facing up. A through focus series of 10-15 images was collected with the CMOS camera to estimate the back focal length of the lens.

OCT Image Data Processing

The OCT images were used to obtain the shape of the anterior and posterior lens surfaces, the thickness of the lens and the optical path accumulated by the rays passing through the lens and arriving to the cuvette, which served (along with the focal length measurements) as input parameters in a GRIN reconstruction algorithm. The image processing algorithms are described in chapter 2. The data were corrected from fan distortion [179]. The distortions due to the preservation medium were corrected to extract the geometry of the first surface of each image. The detected edges of the distorted lens -i.e. subject to the optical distortion produced by refraction by the preceding lens surface [177] and GRIN [224]-, the undistorted surfaces and the cuvette were fitted by Zernike polynomials (up to 7th order) within a 6 mm pupil. However, since the ray tracing algorithm is defined assuming conic surfaces, only symmetric Zernike polynomials and astigmatism were used to implement the lens surfaces in the algorithm. The radii of curvature and conic constants were calculated for 18 meridians. Previous studies using the same instrument estimate an accuracy in radius of curvature estimates of 1% in glass lenses[177]. Some of the variability is associated to interdependencies between radius of curvature and conic constant in a conic fitting [203]. Maximum and minimum radii of curvature and conic constants across meridians are reported. The lens thickness was calculated using the method proposed by Uhlhorn [254].

GRIN Reconstruction Algorithm

The GRIN algorithm described in chapter 2 was used to estimate the 3D GRIN distribution within the central 6 mm pupil diameter of the lens. Input parameters were lens shape and thickness, as well as the optical path obtained from the OCT images. A novelty of the new algorithm, in comparison to the version used by de Castro et al.[47] was the use of focal length data (for 2-mm and 4-mm ring diameters) as additional input parameters. The GRIN was reconstructed searching for the parameters of the refractive index distribution which best fits the experimental data. The procedure was applied in 18 lens meridians, from 0° to 170° in steps of 10°. The algorithm was run five times for each data set. The group refractive index at 849 nm was reconstructed (since the OCT system uses a broadband source) and was then converted to phase refractive index at 633 nm [254] for comparison with previous publications.

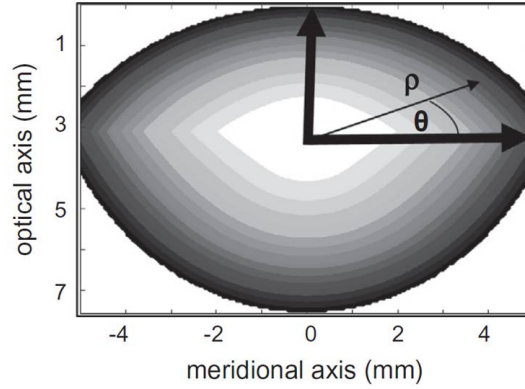


Figure 3.1: The GRIN was described in each meridian by four variables; nucleus and surface refractive index, and exponential decay in optical and meridional axis. The center of the lens was set to 0.41 times the central thickness of the lens, and the refractive index in each point was calculated in polar coordinates with equation (3.2)

GRIN model

The 4-variable GRIN model as described in chapter 2 was used (see figure 3.1):

$$n(\rho, \theta) = n_N - \Delta n \cdot \left(\frac{\rho}{\rho_s} \right)^{p(\theta)} \quad (3.2)$$

The axial direction p1 was constrained to have the same value in all meridians while the meridional direction p2 could vary to account for differences between meridians. The

center of the GRIN was defined to be in the optical axis at a distance from the anterior surface vertex equal to 0.41 times the lens thickness [218].

Search Algorithm

The global search algorithm, described in detail in chapter 2, is based on a genetic optimization algorithm[95], which avoids the solution falling in local minima. The GRIN parameters that best fit the experimental values are searched by minimizing a Merit function. The Merit function is defined by the sum of the RMS (root mean square) of the differences between simulated and measured OPD (optical path difference) for the posterior surface of the crystalline lens and the cuvette[47], and the difference between the simulated and measured back focal length for the two rings of light projected on to the lens.

Computational ray tracing analysis

To investigate the relative effect of the lens geometry and GRIN distribution on the aberrations of the lens, a computational ray tracing analysis was performed at multiple meridians, based on the experimental lens parameters retrieved for each meridian of the lens. The custom developed ray tracing algorithm was programmed in MatLab (MathWorks, Natick MA). Aberrations were calculated for a 6 mm pupil diameter by fitting the wave front up to a 7th order Zernike polynomial expansion using the reconstructed 3D GRIN distribution or the equivalent refractive index. Astigmatism and spherical aberration coefficient (obtained from the wavefront fitting) were compared across lenses, assuming a homogeneous equivalent refractive index or the estimated GRIN. The astigmatism was calculated with the Zernike polynomials using the power vector defined by Thibos et al. [249][250], as described in chapter 2.

3.3 Results

3.3.1 Lens Geometry

All the meridians of the lens surfaces were well fitted with conic sections. The residuals were below 100 μm over the entire central 6 mm pupil. Figure 3.2 shows the measured radii of curvature of the anterior and posterior lens surfaces (from anterior and posterior up lens OCT images, respectively) for all lenses. The posterior lens surface is systematically steeper than the anterior surface. A statistically significant correlation between

anterior and posterior radius of curvature, reported in previous studies [259] was not found. The astigmatism of anterior and posterior lens surfaces was not aligned, i.e. the

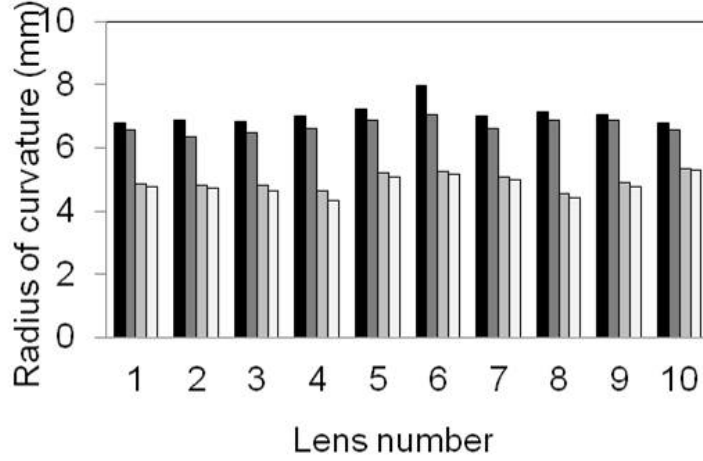


Figure 3.2: Radius of curvature of the 10 studied porcine lenses. Data are maximum and minimum radii of curvature from fits to 18 meridians

axis of astigmatism of posterior surface was rotated with respect to the axis of astigmatism of the anterior surface. The amount of astigmatic axis rotation differed across lenses: in five out of ten lenses the rotation was around 45, and in two of them nearly 90. This crossed astigmatism between anterior and posterior lens surfaces is consistent with previous observations in vitro ([47], porcine lenses). and in vivo ([176], human lenses). Figure 3.3 shows the anterior and posterior lens asphericity (conic constant $k = Q + 1$) for all lenses (average across meridians). The values were in the positive range without exception describing the lenses surfaces as prolate ellipsoids ($0 < k < 1$) or oblate ellipsoids ($k > 1$) [58]. Figure 3.4 shows the estimated central thickness of the lenses. Lens thickness ranged from 7.56 mm to 7.89 mm.

Laser Ray Tracing

The back focal length of the lenses was studied using a laser ray tracing system which projects circles of lights of different diameters on the lens. The measured back focal length in the preservation medium was on average 21.7 ± 0.8 mm (ranging from 20.4 to 22.7 mm) and 22.0 ± 0.8 mm (ranging from 20.8 to 22.9 mm) for the 2 mm and 4 mm diameter rings of light, respectively, corresponding to a power of 61.9 ± 2.4 D and 61.0 ± 2.1 D in solution, respectively. In nine out of ten lenses, a shorter back focal length for a smaller diameter was found. The shorter back focal length for a smaller diameter is consistent with a negative spherical aberration in the crystalline lenses. Using equation

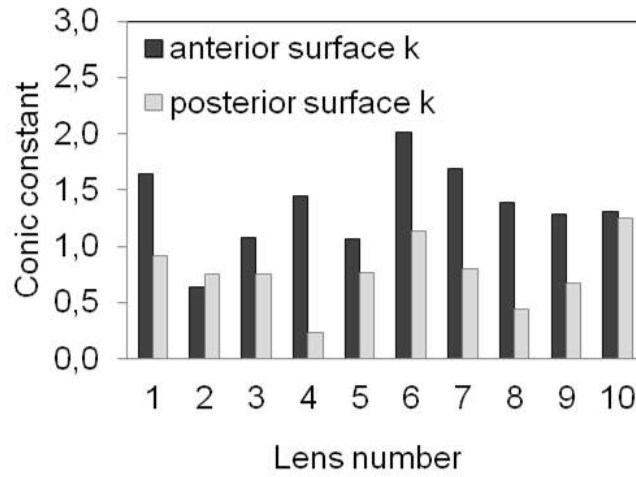


Figure 3.3: Asphericity of the 10 studied porcine lenses. Shown is the conic constant $k=Q+1$ for all lenses. In all cases k was positive for both the anterior and posterior lens surface.

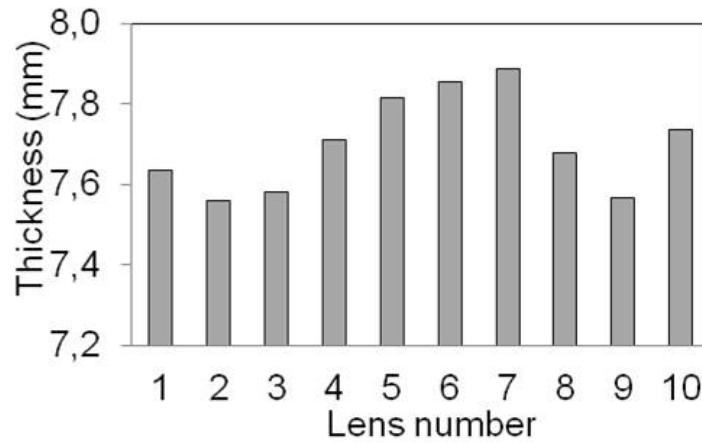


Figure 3.4: Calculated lens thickness for all lens samples.

(1), the fourth order spherical aberration Z_4^0 was calculated to a mean value of $Z_4^0 = -0.137 \mu\text{m}$.

3.3.2 Gradient Index Distribution

Refractive Index of Nucleus and Surface

The reconstructed refractive indices of the lens nucleus and surface are shown in figure 3.5 for all eyes (at a phase refractive index at 633 nm). On average, the refractive index varied from 1.427 in the nucleus to 1.364 in the periphery (the corresponding values at 849 nm are 1.436 and 1.374 for nucleus and periphery, respectively). The standard deviations across repeated reconstructions (averaged across lenses) were 0.0019 for the nucleus, and 0.0030 for the periphery. The standard deviations of the refractive indices across lenses were 0.004 and 0.003 for nucleus and surface, respectively, indicating low differences in the refractive index values across lenses. Figure 3.6 shows examples of

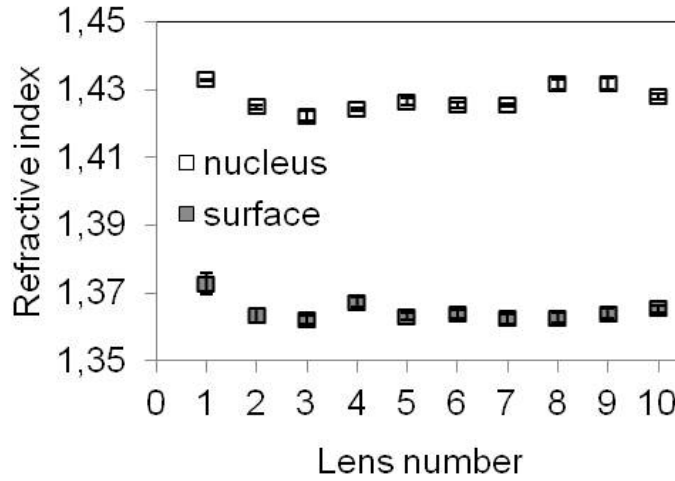


Figure 3.5: Nucleus and surface refractive index for the 10 porcine lenses in the study. The standard deviation across repeated reconstruction is indicated for each individual lens. Averaged across lenses the standard deviation was 0.0019 for the nucleus and 0.0030 for the lens surface.

the index profile for 2 example lenses. The decay along the optical axes is shown for posterior (short dashed line) and anterior (long dashed line) part of the lens. Also shown is the distribution of the decay along multiple meridians (solid line). The meridional variations of the index relate to the meridional GRIN profile differences. The inset shows the corresponding refractive index distribution for one meridian. Across all lenses, the exponential decay of the GRIN in the optical axis (p_1 in eq. (2)) varied from 1.42 to 1.98, and the meridional decay change (p_2 in equation (3.2)) from 1.64 to 5.52 with a change of 1.3 as a mean between meridians. The low value of the exponential decays

means that there is no plateau built up along the axis, this is, that the gradient refractive index is well distributed across all lenses.

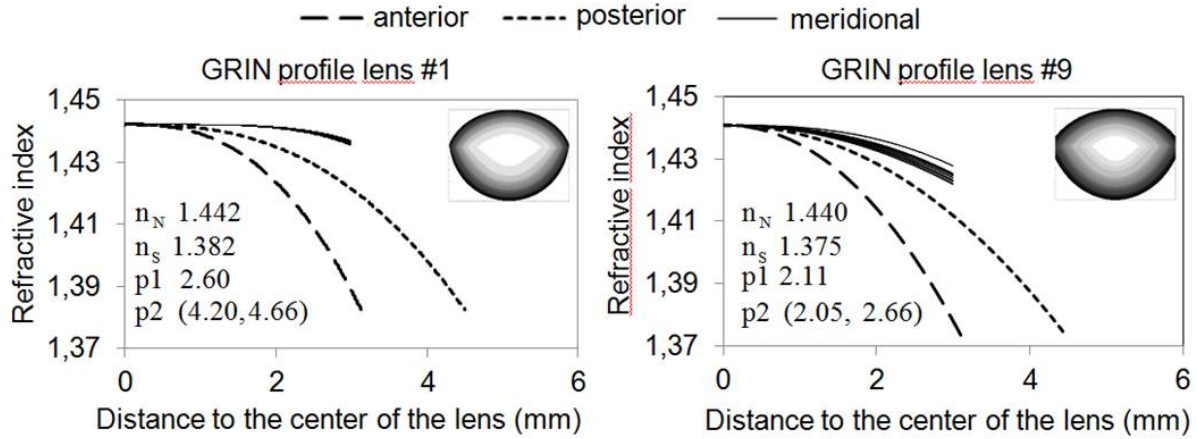


Figure 3.6: Refractive index profile from nucleus to surface along the optical axis and along the meridional direction. Insets show the corresponding reconstructed values of nucleus and surface refractive indices (n_N and n_s) and the exponential decay factor $p1$ and $p2$ from equation (3.2). On the upper right side, the GRIN distribution map is shown for a single meridian.

Equivalent refractive index

The mean equivalent refractive index found across lenses was 1.449 ± 0.003 . The standard deviation of the results of the five repetitions of the search algorithm was below the standard deviation across all lenses. Spherical Aberration Figure 3.7 shows the 4th order spherical aberration for all lenses (6 mm pupil diameter), calculated from the measured lens shape and both the reconstructed GRIN distribution and the equivalent refractive index. Results are shown for both the corresponding equivalent refractive index and the measured GRIN distribution. Spherical aberration was consistently positive with the homogeneous index ($1.17 \pm 0.21 \mu\text{m}$) and shifted towards negative values in the presence of GRIN ($-0.71 \pm 0.67 \mu\text{m}$). GRIN shifted spherical aberration between $1.12 \mu\text{m}$ and $2.7 \mu\text{m}$. A negative spherical aberration was found in 8 out of the 10 lenses.

Astigmatism

Figure 3.8 shows the amount of astigmatism in the lens for an equivalent refractive index lens and the corresponding GRIN lens. Sources of astigmatism in the lens are the

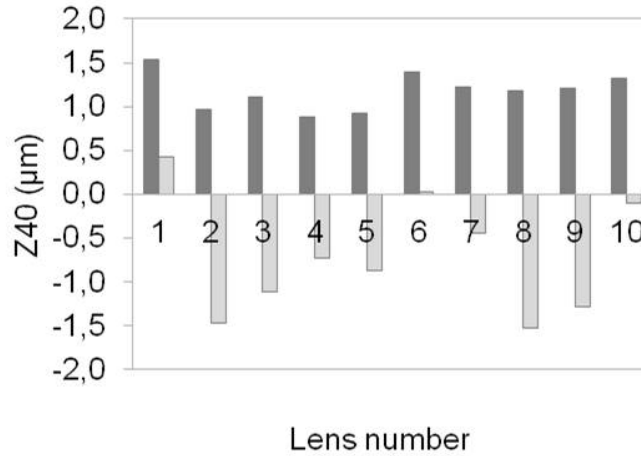


Figure 3.7: Fourth order spherical aberrations calculated with the reconstructed GRIN (light gray bars) and the homogeneous equivalent refractive index (dark gray bars). The presence of GRIN shifts the spherical aberration towards negative values. The shift induced by the GRIN varies between $1.12 \mu\text{m}$ and $2.7 \mu\text{m}$.

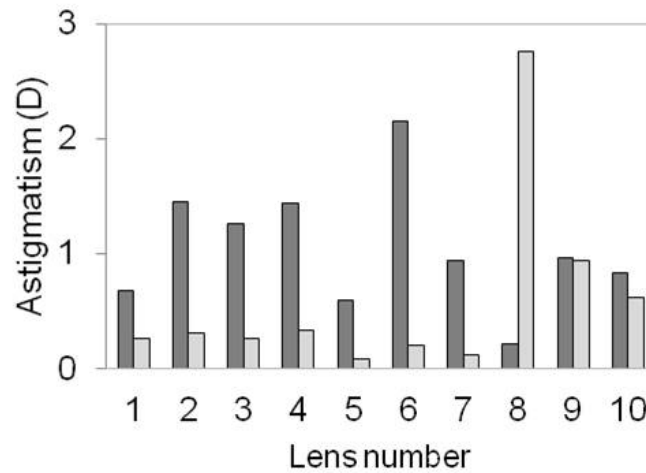


Figure 3.8: Amount of astigmatism for lenses with a homogeneous equivalent refractive index and the corresponding GRIN. In most cases the absence of a GRIN implied a higher amount of astigmatism.

astigmatism of each of the surfaces as well as the meridional variations in the GRIN. In most lenses the presence of GRIN produced a decrease in the total amount of astigmatism in the lens. Only in one lens (#8) astigmatism increased with GRIN, whereas in two lenses (#9 and #10) GRIN seems to play a minor role in lens astigmatism.

Figure 3.9 shows the estimated axis of astigmatism with a homogenous index and GRIN.

The presence of GRIN induces very little rotation in the angle of astigmatism (4.11 deg on average). Except for lens #10, the rotation induced by the GRIN is less than 5 deg.

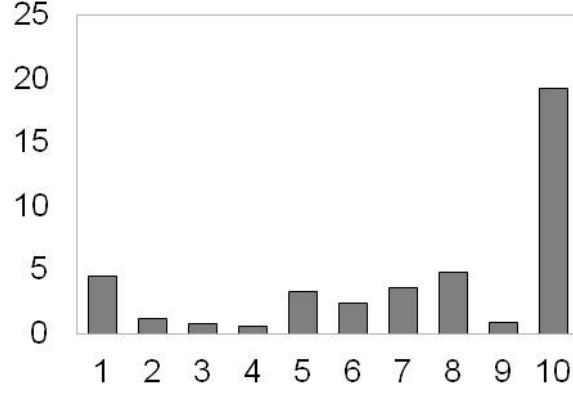


Figure 3.9: Change of astigmatic angle induced by the GRIN. The values are less than 5 deg in almost all cases.

3.4 Discussion

In this chapter, we characterized the geometry, GRIN and optical aberrations of several young porcine crystalline lenses in vitro. To our knowledge, this is the first systematic experimental study of the relative contribution of geometry and GRIN to the aberrations in a mammal lens, and the first time that the study of the crystalline lens geometry and GRIN has been addressed in three dimensions. Experimental measurements of GRIN in mammals are scarce. Vazquez et al. [256] reported 2D GRIN profiles in porcine lenses, using a tomographic reconstruction technique based on lateral ray tracing. Their reported index values ranged from 1.366 to 1.444 (mono-polynomial model) and 1.361 and 1.449 (bi-polynomial model) for the surface and the nucleus refractive index, respectively. In a previous publication, de Castro et al.[47] found 1.362 and 1.443 with a similar reconstruction method to that presented in the current study, but based on OCT images only. Those data are in close agreement with our findings (1.364 and 1.427 for

surface and nucleus indices) using both OCT and focal length as input measurements. Vilupuru et al. studied the biometrical and optical properties of the isolated pig lenses using a scanner laser apparatus [259]. The lens radii of curvature values in our current study (6.5-7.5 mm for the anterior lens and 4.5-5.3 mm for the posterior lens) fall within the ranges reported by Vilupuru et al. (5.5-7.5 mm and 5-7 mm, respectively). Unlike in their study, no significant correlations between the dimensions of the anterior and posterior lens were found, most likely because of the limited age range of the subjects in this study and, therefore, more similar lens geometry across the sample (all eyes were around 6 months of age in our study, compared to 3-8 years in the previous study).

To our knowledge, asphericities in the porcine lens have never been reported. Interestingly, unlike the primate lens, which tends to show negative asphericity in the lens surfaces in vivo and in vitro [32][57], we found positive and negative asphericities in the anterior and posterior surfaces. And unlike the great intersubject variability in lens asphericities found in vivo [57], we found relatively low intersubject variability in the asphericity measured in this group of pigs. The focal length measured in this study, ranging from 20.4 to 22.7 mm and from 20.8 to 22.9 mm for the 2 mm and 4 mm diameter rings of light, respectively), seemed in general lower than the ones by Vilupuru et al. with laser scan measurements who reported a wide range of values (around 23-34 mm) [259]. Contrarily, a study by Jones et al. [105] measured the focal length of six different porcine lenses (age unknown) with both, optical methods and MRI measurements, and reported in average a shorter focal length (19.0 ± 1.2 mm for ray tracing and 19.6 ± 0.77 using the GRIN estimated with MRI) than the one in the present study. The differences may arise from the mentioned differences in age. The results from the ray tracing measurements showed negative spherical aberration in most of the lenses. As expected, since these data were used as input values in the GRIN reconstruction, the ray tracing estimates of spherical aberration based on the lens shape and GRIN are consistent with those negative values. The calculated spherical aberration using the reconstructed GRIN was $-0.71 \mu\text{m}$ (for a 6 mm pupil diameter). The lower value found from laser ray tracing ($-0.137 \mu\text{m}$) is consistent with the smaller sampled lens area in the experimental measurements (4 mm diameter). The calculated value based on ray tracing computations is similar to the spherical aberration reported, using laser ray tracing [213] on a single porcine lens (around $-1 \mu\text{m}$ for a 6-mm pupil diameter) and the results using a point-diffraction interferometer [2] on 12 porcine lenses (around $-1 \mu\text{m}$ for a 6 mm pupil diameter). Also, most porcine lenses (16 out of 20) studied by Vilupuru et al. showed negative spherical aberration but the data showed a high variability [259]. The reported mean spherical was equivalent to a much higher Zernike spherical aberration value ($-2.5 \pm 3.35 \mu\text{m}$) than our results. An earlier study, in which beams from a He Ne laser of varying separations were directed through crystalline lenses of a number of

vertebrates, found pig lenses to be almost free of aberration [226]. We have found that the GRIN appears well distributed within the porcine lenses of the study (low exponential decay values in the GRIN distribution), more similar to the young human lens than to the plateau distribution of the old human lens [48, 104]. Our results show that the GRIN distribution plays an important role in the negative spherical aberration of the crystalline lens in pig. The shift of spherical aberration towards negative values with the GRIN in pig lenses had been suggested by [255] and was also apparent in the experimental study by Wong et al. [269] on polymer refilled pig lenses. Wong et al. reported that the spherical aberration of a lens refilled with a homogeneous refractive index gel was positive, while it was negative before refilling. This compensation of the surface positive spherical aberration by the GRIN is well known in spherical and nearly spherical lenses such as in fish [100][122] where the lens spherical aberration is close to zero. In fish, potential corneal aberrations are irrelevant due to the water immersion. In humans, the negative spherical aberration of the crystalline lens serves to compensate the positive spherical aberration of the cornea [9][23]. This was also reported in primates [213], and may be also the case in pigs. Similarly to spherical aberration, compensatory effects of corneal astigmatism by lenticular astigmatism have been previously reported [9][110]. Population studies conclude that the crystalline lens compensates about 0.5 D of corneal astigmatism in human eyes [209]. As for spherical aberration, the balance of astigmatism tends to be lost with age, by a prevalent shift of corneal astigmatism from with-the rule to against the rule with age [221]. While the internal astigmatism is generally evaluated indirectly from corneal keratometry or corneal topography and ocular refractions, the sources of lenticular astigmatism have generally not been investigated. By measuring the lens shapes and GRIN we have been able to analyze the contribution of the astigmatism of the anterior and posterior lens surfaces and the GRIN distribution to the astigmatism of the porcine lens. Interestingly, the magnitude of the astigmatism was lower in the lens with the reconstructed GRIN than in the lens with the homogeneous index in 8 out of 10 lenses, although the presence of GRIN generally did not entail a relevant change in astigmatism axis. In summary, in this chapter we have investigated the contribution of the gradient refractive index and shape to the crystalline lens spherical aberration and astigmatism for the first time in a mammal lens. Understanding the sources of these aberrations in the human lens is important to gain further insights on the image forming properties of the lens, and is of particular interest in the development of intraocular lenses. Aspheric and toric designs aim at mimicking the compensatory role of the corneal aberrations by the young lens. Better designs can be envisioned by gaining deeper knowledge of the strategies followed in natural lenses, including different animal species.

4

Influence of shape and GRIN in the Accommodative Changes of Spherical Aberration in Non-human Primate Crystalline Lenses

This chapter is based on the paper by de Castro et al. "Influence of shape and Gradient Refractive Index in the Accommodative Changes of Spherical Aberration in Non-human Primate Crystalline Lenses" (Invest Ophthalmol Vis Sci. 2013 Sep 11;54(9):6197-207). The author of this thesis was first coauthor. Other coauthor included Bianca Maceo, Fabrice Manns, Esdras Arrieta, Jean-Marie Parel, and Susana Marcos. The author of this thesis (i) modified the experimental setup, (ii) performed the experiments (in collaboration with Bianca Maceo and Esdras Arrieta), (iii) modified existing Matlab code (in collaboration with Alberto de Castro), (iv) performed data processing (in collaboration with Bianca Maceo and Alberto de Castro), and (v) reviewed the manuscript. This work allowed us to investigate lens shape and GRIN of in vitro monkey lenses under simulated accommodation. As a result it was concluded that the reconstructed GRIN lens has more negative spherical aberration and a larger shift toward more negative values with accommodation.

4.1 Introduction

In the last chapter we measured the shape, the optical path difference, and the focal length of ten porcine lenses using 3D Optical Coherence Tomography and a custom built laser ray tracer. With this data we could reconstruct the 3D Gradient Refractive Index distribution (GRIN) and use computational ray tracing to estimate spherical aberration and astigmatism of the GRIN lenses and lenses with an equivalent refractive index. All lenses were deattached from any zonular tension, therefore, they were in their most "accommodated state". Since porcine lenses only accommodate at a very young age, this term should be understood as a state "without any zonular tension".

In this chapter, we are using in vitro cynomolgus monkey crystalline lenses (which naturally accommodate in vivo), and mimic accommodation experimentally using a computer controlled motorized stretcher.

We present the experimental measurements of the lens geometry and of the reconstructed GRIN profiles as a function of accommodation. The contribution to the spherical aberration and its change with accommodation is calculated by means of computational ray tracing through a lens model, mentioned in the previous chapter.

The aberrations of the eye are known to change with accommodation[141][10][87][42][68], and these changes are related to modifications of the shape and internal structure of the crystalline lens. Several studies have reported the shape of the surfaces and the optics of the accommodating crystalline lens[35][116][58]. However, although of critical importance to understand the optical changes of the lens with accommodation, the relative contribution of the crystalline lens shape and GRIN to the change in the optical properties that the lens undergoes with accommodation is not known.

In addition, the redistribution of the GRIN in the accommodating lens has remained relatively unexplored. Garner and Smith[70] used Purkinje-based phakometry data, in combination with a one-variable bi-elliptical GRIN model to predict the change of the lens focal length with accommodation. A more recent study attempted the use of MRI in vivo[107] to study the changes in axial and equatorial GRIN profile modeled by power functions. Recent studies have investigated the lens power change with accommodation, and the role of the GRIN in cynomolgus monkeys and hamadryas baboons, assuming a value for the outer cortex refractive index [33][142]. The relative contribution of GRIN to lens power appears to remain constant with accommodation[33][142][70]. This suggests that a homogenous index material in lens refilling procedure would be equally efficient in producing a refractive power change (assuming identical lens shape changes) than the natural lens GRIN material.

To contribute to a deeper understanding of the role and relative importance of the

GRIN on the optics of the crystalline lens, this study served to experimentally explore the GRIN redistribution with accommodation in non-human primate lenses and studies its influence in the spherical aberration of the eye.

4.2 Methods

4.2.1 Donor Tissue

We studied 15 young cynomolgus lenses (*Macaca fascicularis*) with ages between 3.0 and 7.3 years old (average 5.7 ± 1.1 years). All experiments adhered to the Association for Research in Vision and Ophthalmology Statement for the Use of Animals in Ophthalmic and Visual Research. The eyes were obtained from the Division of Veterinary Resources at the University of Miami as part of a tissue-sharing protocol and were used in accordance with Institutional Animal Care and Use Guidelines. The eyes were enucleated immediately after euthanasia, wrapped in gauze, and placed in a closed container. No animals were euthanized for the sole purpose of this study. Upon arrival at the laboratory, all eyes were either directly prepared for stretching experiments or refrigerated at 4 °C before testing[168]. The time between euthanasia and use was 11 ± 14 (range 1, 48) hours in this study.

4.2.2 Tissue Preparation

The tissue preparation protocol has been described in chapter 2 and in previous literature[147][62]. In summary, the sclera was bonded on the eight segments of the stretching device using cyanoacrylate. The segments fit 1 mm posterior to the limbus to the equator of the eye. The posterior calotte of the eye was dissected and posterior vitreous was carefully removed to leave the anterior vitreous and hyaloid membrane and ciliary body untouched. The tissue section was then transferred on a Petri dish placed on a retro-illuminated station positioned under the operation microscope, the cornea was dissected at the level of the limbus, meridional incisions were made in the sclera between the mobile segment to produce eight independent segments and the iris was removed. The tissue section was then transferred to Ex Vivo Accommodation Simulator II (EVAS II)[62]. During the dissection and the EVAS II testing experiment, the tissue was immersed in a chamber filled with Dulbecco's Modified Eagle Medium (DMEM)[19].

4.2.3 Stretching

The EVAS II reproduces dis-accommodation by simultaneous radial stretching of the eight scleral segments up to 2.5 mm radially. For this study, steps of 0.25 mm, resulting in 11 accommodative steps, were used except for two of the lenses where only 6 accommodative steps were measured.

4.2.4 OCT Imaging

Cross-sectional images of the crystalline lens were obtained with a custom designed time-domain OCT system that uses a superluminescent diode with a nearly Gaussian spectrum, a specified central wavelength of 825 nm and a bandwidth of 25 nm. The system has an optical scan depth of 10 mm and an axial resolution of 12 m in air (16 m in the preservation medium). Images were recorded with 5000 points/A-line and, 500 A-lines/B-scan, and a lateral scan length of 10 mm[254].

4.2.5 Lens Back Vertex Power Measurements

The lens back vertex power was measured, for each stretching step, using a method described in chapter 2 and previous literature[142][32]. Essentially, the OCT light source was used to project a collimated circular ring of 1.5-mm radius on to the center of the lens. Below the cuvette holding the lens and preservation medium, a CCD camera mounted on a vertical translation stage was used to locate the focus. The mean position resulting from three stretching runs was used in a paraxial optical model to calculate the back vertex power of the lens in diopters (D). The accuracy of the measurements was determined to be ± 0.5 D by calibration tests using a set of glass lenses.

4.2.6 Experimental Protocols

To obtain lens shape and avoid the optical distortions, testing was first performed with the lens anterior surface facing the OCT (anterior up, figure 4.1A) and then the segment was inverted with the posterior segment facing the OCT (posterior up, figure 4.1B). Following the stretching runs to image the lens, more stretching experiments were performed to measure the back vertex power of the lenses both anterior-up and posterior-up. In all the measurements, the lenses were centered using the OCT system such that the specular reflection produced by the lens apex[254] was seen from the anterior or posterior surfaces of the lens. By placing the segments on EVAS II in the same orientation for

both anterior and posterior surfaces up, the same cross-section of the lens was analyzed.

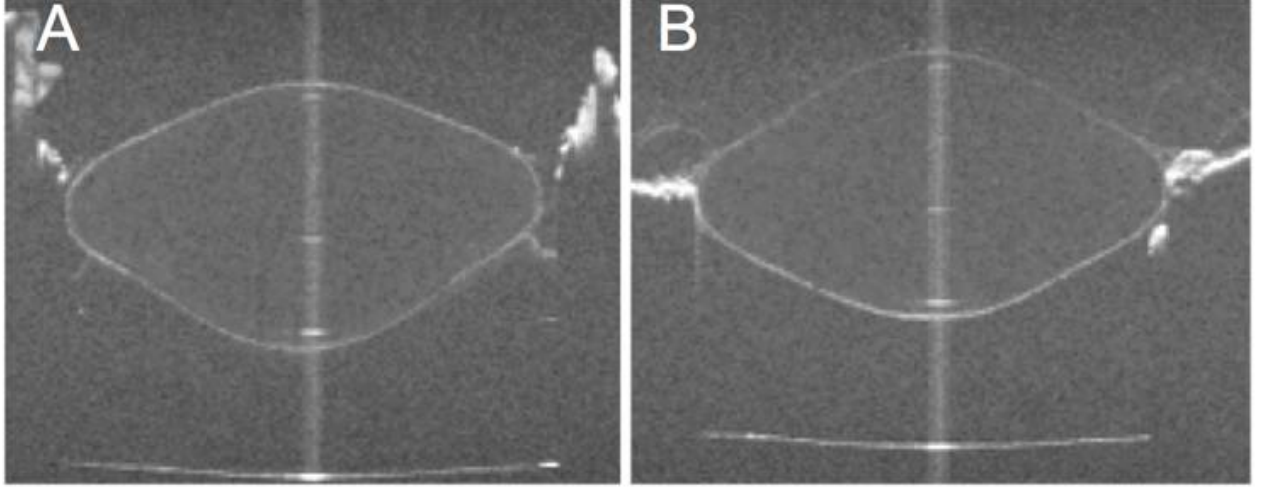


Figure 4.1: OCT images of the crystalline lens with the anterior surface facing the OCT beam (A) and with the posterior surface facing the OCT beam (B). The distorted surfaces contain the information of the optical path of the rays passing through the lens. This data together with power measurements are used in this study to reconstruct the gradient index of refraction of the lens. Images are for a 5.5 year old cynomolgus lens fully accommodated.

4.2.7 OCT Image Analysis

The shape of the lens surfaces was obtained, for each stretching position, from the OCT images. The distortion of the first surface (anterior in anterior-up images and posterior in posterior-up images) is only due to the presence of DMEM therefore its correct shape can be extracted by dividing its height by the DMEM group refractive index. The segmented edges were fitted to conic sections. The optical path difference and lens thickness were calculated as described in chapter xx. For all computations, the central 6-mm area of the lens was evaluated.

4.2.8 GRIN Reconstruction

The OCT-based GRIN reconstruction technique has been described in detail in chapter 2 and prior publications [47][48]. The height of the distorted surfaces in the OCT image

is a recording of the optical path traveled by the light. The search algorithm used to reconstruct the GRIN was the same as in the previous chapter, but to model the GRIN, we used a 3-variable power equation model (due to the 2D nature of the images), described in detail in chapter xx and a previous study[48]. The model is described in polar coordinates with the origin on the optical axis at 0.41 times the central thickness of the lens, and is expressed by the following equation:

$$n(\rho, \theta) = n_N - (n_N - n_s) \cdot \left(\frac{\rho}{\rho_{(s,\theta)}} \right)^p \quad (4.1)$$

The parameters of this model are the nucleus (n_N) and surface (n_s) refractive indices (which are assumed constant across accommodation levels) and a power coefficient (p) (which varied across accommodation levels) to model the decay from nucleus to surface in all directions. The value $\rho_{(s,\theta)}$ is the distance from the center to the conic surface in the direction θ . A ray trace using the Sharma algorithm[223] was programmed in MatLab (Mathworks, Natick, MA) to compute the optical path of the rays, and the results were compared with the experimental data. Input experimental data (optical path differences from OCT and lens power) collected for all stretching levels were pooled, and the optimization was performed simultaneously. The search algorithm was a mixture of a global (genetic algorithm of 20 generations of 200 solutions each) and a local search (simplex)[47]. The reconstruction of the GRIN was achieved in 30 minutes, for a full series of lens data in 11 different stretching positions. The mean value of 10 repetitions of the search algorithm was taken as solution. Since the experimental data are obtained with the OCT light source, the measured GRIN values correspond with the group refractive indices at the OCT central wavelength (825nm) and can be transformed to phase refractive indices at any wavelength[254]. The transformation is explained in detail in a previous publication[254]. In brief, the dispersion data from the literature[13] were used to transform the group refractive index to phase refractive index at 850 nm. The phase refractive index at any other wavelength can be calculated with the constant scaling formula derived by Atchison and Smith[13].

4.2.9 Estimated Thickness of Lens Nucleus and Cortex

To study the contribution of the nucleus and the cortex to the change of lens thickness with accommodation, the nucleus was defined as the central lens region of the reconstructed GRIN for which the refractive index varies within 1% of the peak refractive index, following the definition proposed in previous studies[107].

4.2.10 Estimated Power and Spherical Aberration

Virtual ray tracing was performed on the reconstructed lens (lens shape and GRIN), for 6-mm pupil diameter using the ray trace program written in MatLab (101 rays, ray spacing $30\ \mu\text{m}$). Since the back vertex power was experimentally measured using the OCT beam, the group refractive index was used in the calculations. The back focal length was calculated as the distance from the posterior vertex of the lens to the position where the root mean squares of the ray heights reached a minimum. The estimated back focal length, calculated with the reconstructed GRIN, was used to estimate the equivalent refractive index, defined as the index of a lens with the same geometry and power as the crystalline lens. The error in the approximation of using the back vertex power instead of the power in the calculation of the equivalent refractive index was estimated for the mean geometry of the cynomolgus lenses. The wavefront was computed after the lens (assuming the exit pupil plane at the position of the lens posterior surface vertex). The wave aberration was estimated with respect to a reference sphere centered at the paraxial focal point. Wave aberration was fitted by a 6th order Zernike polynomial expansion. The 4th order spherical aberration Zernike coefficient was evaluated. Calculations were performed for both a homogeneous refractive index (equivalent index) and for the estimated GRIN profile. By definition, the lens power for the homogeneous equivalent refractive index and for the GRIN is the same. The spherical aberration, assuming a homogeneous refractive index or the GRIN, was evaluated in all accommodative steps. Also, the contributions of the lens surface geometry to lens power and to spherical aberration were calculated replacing the GRIN by a homogeneous refractive index equal to that of the surface, i.e. these were not calculated for each surface separately.

4.3 Results

4.3.1 Stretching Versus Refractive Change

Figure 4.2 shows the back vertex power decrease with stretching, in all lenses. On average across lenses, back vertex power changed from $56\pm 4\ \text{D}$ for 0 mm stretching to $29\pm 4\ \text{D}$ at 2.5 mm stretching. In what follows, data will be plotted as a function of the measured lens back vertex power (corresponding to each of the stretching levels).

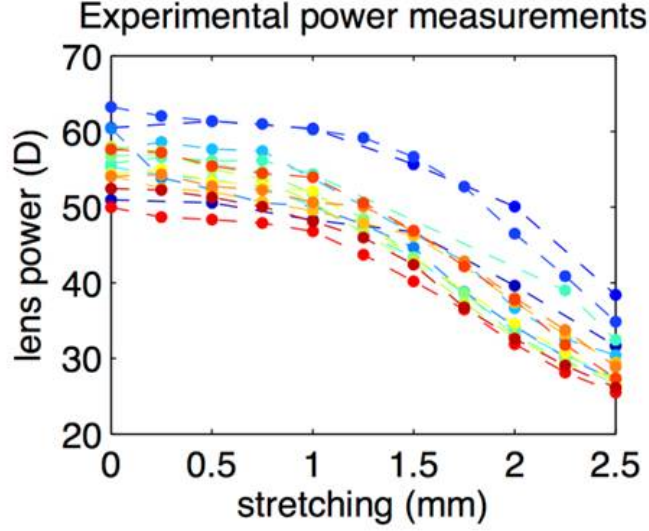


Figure 4.2: Back vertex power of all the studied lenses measured experimentally, as a function of stretching.

4.3.2 Changes of the Lens Geometry With Accommodation

The central 6 mm of the lens surfaces were well fitted to by conic sections[145]. The root mean squares of the residuals of the fittings were below $20 \mu\text{m}$ in all the surfaces. The anterior and posterior radii of curvature decreased with accommodation (Figure 4.3 A, B). Changes with accommodation were larger in the anterior surface than in the posterior surface. The average un-accommodated anterior and posterior lens radii of curvature were 8.7 ± 1.2 and 4.5 ± 0.4 mm, respectively. Anterior lens radius decreased at a rate of 0.19 ± 0.04 mm/D, and the posterior lens radius at a rate of 0.067 ± 0.012 mm/D. The conic constant of the anterior surface showed large variability (not correlated with age or postmortem time) but was negative for all the lenses (with values from -14 to -2 in the un-accommodated state) and shifted towards less negative values with accommodation (Figure 3C). The conic constant of the posterior surface remained near 0, and was rather constant in the entire accommodative range (Figure 4.3 D). Thickness increased linearly with accommodation at a rate of 0.036 ± 0.004 mm/D (Figure 4.3 E). All parameters varied almost linearly when plotted as a function of the back focal length. Except for anterior conic constant, all parameters varied similarly in all lenses.

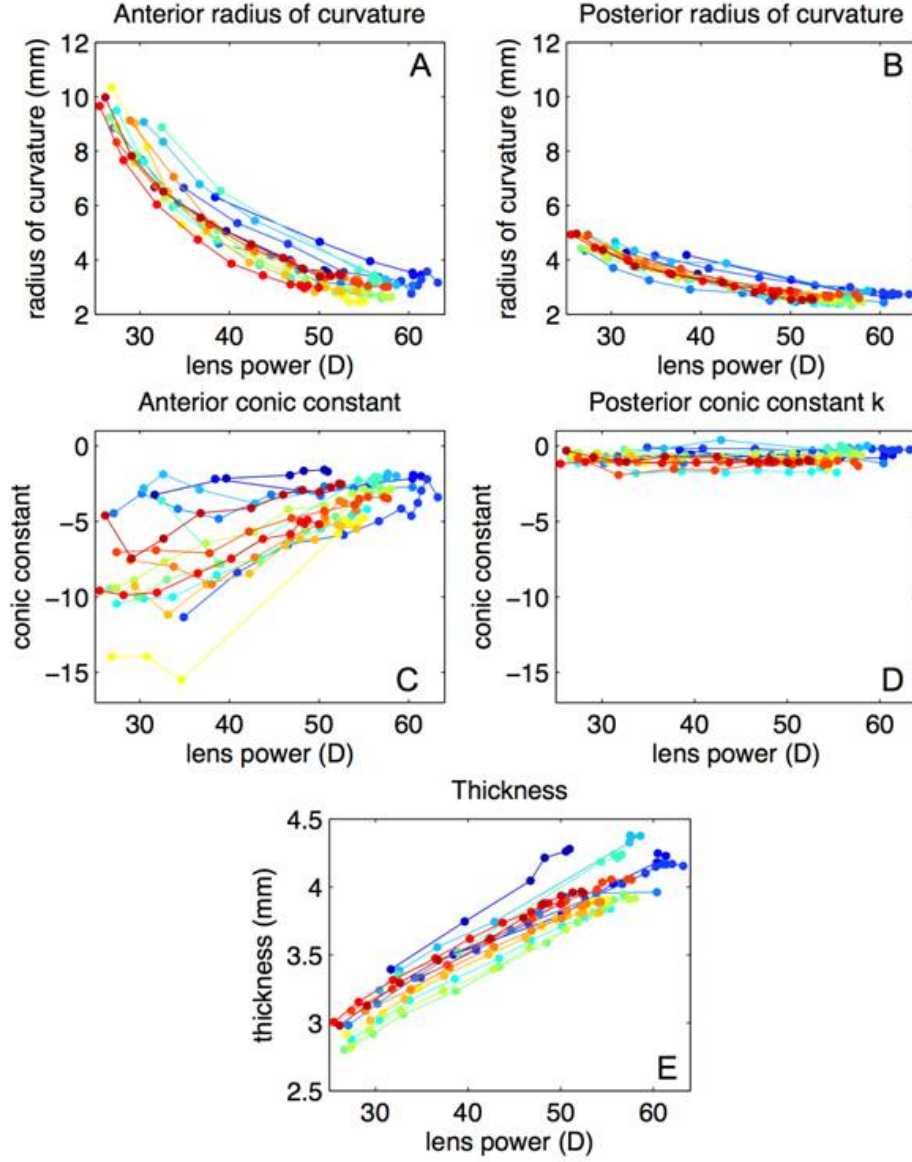


Figure 4.3: Lens geometry as a function of accommodation (stretching-induced changes in lens back vertex power). Using anterior-up and posterior-up images, the surface of the lens can be measured without optical distortion. A. Anterior lens radius of curvature. B. Posterior lens radius of curvature. C. Anterior lens conic constant. D. Posterior lens conic constant. E. Lens thickness.

4.3.3 Changes of the Lens GRIN With Accommodation

Figure 4.4 shows the GRIN parameters from the reconstructed GRIN distributions in all lenses: surface and nucleus refractive indices (figure 4.4A) and the power exponent

(figure 4.4B). The surface and nucleus refractive indices were 1.375 ± 0.003 and 1.429 ± 0.003 , respectively, on average across lenses. The power exponent ranged from 2.1 to 9.1 across lenses. On average, the power exponent remained rather constant with accommodation showing a slight but non-significant ($p > 0.5$) increase with accommodation (average across lenses $p = 4.45 + 0.001 * \text{back vertex power (D)}$). The inter subject variability observed in the exponent was not correlated with age or post mortem time. The mean average values of surface and nucleus phase refractive indices at 633 nm were 1.365 and 1.419 respectively.

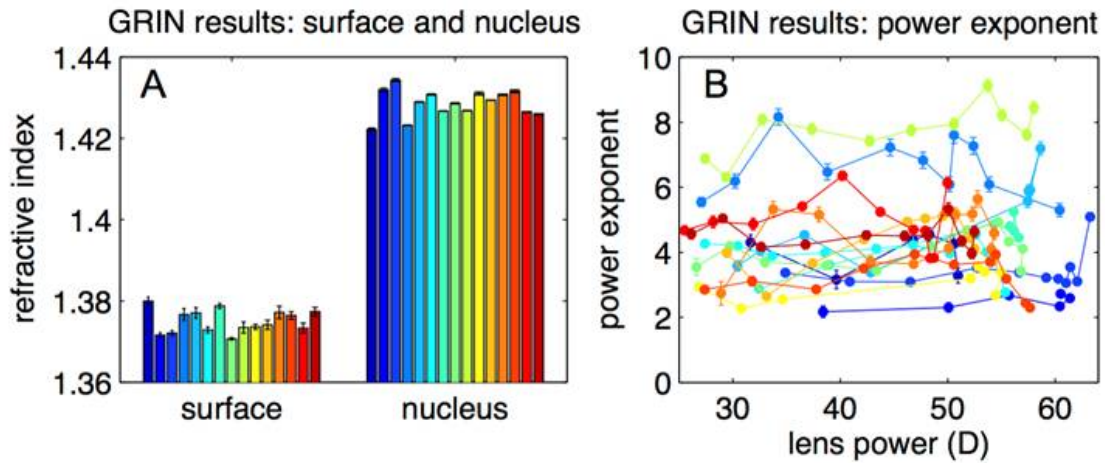


Figure 4.4: Results of the GRIN reconstruction. A. Surface and nucleus refractive index. B. Power exponent as a function of lens back vertex power.

Figure 4.5A shows the lens average refractive index calculated directly from the OCT images. This parameter remains constant with accommodation. The equivalent refractive index (figure fig4-5B) is also constant with accommodation, indicating that a similar power change would be achieved if the GRIN was replaced by the equivalent refractive index. The equivalent refractive index was calculated using the back focal length. Simulations in ZEMAX show that differences in the estimated equivalent refractive index using focal length (measured from the principal plane) or back focal length differed less than 1% in both accommodated and un-accommodated state.

4.3.4 Contributions of Nucleus and Cortex to Lens Thickness

Figure 4.6A shows the thickness of the crystalline lens, nucleus and cortex (as defined in the method section). Due to the constancy of the power exponent through accommodation (figure 4.4B), the nucleus thickness accounted for a rather constant fraction of lens thickness (average value 73%). The change in lens thickness (0.035 mm/D) is mostly

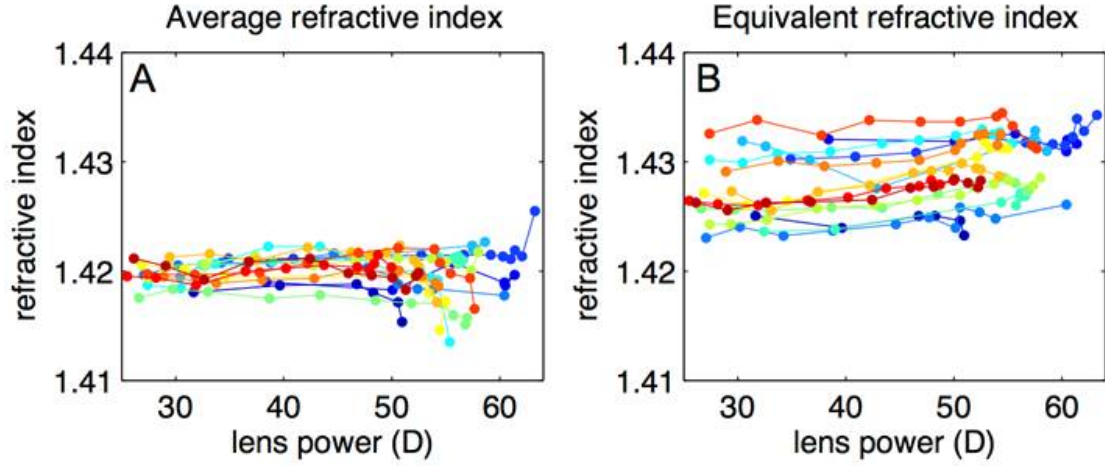


Figure 4.5: A. Average refractive index estimated from the distortion present in the cuvette in the OCT images as a function of back vertex power. B. Equivalent refractive index calculated from the back focal length for each accommodative state as a function of the back vertex power. The first index relates with the GRIN profile in optical axis and the second with the focal power of the crystalline lens (surface and GRIN).

due to an increase in the nucleus thickness (0.024 mm/D) while the contribution to the change of thickness of the cortex is moderate (0.005 and 0.007 mm/D for anterior and posterior cortex respectively) with accommodation (Figure 4.6B). During accommodation the lens increases its thickness by a factor of 1.42. Relatively, the larger thickness change occurs at the cortex. The nucleus, anterior cortex and posterior cortex sections increase by a factor of 1.38, 1.53 and 1.52 respectively).

4.3.5 Contributions to Lens Power Changes With Accommodation

The lens surface shape contributed on average 35-40% to the back vertex power of the crystalline lens in all accommodative states. The fact that this contribution is constant with accommodation indicates that the contribution of the external geometry to the amplitude of accommodation is similar. Surfaces contributed 8.4 ± 1.2 D to the amplitude of accommodation of the lens (27 ± 4 D), while GRIN was found responsible for almost 20 D of accommodation.

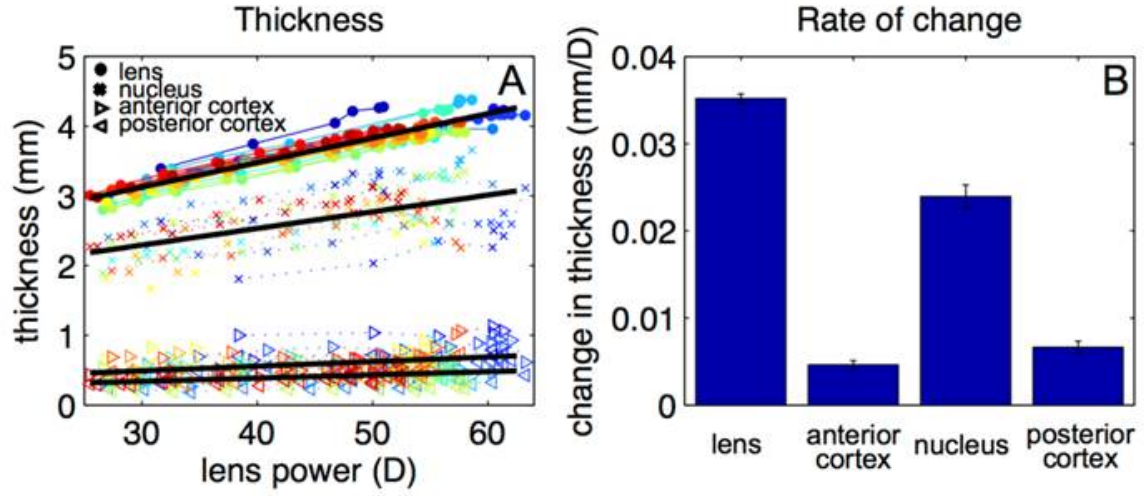


Figure 4.6: A. Thickness of the crystalline lens, nucleus and the anterior and posterior cortex as a function of back vertex power for all lenses. Black lines represent linear fits to the data: total thickness (mm) = $2.070 + 0.035 \times \text{lens power (D)}$, nucleus thickness (mm) = $1.576 + 0.0239 \times \text{lens power (D)}$, anterior cortex thickness = $0.202 + 0.005 \times \text{lens power (D)}$ and posterior cortex thickness = $0.291 + 0.007 \times \text{lens power (D)}$. B. Rate of change of the thickness of the lens, anterior cortex, nucleus and posterior cortex of the lens.

4.3.6 Spherical Aberration

The estimated spherical aberration was negative throughout the accommodative range, with values of -2.3 ± 0.7 m for the un-accommodated state and -5.6 ± 1.5 m in the fully accommodated state. Figure 4.7 shows a comparison between the spherical aberration of the lens with the estimated GRIN and with the equivalent refractive index, assuming the measured lens geometry in both cases. The inter subject variability was not correlated with age or post mortem time. The spherical aberration magnitude assuming the equivalent index is lower compared to that found in the GRIN lens, with a larger difference in the fully accommodated state (53% on average) than in the un-accommodated state (29% on average). The spherical aberration increased in absolute value with accommodation at a rate of 0.124 m/D on average (for the lens with the calculated GRIN) and at a rate of 0.070 m/D on average (for the lens with the equivalent refractive index).

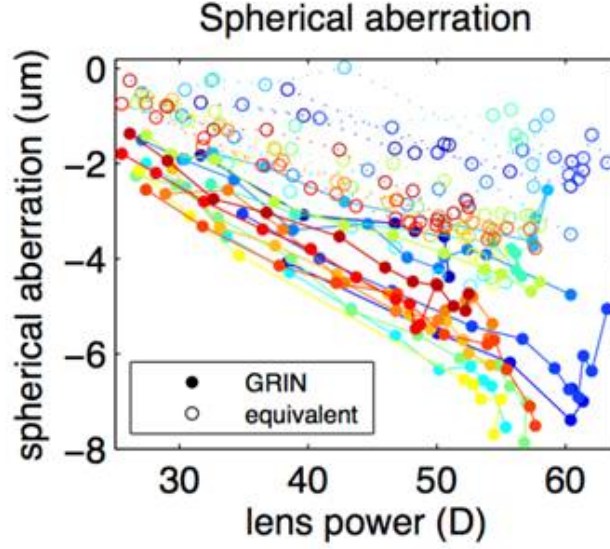


Figure 4.7: Spherical aberration as a function of lens back vertex power through accommodation, for the lens with the reconstructed GRIN (solid circles) and with the equivalent refractive index (empty circles).

4.3.7 Spherical Aberration Changes With Accommodation. Contribution of the Surfaces.

The spherical aberration of a homogeneous lens with an index equal to the surface refractive index (figure 4.8A) varied between -0.25 ± 0.06 and -1.2 ± 0.3 m. The relative contribution of the surfaces spherical aberration to the total spherical aberration of the lens was 19% on average (figure 4.8B). This contribution increased slightly from the un-accommodated state ($12 \pm 4\%$) to the fully accommodated crystalline lens ($21 \pm 6\%$). Figure 4.8 C shows the change of the spherical aberration with the calculated GRIN, $3.8 \pm 1.2 \mu\text{m}$ as an average, and with homogeneous lens with index equal to the surface refractive index, 0.9 ± 0.2 m. Figure 4.8 D shows the contribution of GRIN and surfaces to the change of spherical aberration. On average, the gradient refractive index was found to be responsible for $73 \pm 9\%$ of the spherical aberration change of the lens.

4.4 Discussion

During accommodation, the primate crystalline lens changes its shape and the gradient index of refraction redistributes producing well known changes in the optical properties of the crystalline lens: an increase of power and a shift of the spherical aberration

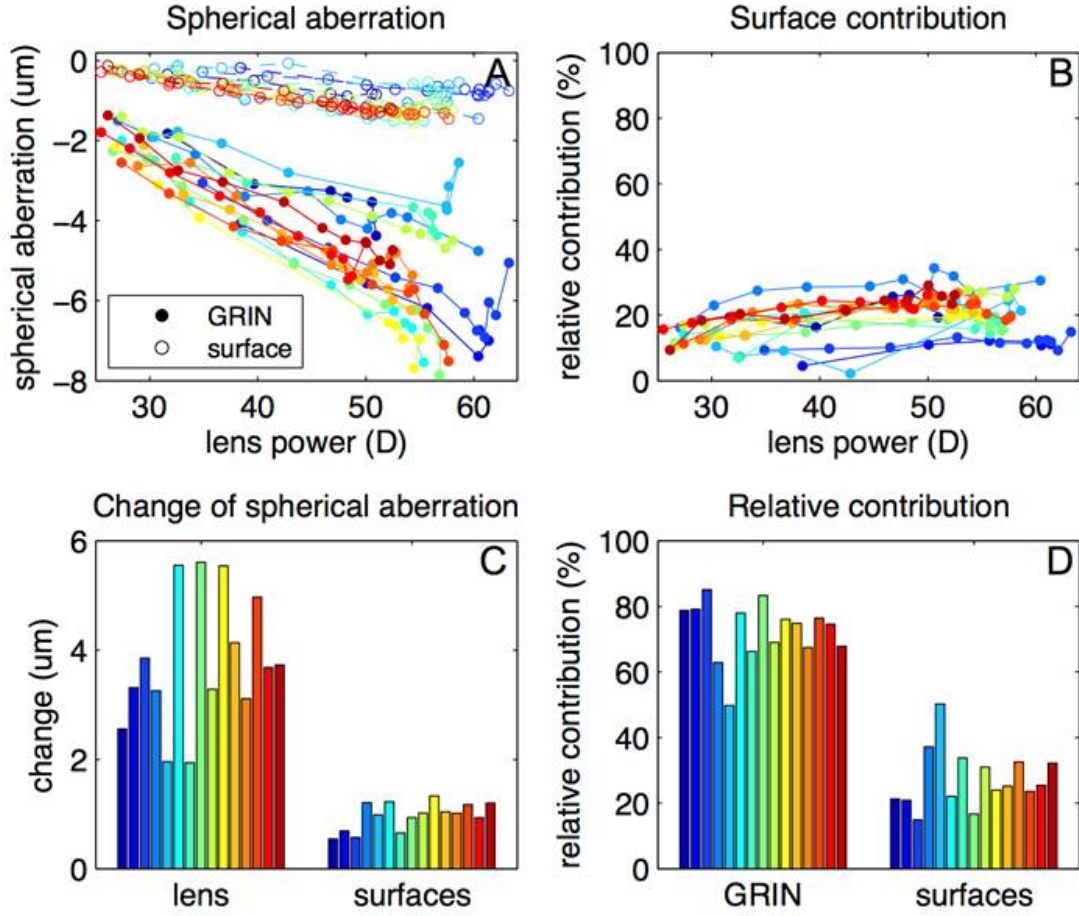


Figure 4.8: A. Spherical aberration of the lens with GRIN compared with that of a homogeneous lens with index equal to the surface refractive index as a function of back vertex power. B. Contribution of the surface to the magnitude of Spherical aberration as a function of back vertex power. C. Change in absolute value of spherical aberration of the lens compared with that of the surfaces from unaccommodated to fully accommodated. D. Relative contribution of the surface and GRIN to the change of spherical aberration with accommodation.

towards more negative values. We have presented experimental measurements of the changes of both the geometry of the crystalline lens, and of the GRIN profile of non-human primate lenses (cynomolgous monkeys) with accommodation. Accommodation was simulated in vitro using a stretching device, a paradigm that has been previously demonstrated to mimic accurately the change in lens shape and power occurring in vivo [72][18][261][32]. These measurements have allowed us to quantify the role of the GRIN in accommodation, and to estimate the relative contribution of surface and GRIN to lens

power, and most notably, to the spherical aberration, as a function of accommodation. All the results presented in this study were calculated using the group refractive index at 825 nm in both the crystalline lens and the preservation medium (DMEM, 1.345). To check the validity of the results in visible light, the ray tracing was repeated with the surface and nucleus refractive index converted to phase refractive index at 555 nm and assuming an aqueous and vitreous refractive index of 1.336. The differences between group refractive index and phase refractive index are around 1% (0.007 as a mean), producing average differences in power and spherical aberration below 3% (1.1 D and below 0.1 m respectively).

In agreement with previously reported in vivo and in vitro data in humans [58][70] and in rhesus monkey lenses[114][217], we found a decrease in the lens surfaces radii of curvature with accommodation, larger for the anterior surface than for the posterior surface. The slight smaller change per diopter of accommodation in cynomolgus monkeys may be due to inter-species differences, age differences and the high asphericity of anterior surface.

We found that the anterior surface conic constant of the un-accommodated cynomolgus lens was consistently negative (although it varied significantly across individuals, ranging from -14 to -2) and decreased in most of the lenses toward a more spherical shape with accommodation. The posterior conic constant was also negative but close to zero and remained constant with accommodation. These data differ from reports in human lenses in vivo [58] measured with Scheimpflug imaging, where both anterior and posterior surfaces had negative asphericity with similar mean values, although the intersubject variability was very large. Also, the change with accommodation differs, as in humans the anterior surface seem to become more curved with the peripheral areas of the lens remaining relatively flatter, and therefore the conic constant changes towards more negative values. However, the results in human lenses in vivo[58] show small changes in the anterior conic constant with accommodation, and no accommodation-related changes in the posterior surface conic constant, due to limitations in the technique. Studies in in vitro human crystalline lenses[145] reported a positive conic constant in the anterior surface and scattered values around 0 for the posterior. The values of asphericity reported here for the fully accommodated state agrees well with those reported by a previous study in isolated cynomolgus monkey lenses[33] where a wide range (from -6 to +4) was reported. The high value of monkey's lens anterior surfaces asphericity in the un-accommodated state, never reported in humans, may be due to the differences in size between the species, to the young age of the monkey lenses compared to the humans from prior literature, and to the larger accommodative range of monkey lenses (up to 30 D) in comparison with the human lens (no more than 10 D).

We found that the lens thickness increased with accommodation at a rate of 0.036 ± 0.004 mm/D. This thickness variation is comparable to that reported in vivo in rhesus monkeys under stimulated accommodation, around 0.041 mm/D[114] and 0.063 mm/D[260], and in vivo in humans, 0.045 mm/D[58] and 0.064 mm/D[212]

In agreement with previous findings we have found that a rather constant value of the parameters defining the GRIN best fitted the experimental data[70][91] and that the equivalent refractive index did not change with accommodation [70][142][91].

Despite the simplicity of the GRIN model, the estimated parameters allow reproducing the experimental input data with great accuracy (mean RMS_i40 m), for all accommodation levels, using our recently developed optimization method based on OCT imaging[47]. Previous studies with simpler GRIN models suggested that no change in the parameters of the model was needed to account for a change of power of the crystalline lens with accommodation [70][91]. We have found a slight trend for an increase in the power exponent parameter of our GRIN model (expansion of the central plateau in the GRIN distribution), although this was not statistically significant. This was in contrast with a report using MRI[107], which suggested a decrease of this parameter (only significant in the meridional direction), and a 50% contribution of the lens thickness nucleus to the change of lens thickness with accommodation. Studies based on Scheimpflug imaging revealed a much higher contribution of the lens nucleus (90%) in humans[35][59][90] and rhesus monkey[113]. While differences arising from the definition of the lens nucleus and the data analysis are expected, our study on cynomolgus monkeys also supports a large contribution of the lens nucleus in thickness changes (69%, following the definition proposed by Kasthurirangan et al.[107]). This conclusion is highly dependent of the adopted definition for the nucleus. While the current definition does not necessarily best describe the physiological area of the lens nucleus, it allows comparison with prior literature using this definition[107]. Nevertheless the estimated nucleus thickness relative to the total thickness (73%) with that definition is only slightly larger than that obtained from direct imaging of the lens (65%, from Scheimpflug imaging in young human subjects in vivo[59]; 57% from OCT imaging as an average in isolated human crystalline lenses of different ages[49]).

We found the spherical aberration of the un-accommodated cynomolgous lens to be negative ($-2.3 \mu\text{m}$, 6-mm pupil diameter), as reported in young human lenses ($-0.16 \mu\text{m}$, 6-mm pupil in vivo[22]) and rhesus monkey ($-1 \mu\text{m}$, 7-mm in vitro[213] and $-0.5 \mu\text{m}$, 8-mm in vivo[261]). Previous studies have reported a compensatory role of the GRIN in different species such as fish[100], rat[40], and porcine lenses[47][255][269]. With those lens geometries, and with a homogeneous index of refraction, the spherical aberration of the lens was positive, and the presence of GRIN shifted the lens towards

less positive values (or close to zero, such as in fish) or towards negative values, which tended to compensate the positive spherical aberration of the cornea. In cynomolgus lenses we have also found a compensatory role of GRIN. However, in this case, the GRIN did not reverse the sign of the spherical aberration, already negative with an equivalent refractive index, but rather doubled its magnitude. The spherical aberration value of the lens with an equivalent refractive index was on average 41% of that of the corresponding GRIN lens. Also, the presence of GRIN emphasized the change of spherical aberration with accommodation. The change in spherical aberration with accommodation estimated with the equivalent refractive index was on average 29% of the change for the corresponding GRIN lens. While there is a large intersubject variability (which we did not find to be correlated with age or post mortem time), we have found that the contribution of GRIN in the spherical aberration is slightly larger in the fully accommodated state.

As in rhesus monkeys [213][261], we found a shift of the spherical aberration during accommodation towards more negative values. We found larger differences (unaccommodated - fully accommodative state) in spherical aberration in cynomolgous ($3.3 \mu\text{m}$ for a 6-mm pupil) than in rhesus monkeys (around $2 \mu\text{m}$, 8-mm, in vivo[261]; 1.7 m , 7-mm, in vitro[213]). However, as the accommodation amplitude is larger in cynomolgus (20-30 D) than in rhesus (around 17 D), the shift of spherical aberration per diopter of accommodation appears relatively similar across species (-0.124 m/D in the current study in cynomolgus monkeys; around -0.11 m/D in vitro[213] and -0.19 and -0.24 m/D in [261] in rhesus). These values are higher than those reported in humans for 6-mm (-0.013 m/D5 , -0.083 m/D3). The calculated spherical aberration value may be affected by errors in the calculation of the geometry of the surfaces of the crystalline lens, in the reconstruction of GRIN and in the accuracy of the GRIN model itself.

As in previous studies[70][142][20][237][104], the contributions of the GRIN to the power and spherical aberration were computed by comparing the lens optical properties with those produced by refraction in the surfaces only. While previous publications assumed a fixed surface refractive index to estimate the contribution of the surfaces to the change of power of the lens, in this study the surface index was obtained directly for each lens from the GRIN reconstructed from the experimental data. The contribution of the GRIN to the power and accommodative amplitude of the lens reported is in agreement with previous studies in vitro in baboon and cynomolgus monkeys[142] and in humans[32]. For spherical aberration, we found that the surface contribution was around 20%, and that the contribution was larger in the fully accommodated state. These results are indicative of a large contribution (almost 80%) of the GRIN to the value of lens spherical aberration in the entire accommodative range. Also the redistribution of the GRIN was found to be responsible for more than 70% of the change of spherical aberration through

accommodation.

For these calculations, we defined the spherical aberration contribution of the GRIN as the difference between the spherical aberration of the reconstructed GRIN lens and the spherical aberration of a lens with a homogeneous refractive index equal to that of the surface, in a similar way as previously evaluated for the surface / GRIN contributions to the power of the lens[70][142][32][104]. There are other possible definitions or values of the index of the homogeneous lens that could be used, which may produce different numerical results, but we expect that the general finding regarding the importance of the GRIN will not change. For instance, our analysis shows that the change in spherical aberration of the lens with homogeneous index equal to the equivalent index (instead of the surface index) is still only 56% of the change found with the GRIN lens.

Theoretical analyses[172][138] have shown that model eyes that simulate the change of spherical aberration with accommodation can be designed using a lens model with a shape, and a homogeneous equivalent refractive index ($n=1.429$), based on the measurements of Dubbelman et al.⁸ However, this lens model does not provide accurate predictions of the actual value of lens spherical aberration, most likely due to uncertainties in the lens surface asphericity values[172]. Our experimental studies on monkey lenses are in disagreement with the conclusion of Lopez-Gil et al.[138] regarding the contribution of the GRIN to the change of spherical aberration with accommodation. We find that the equivalent homogeneous lens predicts the general trend of the change in spherical aberration with accommodation, but it does not reproduce the actual spherical aberration value and its change with accommodation, obtained with experimental measurements of the crystalline lens shape and GRIN.

5

Contribution of shape and gradient refractive index to the spherical aberration of isolated human lenses

This chapter is based on the paper by Birkenfeld et al. "Contribution of shape and gradient refractive index to the spherical aberration of isolated human lenses" (Invest Ophthalmol Vis Sci. 2014; 55:2599-2607) and the paper by Sun et al. "OCT 3-D surface topography of isolated human crystalline lenses". The coauthors of the first paper were Alberto de Castro and Susana Marcos. The author of the thesis was first coauthor in the second paper. Other coauthors were Alberto de Castro, Sergio Ortiz, and Susana Marcos. For the first paper, the author of this thesis (i) modified the experimental setup (OCT), (ii) did the sample preparation (iii) performed the experiments (in collaboration with Alberto de Castro) (iv) performed data processing (in collaboration with Alberto de Castro), (v) analyzed the results (vi) prepared the manuscript (in collaboration with Alberto de Castro and Susana Marcos). For the second paper, the author of this thesis (i) modified the experimental setup (OCT), (ii) did the sample preparation (iii) performed the experiments (in collaboration with Alberto de Castro), (iv) performed data processing (in collaboration with Mengchan Sun) (v) analyzed the results (in collaboration with Mengchan Sun) (vi) prepared the manuscript (in collaboration with Mengchan

Sun and Susana Marcos). The studies allowed us to investigate the influence of shape and GRIN on the spherical aberrations of human lenses. As a result we concluded that both, geometrical changes and GRIN contributes to the age-dependent shift of negative spherical aberration. It also allowed us to study 3D surface topography on human crystalline lenses in vitro, and to describe the surfaces by means of Zernike polynomials, and evaluate the correlations between anterior and posterior surface topography.

5.1 Introduction

In the last chapters we reconstructed the GRIN of porcine and monkey lenses and investigated its influence on optical aberration in different states of accommodation. In this chapter we use 3-D OCT and LRT measurements of 35 isolated human donor lenses (between the age of 19 and 71 years) to fully characterize the lens surface geometry, the GRIN distribution, optical properties, and, in addition, analyze their age-dependencies. With age, the human lens undergoes various physical, biometrical and optical changes[73][36][89][115][57][140][70][12], such as thickness increase, surface steepening and mass and volume linear growth. The changes in optical property result from changes of the lens' shape and its refractive index. In particular, the overall spherical aberration of the eye shifts toward positive values[72][157][6]. This change in lens spherical aberration leads to age-related loss of the corneal/lens spherical aberration compensation and decrease of the optical performance [8][81][234][22].

In contrast to the cornea, there is little information on lens topography and possible compensations within the crystalline lens itself, i.e. between anterior and posterior surface or between its surfaces and its GRIN, although there are also numerous reports of the compensation of the corneal astigmatism by the astigmatism of the crystalline lens using corneal and refractive parameters [110][119][182][136]. As an additional part of this work, we characterized the lens topography, using lens surface elevation maps. As descriptive parameters we included individual surface Zernike coefficients and the root mean square (RMS) of combination of some terms or Zernike orders.

5.2 Methods

5.2.1 Human lens samples and preparation

Human donor eyes were obtained from Transplant Service Foundation (TSF) Eye Bank. The handling and experimental protocols had been previously approved by the Institu-

tional Review Boards of TSF and CSIC. Experiments were performed on 35 eyes from 28 different donor eyes, 1 to 3 days post-mortem. The donor age ranged between 19 and 71 years. All eyes were shipped in sealed vials at 4°C, wrapped in gauze soaked in preservation medium (DMEM/F-12, HEPES). The lens then was carefully extracted from the eye and immersed in preservation medium (DMEM/F-12, HEPES, no phenol red, GIBCO)[19] at room temperature (see figure 5.1. During the measurements, the lens was placed on ring in a cuvette filled with DMEM solution. The whole measurement usually took between 1 to 2 hours. Swollen or damaged lenses would be identified from the OCT images and were excluded from the study.

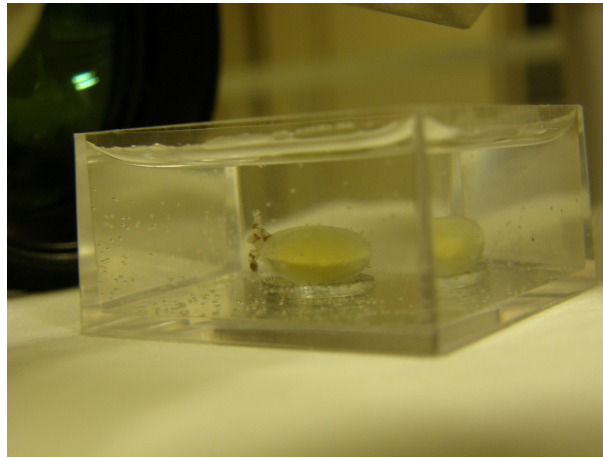


Figure 5.1: Extracted human donor lens in a DMEM filled cuvette

5.2.2 Optical Coherence Tomography Imaging

All lenses were imaged in 3-D using a custom developed high resolution spectral domain OCT system describes. One 3-D image was composed by 1668 A-Scans, in 70 B-Scans on a 12x12 mm lateral area, acquired in 4.5 s. The axial resolution was calculated to 6.9 μm in tissue. The lens axis was aligned with the OCT scan axis such that a specular reflection was seen from the surfaces of the lens. To center and align the lens, real-time display horizontal and vertical A-scans were used. 3-D Images were acquired in two different focal planes of the lens, to allow visualization of the anterior and posterior surface of the lens, and the posterior surface and the holding cuvette in two sequential acquisitions, which were merged using a custom-developed merging algorithm[47]. Images were obtained with the anterior surface up and with the posterior surface up, after carefully flipping the lens in the holder.

5.2.3 Laser Ray Traing

The paraxial and non-paraxial focal length of each lens were measured using a 2-mirror galvanometric scanning laser ray tracing system introduced in chapter 2. It scanned rings of light of different diameters (2 and 4 mm) onto the crystalline lens and the CMOS camera captured a series of through-focus images around the focal plane (full range: 35 mm; step size: 0.5 mm). The lens was placed in a cuvette (posterior up) and was aligned with the LRT system, such that the lens and principal ray was collinear to the center of the CMOS in the entire focus range. Measurements of the focal length at two different pupil diameters allowed an approximate experimental estimate of the 4th order Zernike spherical aberration, using the following equation [247]

$$Z_4^0 = \frac{M r^2}{12 \sqrt{5}} \quad (5.1)$$

where r is the paraxial pupil radius, and $M = M(\text{non-paraxial power}) - M(\text{paraxial power})$. The group refractive index of the solution was taken as $n = 1.345$ at 825nm [34].

5.2.4 Image Processing

All OCT images were processed to obtain the shape of the lens surfaces and the surface of the cuvette, using algorithms described in detail in chapter 2. Images at different focal planes were merged into one complete 3-D image of the lens. Images were then corrected from fan distortion (Ortiz2009) and optical distortion induced by the presence of the preservation medium. The full geometry of the lens was obtained from analysis of the anterior surface measured in each condition. All surfaces (lens and cuvette) were fitted by Zernike polynomials (up to 7th order) within a 6-mm pupil. For the purposes of this study, only symmetric Zernike polynomials and astigmatism were used[145]. Additionally, the radii of curvature and conic constant of the surfaces were estimated for 18 meridians, using the following equation

$$z = z_0 - \frac{(x - x_0)^2}{r + \sqrt{r^2 - k(x - x_0)^2}} \quad (5.2)$$

where z is the height, x the radial position along the meridian, r the apical radius of curvature at the vertex, and k the conic constant. The lens thickness was calculated from the distortion of the image of the cuvette surface[254].

5.2.5 GRIN reconstruction

The 3-D GRIN was reconstructed from lens geometry, optical path difference, and lens focal length using the global search algorithm described in chapter 2. The GRIN was estimated for a 6-mm pupil. The algorithm was run five times for each data set, and was applied in 18 cross-sectional meridians (0-170, in steps of 10°). The reconstruction was done for a refractive index at 849 nm and was converted to a phase refractive index at 633 nm (Uhlhorn2008), to allow comparison with prior literature.

5.2.6 GRIN model and Search Algorithm

The GRIN model used was described in detail in chapter 2. In short, the refractive index in the lens is modeled so that it varies continuously from nucleus to the surface in both, axial and meridional directions. The center of the lens is assumed to be in the optical axis at a distance from the anterior surface vertex equal to 0.41 times the lens thickness[218]. The GRIN is described as a 4-variable model in polar coordinates:

$$n(\rho, \theta) = n_N - \Delta n \cdot \left(\frac{\rho}{\rho_s} \right)^{p(\theta)} \quad (5.3)$$

where n_N is the refractive index of the nucleus, Δn the difference between the refractive index of surface and nucleus, ρ_s is the distance between nucleus and surface, and $p(\theta)$ is the exponential decay in axial (p1) and meridional (p2) direction; p1 is constant across meridians, while p2 can vary to account for differences between meridians. The GRIN distribution which best fits the experimental data is searched through minimization of a merit function. A genetic optimization algorithm[95] was implemented as a global search algorithm to prevent the solution falling in a local minimum[47].

5.2.7 Equivalent refractive index and average refractive index

The homogeneous equivalent index was calculated by matching the refractive index producing the same focal length than the lens with the estimated GRIN and the same 3D geometry. Unless otherwise noted, the equivalent index was estimated for 6-mm pupil diameters. The average refractive index was estimated as is the mean value of the GRIN profile along the optical axis of the lens

5.2.8 Computational ray tracing analysis

The optical aberrations of individual lenses were investigated by means of a computational ray tracing analysis (performed at multiple meridians) using the experimental geometric data of the lens and the reconstructed GRIN. Additionally, aberrations were estimated for lenses of similar geometry and a homogeneous equivalent refractive index, allowing accounting for the relative contribution of lens geometry and GRIN to the optical properties of the lens. The custom-developed ray-tracing algorithm used is describes in detail in chapter 2. In short, the Stavroudis formula[241], and Sharma algorithm[223] are used to to trace rays through conical surfaces and GRIN, respectively. The calculated wave aberrations were fit by Zernike polynomial expansions. In particular, 4th order spherical aberration were calculated for GRIN lenses and their corresponding equivalent refractive index lenses.

5.2.9 Lens surface elevation

Lens surface elevation maps were obtained by subtraction of the best fitting spheres from the segmented surfaces and fitted to 6th order Zernike polynomials. All fittings were done for a 6-mm diameter optical zone. Descriptive parameters of the surface elevation maps include individual surface Zernike coefficients (up to 6th order), and the root mean square (RMS) of combination of some terms (RMS astigmatism, RMS trefoil, RMS tetrafoil, RMS spherical terms, RMS coma and RMS 4th order Astigmatism), or Zernike orders (from 2nd order to 6th order RMS). We studied the changes of those parameters with age as well as relationships between the anterior and the posterior surfaces. As the relative angle of the lens with respect to the cornea and the body is unknown, for convention, all lenses were aligned such that the astigmatic axis of the anterior lens was on the vertical meridian. The results of the lens surface elevation maps are based on a more recent study, therefore, one more lens is included.

5.3 Results

5.3.1 Lens surface elevation maps

Figure 5.2 shows all Zernike fittings to the anterior and posterior lens surface elevations (relative to the best fitting sphere) for all lenses of the study, sorted by age. The first and third columns show the astigmatic Zernike terms only, and the second and fourth columns represent astigmatism and high order aberrations (up to the 6th order terms) of

anterior and posterior surface respectively. Because all lenses have been aligned such that the anterior lens steepest meridian is vertical, the anterior maps show only astigmatic component. In many cases the astigmatic axis of the posterior lens surface is rotated with respect to the anterior lens surface. A certain degree of similarity occurs in the high order aberration maps of both eyes from the same donor eyes, marked with asterisks (lens #7&8, lens #11&12, lens #15&16, lens #17&18, lens #23&24, lens #27&28, lens #35&36). On average, for eye pairs from the same donor, the average correlation coefficients for the high order Zernikes between eyes are $r=0.78$ ($p=0.002$) for the anterior, and $r=0.53$ ($p=0.03$) for the posterior lens surface, while the average correlation coefficients when compared to eyes from other eyes are $r=0.10$ and $r=0.17$ for the anterior and posterior lens surface, respectively.

5.3.2 Age-dependence of crystalline lens shape

All lens surfaces were well fitted by conic sections for the central 6 mm of the surface. The root mean square of the residuals of the fittings was below 30 m. Lens thickness, radii of curvature and conic constants of the anterior and posterior surfaces were analyzed as a function of age. The thickness increased linearly with age at a rate of 0.0196 mm/year (figure 5.3). Anterior and posterior radii of curvature (Figure 5.4) were fit by a 3rd order and 2nd order polynomial, respectively. The anterior lens radii of curvature increased almost linearly up to age 60 (a linear fit between 19 and 60 years showed an increase of 0.119 mm/year), and tended to decrease beyond that age. Posterior lens radii of curvature tended to increase slightly with age.

Anterior and posterior lens conic constants (Figure 5.5) were fit by linear regressions. The anterior surface conic constant was negative for 26 of the 35 lenses (ranging from -16.2 to -0.08), and for all lenses under 47 years. It shifted towards positive values in older lenses, increasing linearly with age at a rate of 0.228/year ($p<0.001$). The posterior surface conic constant was -0.17 on average with values between 1.17 and -1.89, and showed a small increase of 0.0275/year ($p=0.002$).

Across meridians, the radius of curvature changed by 0.8 mm and 0.32 mm, and the conic constant changed by 2 and 0.12, for anterior and posterior surface, respectively (average values). We did not find any correlation of meridional changes with age. In addition, the magnitude of the changes in the anterior surface was not correlated with those of the posterior surface.

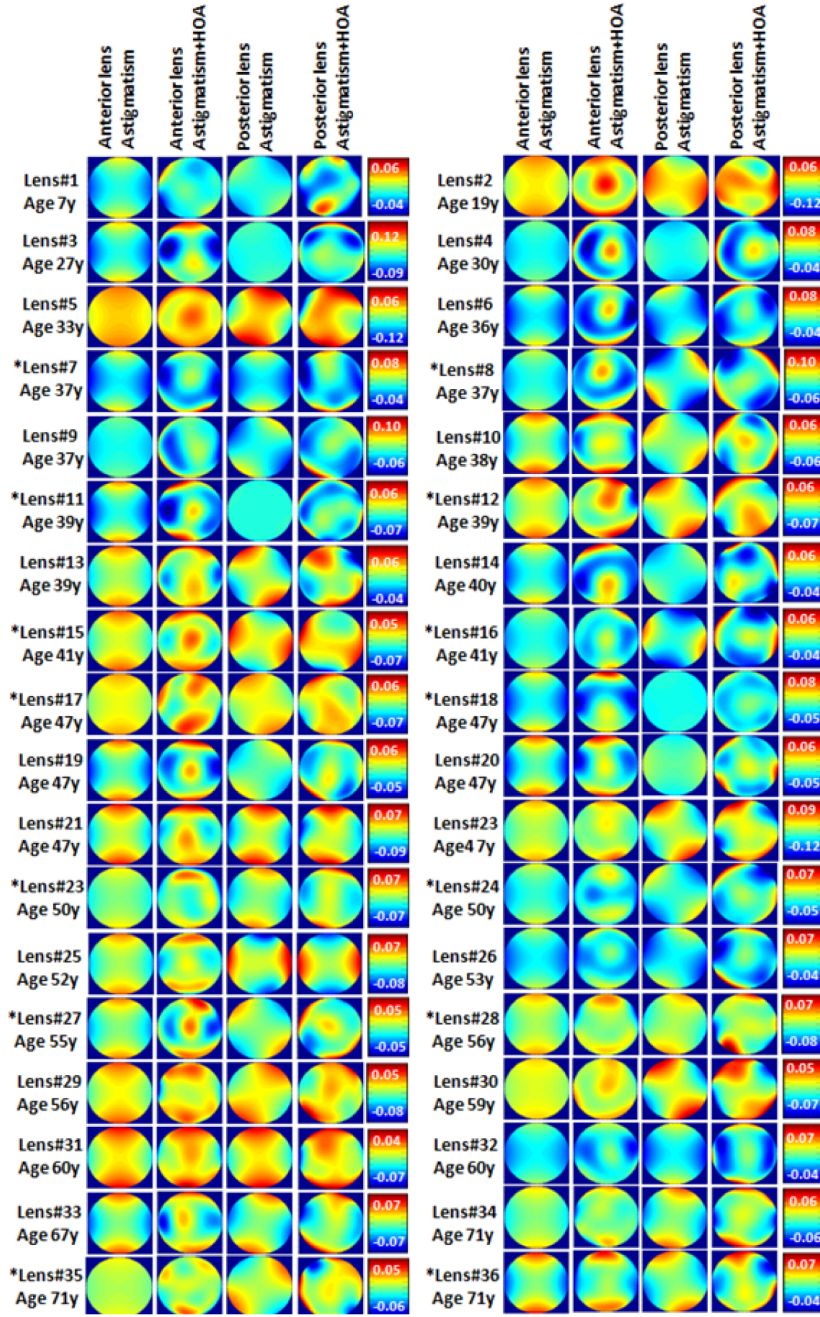


Figure 5.2: Lens surface elevation maps for all lenses, ordered by age. By convention, the maps are aligned so that the steepest meridian of the anterior lens surface lies in the vertical axis. Anterior and posterior images are shown as mirrored in the vertical axis. Asterisks indicate pairs of lenses from the same donor.

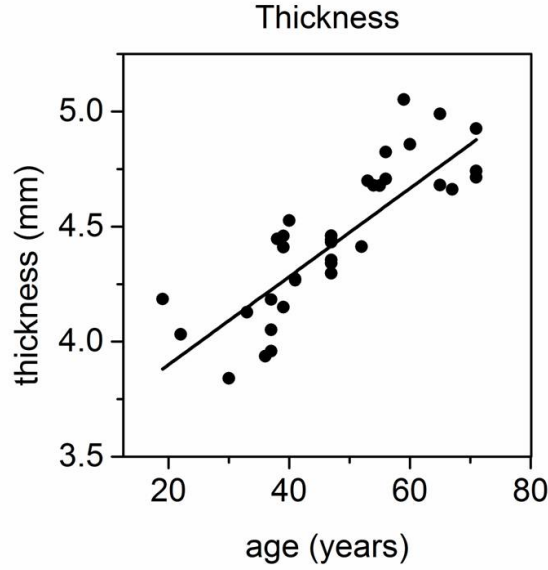


Figure 5.3: Central thickness of isolated human lenses as a function of age, and a linear fit to the data ($\text{Thickness}(\text{mm}) = 0.0196 \cdot \text{Age} + 3.5187$, $r = 0.84$, $p < 0.001$).

5.3.3 Distribution of the gradient refractive index

The reconstructed parameters for the GRIN are nucleus refractive index (n_N), surface refractive index (n_S), and the power exponents for axial (p1) and meridional (p2) decay. Surface refractive index varied from 1.3601 to 1.3896, and nucleus refractive index between 1.3985 and 1.425 (Figure 5.6). The surface refractive index increased by a statistically significant amount with age ($p=0.015$), at a rate of 0.0002/year. The nucleus refractive index tended to decrease with age, but its change was not statistically significant ($p=0.3$). Figure 5.7 shows the power exponents in axial (p1) and meridional (p2) directions. The power exponents in the meridional direction (p2) increased ($p=0.1$), indicating a flattening of the GRIN distribution with age. The power exponent in the axial direction (p1) remained almost constant with age. The power exponent p2 changed across meridians in some lenses, indicating a contribution of the GRIN to the astigmatism of these lenses. The mean change of p2 was 1.09 with a maximum change of 2 in one of the lenses (changes were on average 34%, and always below 58% of the mean value of the meridional exponent).

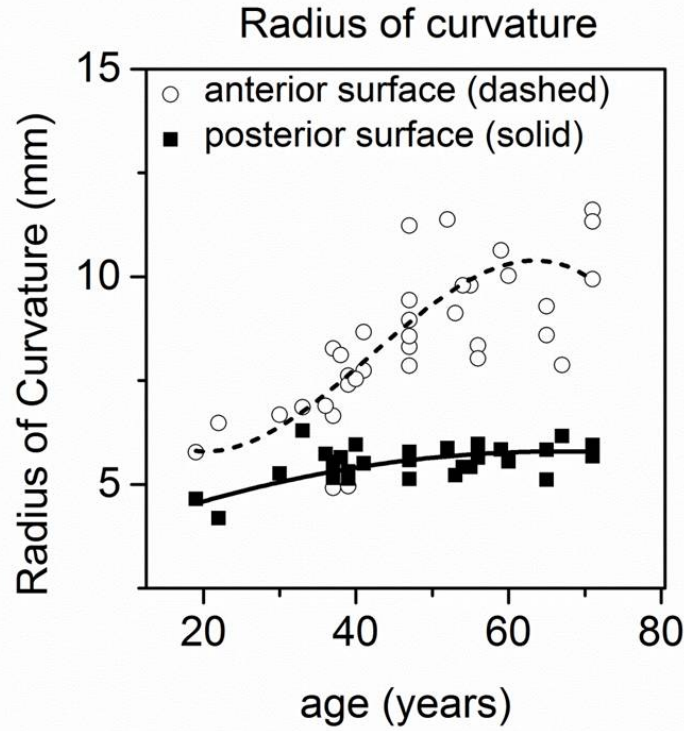


Figure 5.4: Anterior and posterior lens radii of curvature as a function of age. The anterior surface radius was fit by a 3rd order polynomial (R_{ant} (mm) = $-6E^{-5}x^3 + 0.0082x^2 - 0.2203x + 7.4769$; $r = 0.727$, $p < 0.001$; $x = \text{age}(\text{years})$). The posterior radius of curvature with age was fit by a 2nd order polynomial (R_{pos} (mm) = $-0.0006x^2 + 0.073x + 3.5139$; $r = 0.518$; $p = 0.001$; $x = \text{age}(\text{years})$).

5.3.4 Equivalent refractive index and average refractive index

The equivalent refractive index (figure 5.8A) ranged between 1.391 and 1.417 across lenses, and was 1.406 on average. The average refractive index (along the optical axis, figure 5.8B) ranged between 1.388 and 1.403, and was 1.394 on average. No statistically significant level of correlation was found between age and the refractive indices.

5.3.5 Experimental back focal length

The experimental back focal length (BFL) of the isolated lenses, obtained with LRT, increased significantly ($p < 0.001$) with age, both for a 2-mm (paraxial) and 4-mm (non-paraxial) pupil diameter (Figure 5.9). The non-paraxial focal length was relatively

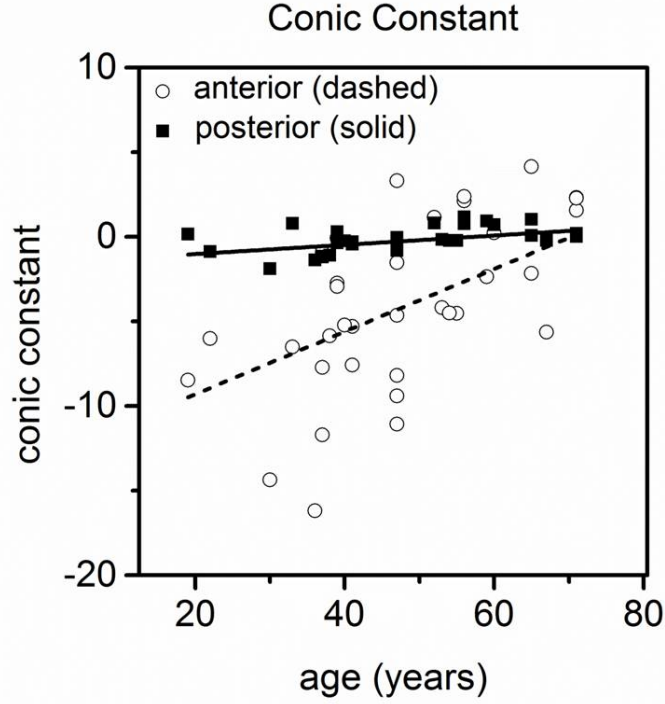


Figure 5.5: Anterior and posterior lens asphericity as a function of age, and linear fit to the data for the anterior lens ($k_{\text{ant}} = 0.228 \cdot \text{Age} - 14.853$; $r = 0.604$; $p < 0.001$), and for the posterior lens ($k_{\text{pos}} = 0.0275x - 1.4798$; $r = 0.5$; $p = 0.002$).

higher in younger lenses and changed with age at a lower rate (0.285mm/year) than the paraxial focal length (0.466mm/year). Therefore, the difference between paraxial and non-paraxial lens power decreased with age, approaching zero at around 60 years of age. The higher central lens power (i.e. a shorter BFL) in the paraxial zone for younger lenses is consistent with negative spherical aberration. Therefore, the decrease in the paraxial and non-paraxial BFL difference is consistent with of a shift of spherical aberration from negative to more positive values. Using equation 5.1, the 4th order spherical aberration Z_0^4 the Zernike was calculated for the lenses, showing a linear shift with age from negative values in young lenses to closer to zero in older eyes ($Z_0^4 = 0.0029 \cdot \text{Age} - 0.1904$; $r = 0.455$, $p = 0.022$).

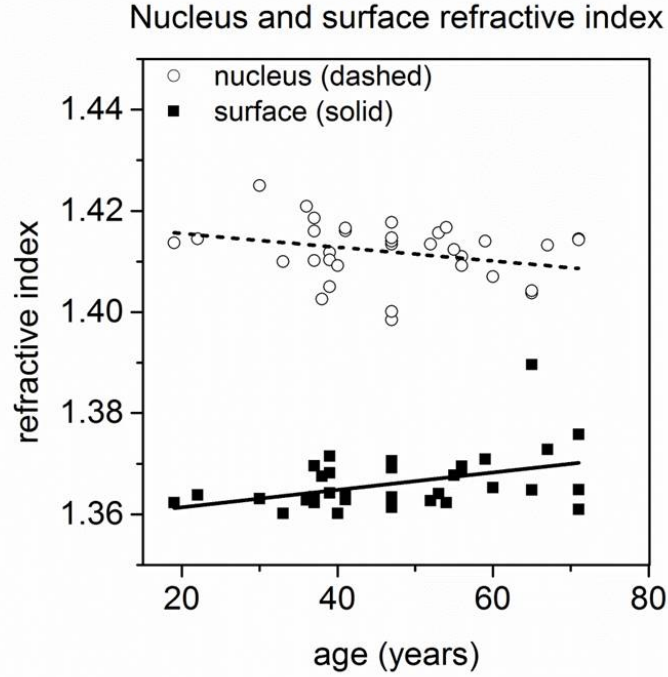


Figure 5.6: Nucleus and surface refractive index as a function of age, and linear fits to the data for the nucleus refractive index ($n_N = -8E^{-5} \cdot \text{Age} + 1.4157$; $r = -0.180$, $p = 0.3$) and the surface refractive index ($n_S = 0.0002 \cdot \text{Age} + 1.358$; $r = 0.409$, $p = 0.015$).

5.3.6 Spherical aberration: surface and GRIN contributions

Figure 5.10 shows ray-tracing estimates of the lens spherical aberration (SA) for a 6-mm pupil diameter as a function of age. Calculations were performed for the measured 3-D lens geometry and estimated GRIN distribution as well as for the same lens with a homogeneous equivalent refractive index. The spherical aberration shifts significantly towards less negative values with age, at higher rates when considering the reconstructed GRIN (0.041/year, $p < 0.001$) than with a homogeneous refractive index (0.0249/year, $p < 0.001$). The GRIN therefore plays a significant role in the negative values of the crystalline lens spherical aberration at all ages, but primarily in young lenses.

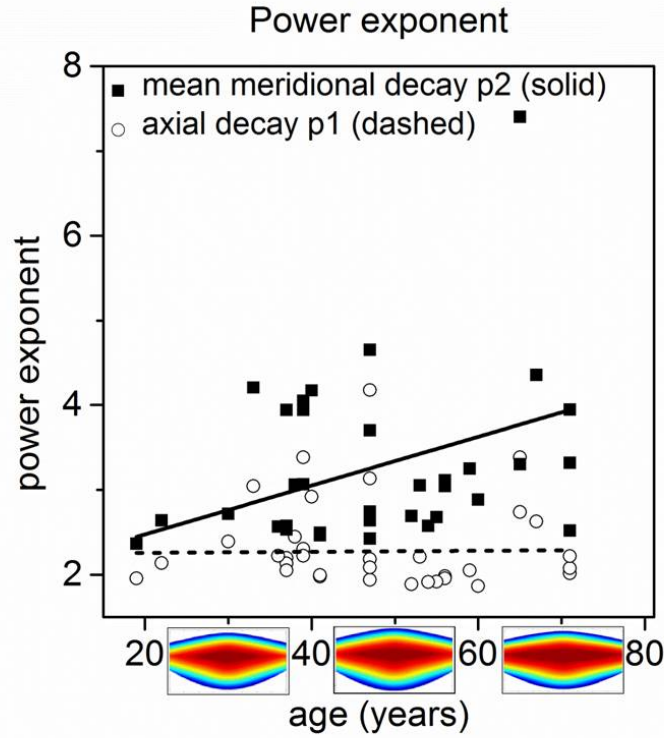


Figure 5.7: Power exponent in the axial direction (p1) and in the meridional direction (p2) as a function of age, and linear fit to the data for p1 (axial decay = $-9.9E^{-4} \cdot \text{Age} + 2.3878$; $r = -0.025$, $p = 0.89$) and p2 (meridional decay = $0.021 \cdot \text{Age} + 2.2973$; $r = 0.281$, $p = 0.1$). For p2, each symbol is the average across 18 meridians. A higher p is consistent to a flatter plateau in the GRIN, and rapid changes towards the surface. The colored insets illustrate a cross-section of the GRIN distribution in one meridian. Examples are for lenses of 30, 52 and 71 years

5.3.7 Lens surface topography: HO Zernike terms

Lens elevation high order Zernike terms: anterior and posterior lens relationships We tested the similarity of anterior and posterior lens topographies by evaluating the correlation of the high order Zernike terms in the anterior and posterior lens surfaces. We found that several high order Zernike terms were statistically significantly correlated in anterior and posterior lens surfaces. As an example, Figure 5.11 shows linear regressions between anterior and posterior surface for vertical trefoil (Z_3^-3), 3rd order RMS, RMS Coma and 4th order RMS. Table 5.1 shows the mean and standard deviation of different parameters describing the shape of the anterior and posterior surface in terms of RMS,

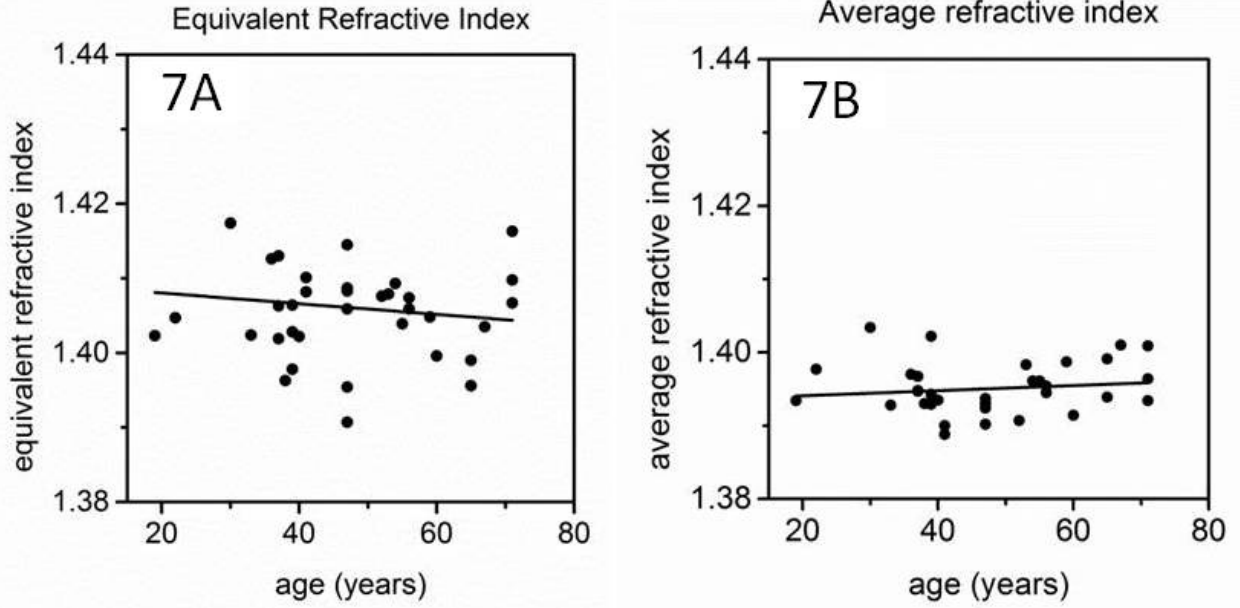


Figure 5.8: (A) Equivalent refractive index estimated from the reconstructed GRIN as a function of age, and linear fit to the data (equivalent $RI = -5E^{-8} \pm 0.06 \cdot \text{Age} + 1.405 \pm 3.1$; $r = -0.0001$, $p = 1$). (B) Average refractive index, calculated from the OCT images along the optical axis of the lens as a function of age, and linear fit to the data (average $RI = 3E^{-5} \cdot \text{Age} \pm 0.11 + 1.393 \pm 5.4$; $r = 0.14$, $p = 0.44$).

and its level of correlation with age. Also the last column shows the correlation between anterior and posterior lens surface elevation. We have found a strong correlation between surfaces of the RMS for 3rd order, 4rd order, coma, trefoil and spherical terms.

5.3.8 Lens elevation high order Zernike terms: changes with age

Figure 5.12 shows several Zernike surface elevation terms (in terms of RMS) as a function of age: (A) RMS coma (B) RMS spherical (C) 3rd order RMS, and (D) 5th order RMS, both for the anterior and posterior lens surfaces. RMS coma (figure 5.12 A) decreased highly statistically significantly with age at a rate of $0.087 \mu\text{m}/\text{year}$ and $0.123 \mu\text{m}/\text{year}$ for anterior and posterior surface, respectively. RMS spherical (figure 5.12 B) and 3rd order RMS (figure 5.12C) decrease significantly with age for the anterior surface, but not for the posterior surface (RMS spherical: slope= $-0.175 \mu\text{m}/\text{year}$ for anterior surface,

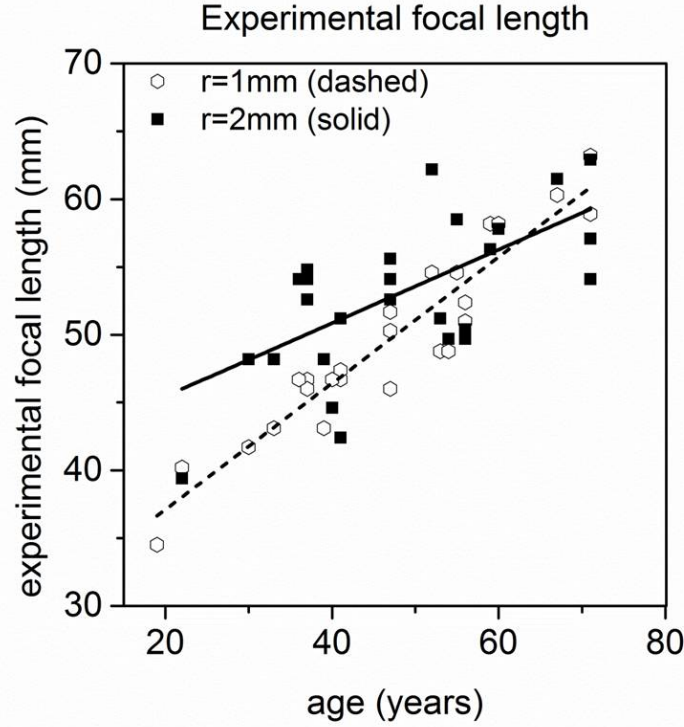


Figure 5.9: Experimental back focal length for a 1-mm pupil radius (paraxial) and a 2-mm pupil radius (non-paraxial), and linear fit to the data for the paraxial ($\text{exp.bfl}_{1\text{mm}}(\text{mm}) = 0.466 \cdot \text{Age} + 27.81$; $r = 0.95$, $p < 0.001$) and non-paraxial data ($\text{exp.bfl}_{2\text{mm}}(\text{mm}) = 0.285 \cdot \text{Age} + 39.070$; $r = 0.66$, $p < 0.001$).

slope $= -0.05 \mu\text{m} / \text{year}$ for posterior surface; 3rd order RMS: slope $= -0.083 \mu\text{m} / \text{year}$ for anterior surface). 5th and 6th order RMS do not show statistically significant changes with age. Table 1 (4th and 5th columns) shows the correlations with age for different RMS orders.

5.3.9 Relative contribution of different Zernike terms to the lens surface elevations

Figure 5.13 shows the average relative contribution of lower and higher order Zernike terms in both surfaces. Relative contributions are accounted for in terms of variance (RMS2). Astigmatism is the predominant term both in the anterior surface (54.96%) and in the posterior surface (62.95%), followed by spherical (25.10%, 17.85%), coma (10.91%, 6.38%), trefoil (6.01%, 8.73%), tetrafoil (1.76%, 2.80%) and 4th order astig-

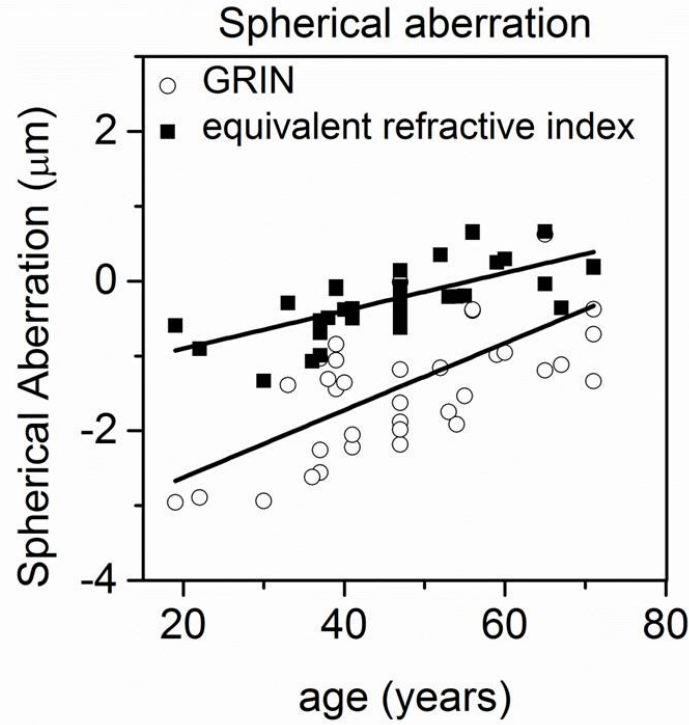


Figure 5.10: Estimated spherical aberration as a function of age for GRIN lenses and lenses with a homogeneous refractive index, and linear fits for GRIN lenses ($SA_{\text{GRIN}} = 0.041 \cdot \text{Age} - 3.4075$; $r = 0.654$, $p < 0.001$) and homogeneous refractive index lenses ($SA_{\text{homogeneous}} = 0.0249 \cdot \text{Age} - 1.4043$; $r = 0.696$, $p < 0.001$). The GRIN shifts the SA towards more negative values.

matism (1.27%, 1.30%), for anterior and posterior surface, respectively. We also studied the change of these relative contribution with age (not shown) and found that for the anterior surface, the two predominant terms, astigmatism and spherical, change with age. The relative proportion of astigmatism increased with age at a rate of 0.702 /year ($r=0.427$, $p=0.009$), while the percentage of spherical term decreased significantly with age at a rate of -0.616/year ($r= 0.474$, $p=0.004$). For the posterior surface, we only found a significant change in the proportion of coma (slope=-0.307/year, $r= 0.450$, $p=0.006$).

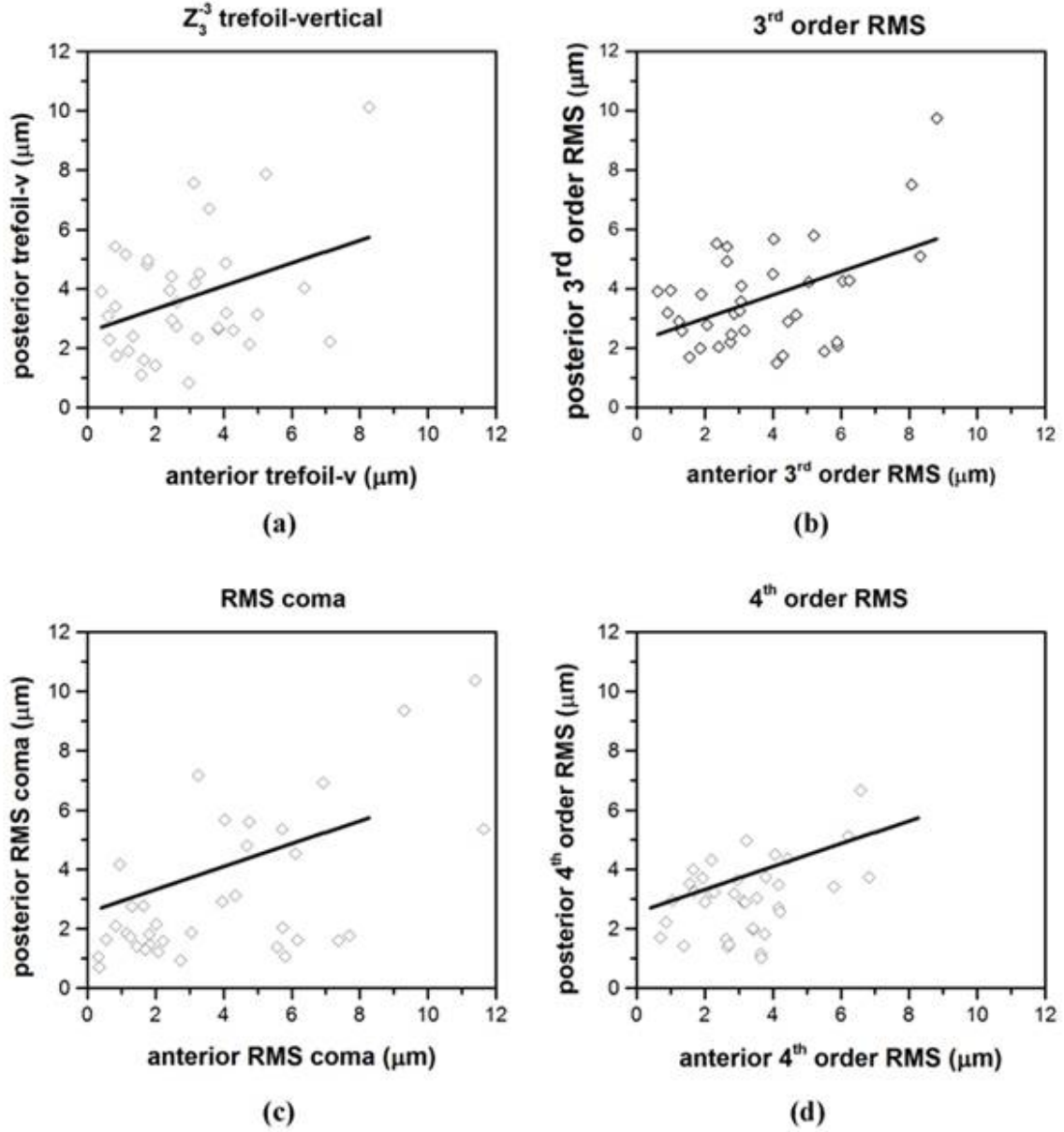


Figure 5.11: (a) Trefoil-v : Anterior vs. Posterior (Slope= 0.387, $r = 0.467$, $p = 0.004$) (b) RMS 3th order terms: Anterior vs. Posterior (Slope= 0.019; $r = 0.477$, $p = 0.003$; (c) RMS Coma: Anterior vs. Posterior (Slope= 0.387; $r = 0.617$, $p = 0.0001$) (d) 4th order RMS: Anterior vs. Posterior (Slope= 0.387; $r = 0.423$, $p = 0.010$).

5.4 Discussion

We have used spectral Optical Coherence Tomography to image isolated crystalline lenses of different ages. These measurements allowed quantification of the lens shape

Table 5.1: **Parameters for the LPL model**

HO Zernike coefficients	anterior (μm)	posterior (μm)	anterior vs. age r and p	posterior vs. age r and p	ant. vs. pos r and p
RMS 3rd order Zernike Terms	3.294 ± 1.822	3.697 ± 1.770	r= 0.564 p= 0.0001*	r= 0.373 p= 0.025	r= 0.477 p= 0.003*
RMS 4th order Zernike Terms	3.227 ± 1.779	3.077 ± 1.247	r= 0.721 p= 0.0001*	r= 0.277 p= 0.102	r= 0.423 p= 0.010*
RMS 5th order Zernike Terms	1.151 ± 0.407	1.381 ± 0.727	r= 0.290 p= 0.087	r= 0.338 p= 0.023	r= 0.047 p= 0.786
RMS 6th order Zernike Terms	1.216 ± 0.453	1.076 1.076	r= 0.149 p= 0.384	r= 0.061 p= 0.723	r= 0.236 p= 0.166
RMS Trefoil	2.553 ± 1.300	3.683 ± 2.038	r= 0.266 p= 0.117	r= 0.011 p= 0.949	r= 0.363 p= 0.030*
RMS Tetrafoil	1.380 ± 0.824	2.085 ± 1.299	r= 0.278 p= 0.101	r= 0.105 p= 0.544	r= 0.004 p= 0.980
RMS Spherical	5.219 ± 5.708	5.266 ± 3.232	r= 0.439 p= 0.007*	r= 0.226 p= 0.185	r= 0.375 p= 0.024*
RMS Coma	3.440 ± 2.918	3.147 ± 2.467	r= 0.582 p= 0.0001*	r= 0.515 p= 0.001*	r= 0.617 p= 0.0001*
RMS 4th order astigmatism	1.173 ± 0.826	1.420 ± 0.728	r= 0.189 p= 0.270	r= 0.022 p= 0.897	r= 0.174 p= 0.310

* Statistical significance (Following Bonferroni correction).

and 3D GRIN distribution, and their changes with aging. Computational ray tracing on these experimental data allowed evaluating the relative contribution of lens shape and GRIN to spherical aberration and astigmatism, as a function of age.

5.4.1 Crystalline lens shape changes with age

Radius of Curvature

In agreement with previous reports [73] [32], we found a flattening of the isolated lens with age, with a larger increase of the radius of curvature by 45% in the anterior lens and 20% in the posterior lens, between the ages of 20 and 60. Our data also show the biphasic behavior reported by Glasser and Borja on isolated crystalline lenses, peaking

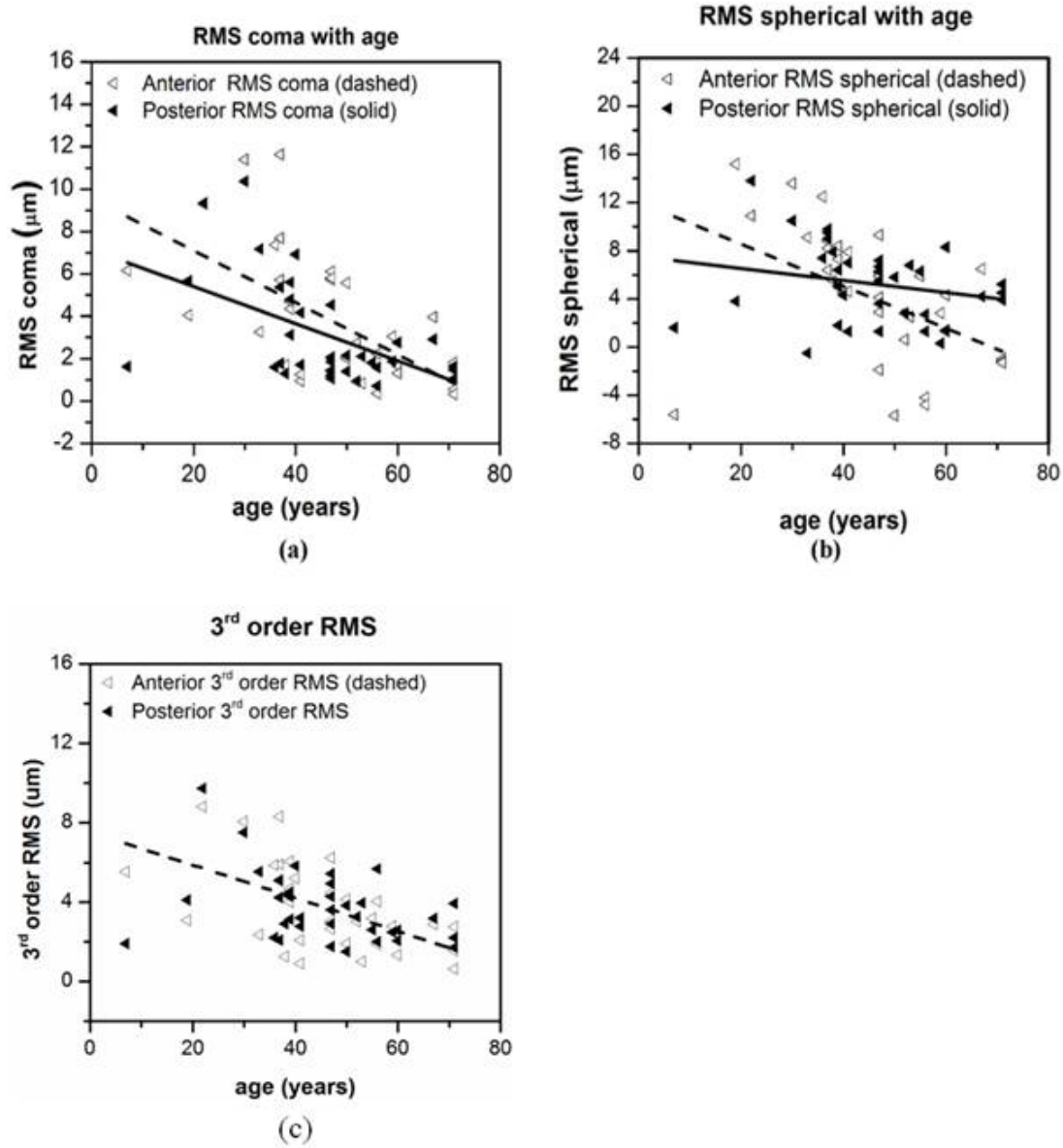


Figure 5.12: Lens elevation high order RMS terms; RMS coma: slope= $-0.087\mu\text{m}/\text{year}$, $r = 0.582$ and $p=0.0001$ for anterior surface, slope= $-0.123\mu\text{m}/\text{year}$, $r = 0.515$ and $p=0.001$ for posterior surface; (b) RMS spherical: slope= $-0.175\mu\text{mm}/\text{year}$, $r = 0.439$, $p=0.007$ for anterior surface only (c) 3rd order RMS: slope= $-0.083\mu\text{m}/\text{year}$, $r = 0.564$ and $p=0.0001$ for anterior surface only

around 60 years of age. Data are also consistent with measurements in vivo using corrected Scheimpflug imaging[116], Purkinje imaging[214] and OCT [67] as function of accommodation and/or aging.

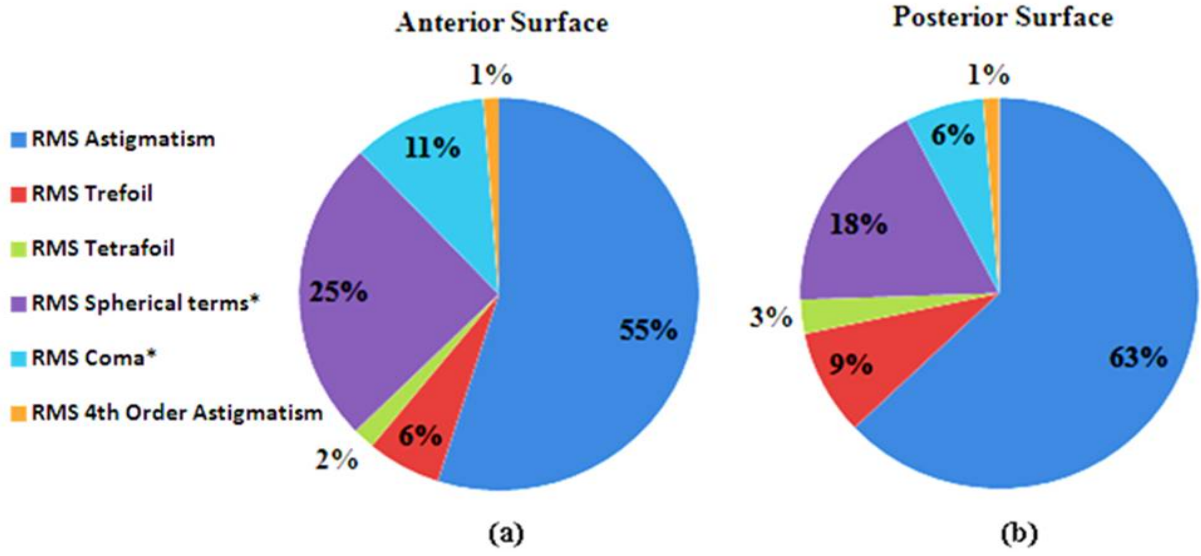


Figure 5.13: Relative contribution of different Zernike terms to the overall surface elevation maps (in terms of RMS2) with an asterisk the terms that change statistically significantly with age.

In the young lens, the radii of curvature that we report here in vitro are close to those reported in vivo for the maximally accommodated state. Dubbelman et al.[58] reported anterior and posterior lens radii of curvature of 7 mm and 5 mm, respectively, in a 6-D accommodating 25-year old lens in vivo, in good agreement to the 6.2 mm, and 5 mm for anterior and posterior radius of an interpolated 25-year old lens in our study in vitro. Radii of curvature found in vitro are similar in magnitude to those reported in vivo for older lenses by Dubbelman who reported anterior and posterior lens radii of curvature of 9.5 mm and 5.7 mm, respectively, studying a 60-year old lens, in agreement with 10.9 mm, and 5.6 mm for a lens of the same age in our study.

Conic Constant

Our data are comparable to the conic constants reported by Dubbelman et al.[57] for in vivo unaccommodated lenses (-4 to 0.06 for the anterior surface and -3 to 1.96 for the posterior surface, respectively), between 16 and 65 years. As expected, the agreement between our in vitro data and in vivo data from Dubbelman is better in older eyes, due to the differences in lens shape between the relaxed accommodation state (in vivo) and the maximally accommodated state (in vitro) in young lenses. In fact, the anterior conic lens constants in the isolated young lenses in our study are in excellent agreement with those reported in maximally accommodated young lenses in (-12 to -3, in an age range

of 24 to 34 years). The positive anterior lens conic constant in the older eyes is also consistent with results from a study by Manns et al.[145] who reported an average value of 4.27 in a group of old isolated human lenses (average age 76.4 years).

5.4.2 GRIN distribution and equivalent refractive Index

Age dependence on nucleus/surface refractive index

We found a tendency for the nucleus refractive index to decrease with age, and the surface refractive index to increase with age, although only the changes in the surface refractive index showed statistical significance. The range of nucleus and surface refractive index agree with those reported by de Castro et al.[48] using a similar GRIN reconstruction technique, but based on 2-D cross sectional OCT images and a smaller lens sample (n=9), although the relatively high scattering of the values prevented from observing a clear behavior. The higher number of lenses of the current study (n=35) and smaller uncertainties due to the 3-D nature of the data allowed us to reach near statistical significance. These data are in agreement with the change in surface refractive index of human lenses (anterior pole only) reported by Pierscionek[194] (n=14), using destructive methods, and the decrease of nucleus refractive index reported by Jones et al.[104] (n=20), using an MRI approach. Also, the decrease of the nucleus refractive index with age is consistent with the hypothesis by Moffat et al.[161] that the lens nucleus experiences an age-related decrease in the soluble protein.

Power exponent change

The shape of the GRIN profile changed with age, predominantly in the meridional direction. For some lenses the meridional power exponent p_2 changed across meridians, but the magnitude of change was not correlated with age or surface shape (radii, conic constant). The findings of a more distributed index in younger lenses and an increase of the central plateau with age agree with earlier findings by Hemenger et al.[89], Moffat et al.[161], Jones et al.[104] and de Castro et al.[48] Also, Kasthurirangan et al.[107] reported higher exponents in the equatorial direction, than in the axial direction in both young (accommodated) lenses and older lenses, with a larger increase for the equatorial power exponent (from 5.09 to 10.28) than for the axial power exponent (from 4.04 to 6.7) with age. Although the reported power exponent values of the GRIN should better describe those of fully accommodated lenses, as isolated lenses approach a fully accommodated state, we anticipate that the power exponent changes with age in the current study are largely associated with age, and not the accommodation state. This

is supported by the finding in chapter 4, of the GRIN distribution with simulated accommodation of cynomolgous monkey lenses in vitro, showing no significant differences in the meridional or axial power exponent with accommodation. It is noticeable that, despite the variability of the results, our estimated power exponents are lower (particularly in older lenses) than earlier reports[48][107], although they support the increase of the meridional exponent with age. To rule out potential differences associated with the reconstruction model we repeated the reconstruction using a subset of 2D cross-sections from our complete 3D dataset, and a 3-variable model with only one power exponent. In comparison, the power exponent derived from the 2D model showed an age-dependence similar to the meridional power exponent derived from the 3D model. This finding suggests that the differences with respect to earlier results obtained from 2D data are likely related to the measured sample, and not to the 2D versus 3D nature of the data. Our power exponent findings reconcile the apparent contradiction in the expression reported by Charman and Atchison in a recent article[41], where if the power exponent increases (as typically reported to happen with age), then the average refractive index should show a major increase with age. In contrast, another study reports a decrease of the average refractive index with age[45]. Since the average refractive index is only dependent on the axial power exponent, our finding that axial power exponent is practically constant with age predicts only small differences in the average refractive index with age. Our experimental measurements of average refractive index show a minor increase with age (increase +0.002 from 20 to 70-year old lenses) in good agreement with the predicted increase (+0.003 in the same age range) by Charman and Atchison. It should be mentioned that our average refractive index results (slightly lower than previous reports) are derived directly from the OCT images (i.e. independent from the GRIN reconstruction).

Equivalent refractive index

We only found small, not statistically significant, decrease of the equivalent refractive index with age. Earlier literature found evidence for a decreasing equivalent refractive index with age in vivo[57][70][12] but the in vitro results vary[72][32]. While Pierscionek[194] (isolated lenses) and Glasser et al.[73] (isolated lenses) did not find an age dependency of the equivalent refractive index, Borja et al.[32] (isolated lenses) found a significant decrease with age. We expect that this variation is largely associated with age-related changes, and not due to accommodation, and therefore possibly extrapolated in vivo, given the independence of refractive index with accommodation recently reported in monkey lenses[142]. A reason for the discrepancy across data in the equivalent refractive index values in the literature (and its change with age) may be that indirect measurements of the lens power, such as those obtained by comparing lens

shape and eye refraction (in vivo) and lens shape and power (in vitro), may be affected by the pupil diameter used (due to the spherical aberration of the lens). For example, Dubbelman et al.[57], calculated the equivalent power from refraction and geometrical data using a 3 mm-pupil; Borja et al.[32] calculated the refractive index of crystalline lenses in vitro from the measured equivalent power, thickness, diameter, and radius of curvature, using a 4 mm ring of light as an entrance pupil for the power measurements. Unlike in previous studies, our equivalent refractive index is obtained directly from our estimated GRIN distribution. In fact, a simulation of the estimation of the equivalent refractive index based on a comparison of the lens power estimated from the lens geometry and GRIN, and the lens geometry and an equivalent refractive index reveals that the estimated equivalent refractive index (and its change with age) is dependent on the pupil diameter used in the calculation. When comparing the results for a 1.5-mm pupil diameter with the results for a 3-mm pupil diameter, we found the decrease in the equivalent refractive index to be 35% larger for the smaller pupil diameter. The actual values of the equivalent refractive index changed around 1%. Similarly, using the results from the LRT measurements in the estimation, the values for the equivalent refractive index were also higher for smaller pupil diameter (2-mm), and increasing the diameter (to a 4-mm pupil) led to a larger decline (by 20%) with age. Differences in the rate of variation of the equivalent refractive index may therefore be explained by differences in the pupil diameter in the measurements and/or model.

subsectionLens GRIN and shape contributions to Spherical Aberration The spherical aberration estimated for the isolated lenses assuming a homogeneous equivalent refractive index was negative in 71.4% of the lenses. When the reconstructed GRIN was assumed instead, the SA was shifted to more negative values, with 97.1% of the lenses showing negative spherical aberration. In both cases the SA shifted towards less negative values with age (at rates of $0.025\mu\text{m}/\text{year}$ for the homogeneous lens and $0.041\mu\text{m}/\text{year}$ for the GRIN lens). The linear fit crossed zero spherical aberration at approximately ages 57 and 84, for homogeneous and GRIN lens, respectively. Although the compensatory effect of GRIN occurred even in older lenses, this effect decreased significantly with age, and it should be associated to the change in the GRIN distribution (flatter central index). The negative SA in young human lenses has been reported in numerous studies[8][232], as well as its shift towards less negative values with age [158, 9]. Furthermore, previous studies have reported a compensatory role of GRIN in different species[32][47][40][269]. The current study confirms experimentally the relative contribution of shape and GRIN to spherical aberration in human lenses. It should be noted that the current study exaggerates the negative spherical aberration in young lenses, as these are maximally accommodated. With accommodation the SA shifts towards more negative values[87][67].

5.4.3 Lens Topography

In a later study, which is included in this chapter, we have studied the 3-D surface topography of 36 isolated human crystalline lenses and evaluated the correlations between the anterior and posterior lens surface shape and their changes with age. The anterior and posterior crystalline lens surfaces have been usually described by the fitting parameters of conic functions, using measurements normally obtained from single meridians[58] [67] [69]. However, the direct measurements of lens surface elevations in 3-D in our study reveal that the human crystalline lens shows in fact non-spherically symmetric surfaces. Our data show the presence of astigmatism and, to a lesser extent, other high order contributions (trefoil, spherical, coma and tetrafoil). Previous studies comparing corneal and total astigmatism predicted the presence of astigmatism in the lens, which in many cases, has a compensatory effect for corneal astigmatism [8][110].

The surface astigmatic power in our lens sample ranged from 0.046 to 1.185 D in the anterior lens surface, and from 0.013 to 1.118 D in the posterior lens surface. However, the presence of a gradient refractive index (GRIN) distribution may play a role in the overall astigmatism of the lens. Whether GRIN has a compensatory role in lens astigmatism will be investigated further in chapter 6. We can only speculate on the factors contributing to lobular Zernike surface terms such as trefoil and tetrafoil, which might be related to suture branching. In a work by Gargallo et al.[69], lens aberrations in several species with Y-suture branches were analyzed for their relationship with suture distribution. A high degree of correlation between suture orientation and the axis of no rotationally symmetric wavefront aberrations was found. We did not find a strong correlation between surface astigmatism and age, but there were statistically highly significant correlations in the anterior lens surface for RMS 3rd and RMS 4th order Zernike term, RMS spherical, and RMS coma ($p < 0.008$), which all decreased with age. No significant correlation was found with age for surface tetrafoil and high order lobular terms, in general. Although, if lobular terms are associated with the presence of lens sutures, we would have expected an increase of the higher order terms with age, as the number of suture branches increases with aging. However, it has been shown that the formation of new branching structures is highly slowed down at older ages[80]. Our OCT data did not provide a three-dimensional view of the lens branching, although more refined imaging modalities revealing these structures could allow modeling of the potential relationships between branching and trefoil and tetrafoil terms. A previous study from our laboratory showed the ability to measure crystalline lens anterior and posterior surface elevation maps in vivo using SD-OCT, and demonstrated it on three human young subjects in vivo. Besides the fully accommodated state of the isolated lens and potential changes post-mortem, there may be other differences between in vivo and ex vivo conditions.

For example, the zonular tension in the un-accommodated condition may cause further high order contributions to lens surface elevation in vivo (besides the flattening of the lens). Also, the up-right orientation of the lens in vivo (as opposed to the horizontal one in our ex vivo measurements) may create some changes associated to gravity. Two of the three eyes measured in vivo in our earlier study showed perpendicular astigmatism axes in the anterior and posterior lens, whereas the astigmatic axes in our study tended to differ on average only 36.46 deg. In 13 of the 36 lenses the relative angle was ≥ 45 deg, and 6 lenses had a relative angle ≥ 70 deg. The average relative angle was larger for younger lenses (41.74 deg for lenses ≤ 50 years) than for older lenses (28.1 deg for lenses ≥ 50 years), indicating a higher compensation of astigmatism between the lens surfaces at a younger age. Finally, the lack of a reference for the lenses ex vivo prevented from analyzing the relative axis of astigmatism with corneal astigmatism (and therefore evaluation of potential compensatory effects between the surfaces of the cornea and lens and their changes with aging) and possible changes in the anterior lens surface astigmatic axis with aging. For example, it is well known that corneal astigmatism changes from with-the-rule to against-the-rule astigmatism with aging[85][92][5]. Previous studies suggested that the astigmatism axis of the anterior lens surface is vertical[176][149][112]. We did not find a significant correlation between anterior or posterior astigmatism axis and age. However, the tendency of the relative angle to decrease with age could be interpreted as a decrease of the anterior by posterior astigmatism balance with age. Regardless potentially additional information provided by future in vivo studies of lens surface topography and its changes with age, our study shows that astigmatism and high order terms may be considered when trying to predict optical quality of a phakic eye at an individual level, based on anatomical information. This information could be used in a ray tracing analysis that incorporates data from cornea and lens GRIN distribution[145][112][2][118].

6

Astigmatism of the in vitro human lens: Surface and gradient refractive index contributions

This chapter is based on the paper in revision by Birkenfeld et al. "Astigmatism of the in vitro human lens: Surface and gradient refractive index contributions" (Invest Ophthalmol, 2014). The author of this thesis (i) did the sample preparation, (ii) performed the experiments (in collaboration with Alberto de Castro), (iii) performed data processing (in collaboration with Alberto de Castro), (iv) wrote additional Matlab code, (v) analyzed the results, and (vi) wrote the manuscript (in collaboration with Alberto de Castro and Susana Marcos). In this study, we distinguished between lens astigmatism and surface astigmatism to find how the different lens characteristics relate to each other. We found a significant decrease of the astigmatic magnitude of the anterior surface with age, while the posterior surface and the GRIN of the lens did not change significantly with age. The total astigmatism of the lens decreased from negative values towards zero. These results suggest that the anterior lens surface has a larger impact to the total lens astigmatism, and that the influence of the GRIN to the lens astigmatism is limited.

6.1 Introduction

In the last chapter, we used 3-D OCT and LRT measurements to fully characterize the surface geometry, topography, the GRIN distribution, and optical properties of 35 human donor lenses with age. We were able to investigate the influence of shape and GRIN on the spherical aberrations of human lenses. As a result we concluded that both, geometrical changes and GRIN contributes to the age-dependent shift of negative spherical aberration. In this chapter, we use the collected data to draw conclusion about the impact of lens surfaces and GRIN on the total astigmatism of the lens.

While the contributions of the cornea and lens, and the lens shape and GRIN to spherical aberration have been relatively well studied, a similar analysis has not been performed on the relative contributions to astigmatism, likely because most lens studies to date only had access to 2-D data. Compensatory effects of anterior and posterior corneal astigmatism have been reported[55, 60][61, 180, 205, 209]. Also, longitudinal studies report changes in the cornea from with-the-rule (direct) to against-the-rule (indirect) astigmatism with age[21]. Measurements of total and corneal astigmatism from ocular refraction[108, 109] or ocular aberration measurements[8, 110] suggest at least partial compensation of total and internal astigmatism. In those studies, the contributions of the posterior cornea and lens to the internal optics cannot be isolated. In an interesting study using an ophthalmophakometric technique, Elawad et al.[63] estimated the ocular component contributions to residual astigmatism in human eyes, and found that, whilst the astigmatic contributions of the posterior corneal and lens surfaces were found to be predominantly inverse, that direct astigmatism came from the anterior lens surface. Similar conclusions were reached by Dunne et al.[61] in a later work, although the authors recognized that the method was indirect and prone to accumulated experimental errors. To our knowledge, the only direct measurement of crystalline lens surface astigmatism on ex vivo human lenses comes from surface topographic analysis (using OCT)[243]. In that study we found that astigmatism was the predominant lens surface aberration. We did not find a significant change in the amount of surface astigmatism aberration with age, although the relative angle of astigmatism between the anterior and posterior lens surfaces tended to decrease with age, indicating a potential decrease in the compensatory effects of anterior and posterior lens astigmatism with age. The study did not consider potential effects of the GRIN distribution with age. Access to both, lens shape and GRIN in 3-D opens the possibility of evaluating the relative role of lens surface and GRIN astigmatism to lens astigmatism, and the potential changes with age, as in similar analysis of contributors of spherical aberration.

6.2 Methods

6.2.1 Human lens samples and preparation

All human donor eyes were received from the Transplant Service Foundation (TSF) Eye Bank in Barcelona, Spain. During the transportation eyes were packed individually in sealed vials at 4C, wrapped in gauze which was well soaked in a preservation medium (DMEM/F-12, HEPES, no phenol red, GIBCO). Presence of any form of cataract was considered an exclusion criterion for the study. Prior to shipment the corneas had been removed (for corneal transplant purposes) and in some cases sections of the sclera. However, the vitreous and the choroid were preserved and provided a safe transportation environment for the lens. All lenses arrived 1 days post-mortem, and were measured within 24 hours. Thirty-five eyes from 30 human donors were used in the study. Ages ranged between 19 and 71 years. Prior to the experiment, the lens zonules were carefully cut and the lens was extracted from the eye with soft tweezers and handled mainly using the remaining zonules rather than touching the lens capsule. After extraction the lens was immediately immersed in DMEM at room temperature. During the measurements, the lens was placed on a ring in a DMEM filled cuvette. The whole measurement took up to 2 hours. Swollen or damaged lenses were identified with the OCT images and excluded from the study. Handling and experimental protocols had been previously approved by the Institutional Review Boards of TSF and CSIC. Methods for securing human tissue were in compliance with the Declaration of Helsinki.

6.2.2 Optical Coherence Tomography Imaging

The cuvette containing the lens was placed on a horizontal platform and imaged using a custom developed high resolution spectral OCT system described in detail elsewhere (Grulkowski2009). A mirror system above the platform assured that the OCT beam was directed towards the upper lens surface. The system uses an 840-nm SLD diode as illumination source and obtained 3D images composed of 1668 A-Scans, and 70 B-Scans on a 12x12 mm lateral area. The acquisition time was 4.5 s, and the axial resolution was calculated to 6.9 μ m in tissue. The lens axis could be easily centered with the OCT set to visualization mode, which displays two orthogonal B-Scans in an interval of 0.5 s. The manually adjustable platform on which the cuvette was set up, was aligned until a specular reflection was seen from the surfaces of the lens. Because of the thickness of the lens, we had to take images in two different focal planes, to allow visualization of both lens surfaces and the cuvette holding the lens. The images were merged into one complete 3D image (anterior surface, posterior surface and cuvette surface) using

a custom-developed merging algorithm. The lens was first completely imaged with the anterior surface facing the OCT beam, and was then flipped around a predetermined axis and imaged again with the posterior surface up.

6.2.3 Laser Ray Tracing

After the OCT measurements, the focal length of each lens was calculated with the custom developed Laser Ray Tracer (LRT). The crystalline lenses were placed horizontally in a cuvette (anterior up position), positioned on a stable platform. The bottom of the cuvette consisted of optical glass, and a motorized CMOS camera was placed right under the cuvette to capture a series of through-focus images from the beam directed through the lens. The lenses were aligned with the LRT, such that the lens and principal ray were collinear to the center of the CMOS camera, in the entire focus range (which was the range of the motor, 35mm). With this set up the paraxial and non-paraxial focal lengths of each lens were measured projecting rings of light of different diameters (2 and 4 mm) onto the crystalline lens' upper surface. The estimated precision of the focal length measurements was 0.8 mm. The focal length data was used as input data for the GRIN reconstruction algorithm as described in.

6.2.4 Image Processing and GRIN Reconstruction

All OCT images were corrected from optical distortions and all surfaces (lens and cuvette) were fitted by Zernike polynomials (up to 7th order) within a 6-mm pupil. The GRIN was reconstructed by means of a search algorithm using the optical path measured from the OCT images and the measured back focal length. The gradient refractive index is described as a 4-variable model expressed in polar coordinates centered in the center of the lens as

$$n(\rho, \theta) = n_N - \Delta n \cdot \left(\frac{\rho}{\rho_S} \right)^{p(\theta)} \quad (6.1)$$

where n is the refractive index of the nucleus, Δn the difference between the refractive index of surface and nucleus, ρ is the distance between nucleus and surface, and p is the exponential decay for axial ($p1$) and meridional ($p2$) direction. The axial decay $p1$ is constant across meridians while $p2$ can vary to account for differences between meridians. The change of the power exponent $p2$, the decay of the GRIN in the meridional direction, is an indicator of the GRIN contributor to astigmatism. A constant $p2$ across the lens is indicative of no GRIN astigmatism. A sinusoidal variation of $p2$ across lens meridians is indicative of the presence of GRIN astigmatism. In all computations The center of the lens is assumed to be placed at a distance from the anterior vertex equal to 0.41 times

the central thickness of the lens (Rosen2006).

6.2.5 Calculation of Lens Astigmatism (magnitude and axis)

Lens GRIN astigmatism

The lens total astigmatism was calculated using the measured lens shape and the reconstructed GRIN by means of a computational ray tracing analysis[47]. The calculated wave aberrations were fit by Zernike polynomials and the magnitude of astigmatism C and angle α were calculated[250] as:

$$C = -2 \cdot \sqrt{J_{45}^2 + J_{180}^2} \quad (6.2)$$

$$\alpha = \frac{1}{2} \tan^{-1} \left(\frac{J_{45}}{J_{180}} \right) \quad (6.3)$$

Where J_{180} is the power at axis $\alpha = 0^\circ = 180^\circ$ and J_{45} is the power at axis $\alpha = 45^\circ$,

$$J_{180} = -\frac{2 \cdot \sqrt{6} \cdot Z_2^{-2}}{r^2} \quad (6.4)$$

$$J_{45} = -\frac{2 \cdot \sqrt{6} \cdot Z_2^2}{r^2} \quad (6.5)$$

where Z_2^{-2} and Z_2^2 are the corresponding astigmatism Zernike terms, and r the pupil radius (3 mm in this study).

Equivalent index lens astigmatism

The lens total astigmatism assuming an equivalent index of refraction was calculated (equivalent index lens astigmatism). By comparing the lens total astigmatism with GRIN and with a homogeneous index we will assess the contribution of GRIN to the lens astigmatism.

Surface Lens Astigmatism

The lens surface astigmatism C_S was calculated as

$$C_S = (n_2 - n_1) \cdot \left(\frac{1}{R_x} - \frac{1}{R_y} \right) \quad (6.6)$$

where R_x is the astigmatic angle and R_y is the axis perpendicular to R_x . n_1 and n_2 are the refractive indices of the lens immersion medium and the lens cortex, respectively. The index of refraction of the lens cortex (n_2) was taken directly from our results from the GRIN reconstruction, individually for every lens. The refractive index of aqueous (n_1) was assumed to be 1.336[142].

Relative Astigmatic Angle

Since the orientation of the isolated lens during the measurements is arbitrary, i.e. not corresponding to its actual orientation in vivo (up, down, nasal, temporal), the calculated axis of astigmatism of the lens surfaces is arbitrary. However, the relative angle between the different axes of astigmatism (anterior, posterior, GRIN, and total astigmatic axis (the latter for GRIN lens and equivalent refractive index lens)) can be computed, and defined, in a range between 0 and 90 deg. The axis of astigmatism of the GRIN is defined as the meridian with the minimum value of p_2 (exponential decay of the refractive index in meridional direction) which is associated with the most increase in lens power.

Power Vector Analysis

To illustrate the magnitude of astigmatism and angle, the results are presented using a power vector graph (Thibos01, Liu11). For this we assumed that the axis of astigmatism of the anterior lens surface is aligned vertically in all lenses. This way we could see how the GRIN astigmatism and the posterior surface lens astigmatism is rotated with respect to the vertical aligned anterior surface, and how these relative angles change with age. Furthermore, we could compare the lens astigmatism of the GRIN lens and the corresponding homogenous lens.

6.3 Results

6.3.1 Change of magnitude of astigmatism with age

Figure 6.1 shows the change of the anterior and posterior lens surface astigmatism with age, fitted by linear regression. While the posterior lens surface astigmatism does not change significantly with age ($r=0.189$, $p=0.276$), the anterior lens surface astigmatism decreases significantly with age ($r=0.397$ and $p=0.018$). Figure 6.2 a shows the change of the lens magnitude of astigmatism with age for the crystalline lenses with the reconstructed GRIN ($r=0.359$, $p=0.034$) in comparison to the lens magnitude of

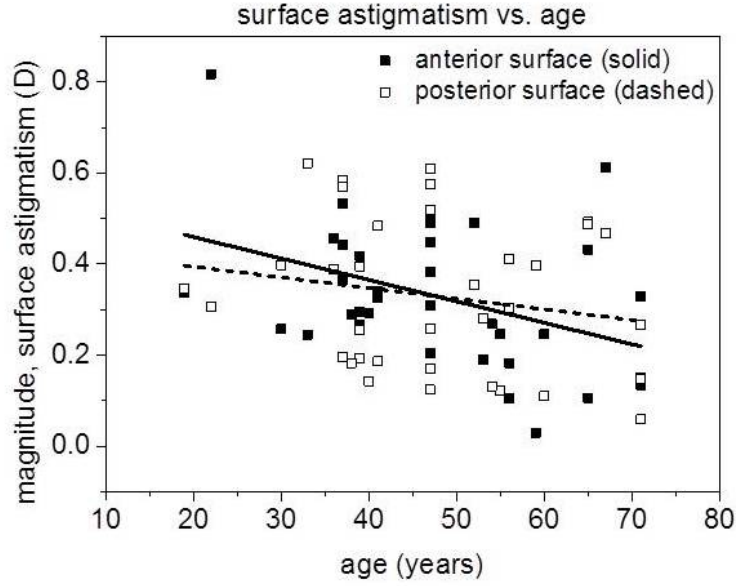


Figure 6.1: Change in the lens surfaces astigmatism with age. (anterior surface: slope= -0.005 D/year; $r = 0.397$, $p=0.018$; posterior surface: slope= -0.002 D/year, $r=0.189$, $p=0.276$)

astigmatism in the same lenses assuming a homogeneous equivalent refractive index ($r=0.380$, $p=0.024$). The mean value of the magnitude of astigmatism was -1.46 D for the GRIN lens and -1.02 D for the homogeneous refractive index lens. The ratio of magnitude of astigmatism of the GRIN lens to the homogeneous refractive index lens with age was not found to be statistically significant ($r=0.062$ $p=0.722$ slope= 0.004 D/year, not shown). In general, the absolute decrease of astigmatism with age seems to be driven by a decrease in the astigmatism surfaces, with a minor contribution of GRIN. In fact, the high correlation between the anterior lens surface astigmatism and the total astigmatism, even more so for the GRIN lens than for the lens with an equivalent refractive index (GRIN lens: $p=3.9E-6$ $r=0.693$; equivalent refractive index lens: $p=4.1E-4$, $r=0.565$, not shown) indicates, on average, a predominant role of the anterior lens surface. However, additional compensatory effects may be driven by the relative astigmatic angle between the two lens surfaces or the lens surfaces and the GRIN, or both, with potential age-dependent contributions.

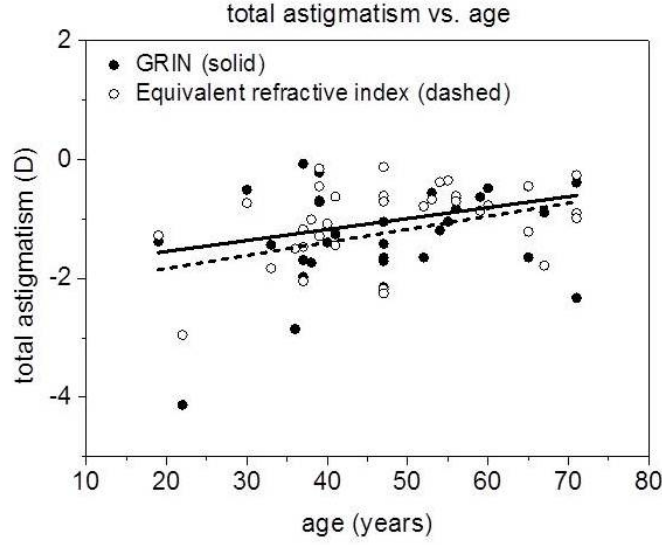


Figure 6.2: Change in the lens total astigmatism with age; lens with GRIN: slope: 0.022 D/year $r=0.359$, $p=0.034$; Lens with equivalent refractive index: slope: 0.018D/year, $r=0.380$, $p=0.024$).

6.3.2 Change of relative astigmatic angles with age

The lack of information on the actual orientation of the lens in vivo resulted in the definition of the astigmatic relative angle. We studied potential changes with age of the relative astigmatic angle between anterior and posterior surface axis of astigmatism, GRIN axis and anterior axis of astigmatism, GRIN axis and posterior axis of astigmatism, and between the lens total astigmatism axis with GRIN and the lens total astigmatism axis with a homogeneous refractive index. Figure 6.3 and 6.4 show polar plots (using the power vector analysis notation) for astigmatic magnitude and axes of the lens surface, the GRIN (the axis of GRIN astigmatism was associated with the axis of minimum value of the power exponent p_2 which would give the most increase in lens power), and the lens with both, GRIN and equivalent refractive index. For illustration purposes we assumed that all lenses were aligned such that the anterior lens axis was aligned with the steepest meridian along the vertical axis. The magnitude of astigmatism is indicated on the left scale in absolute values, and the angle, indicated in the polar plot, represent the aligned angle or the relative angle. Each dot represents one lens. Figure 6.3 a shows magnitude and aligned (steepest) meridian of the anterior surface astigmatism. Figure 6.3 b shows the shift in angle of the posterior surface astigmatic axis, relative to the aligned anterior axis, and the values of astigmatic magnitude of the posterior lens surface astigmatism. The average relative angle between anterior and posterior surface was 48.3

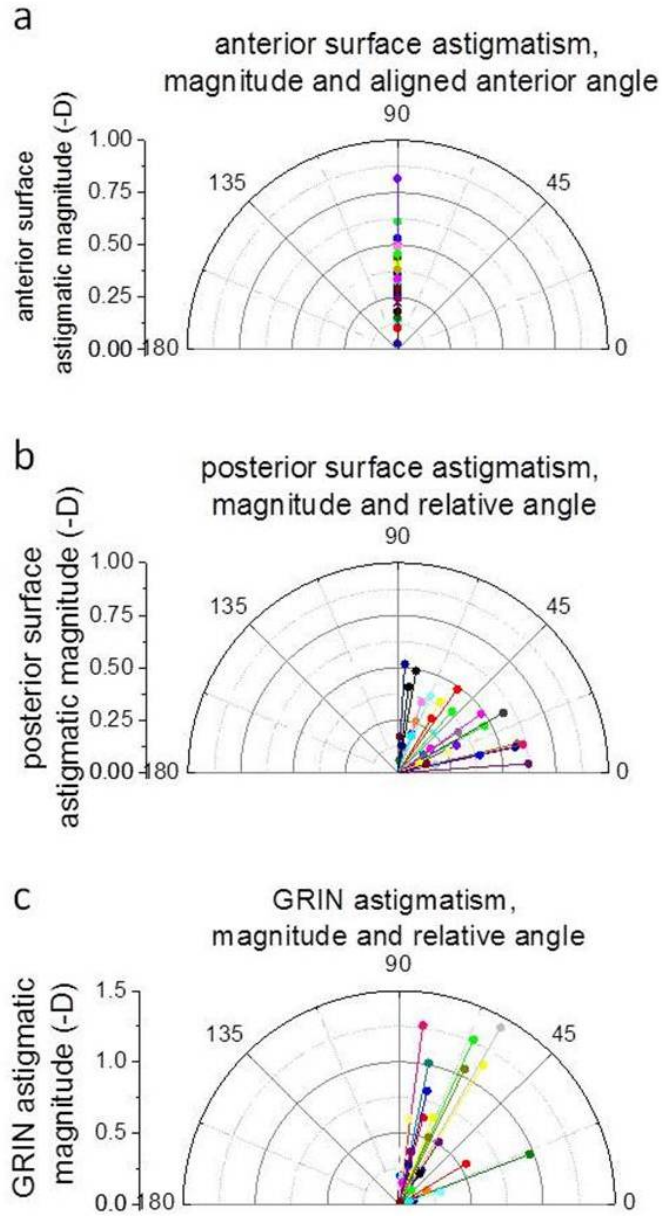


Figure 6.3: Power vector plot of a) anterior surface astigmatism (the astigmatic axis is vertically aligned) b) posterior surface astigmatism (with the relative angle between anterior and posterior astigmatic axis), and c) GRIN astigmatism (with the relative angle between anterior axis and the astigmatic axis of the GRIN).

deg, and for 45.7% of the lenses the relative angle was <45 deg. For the GRIN astigmatism (figure 4 c, derived from p2), the average relative angle to the anterior surface was 66.3 deg, and for 14.3% of the lenses the relative angle was <45 deg. Figure 6.4 shows

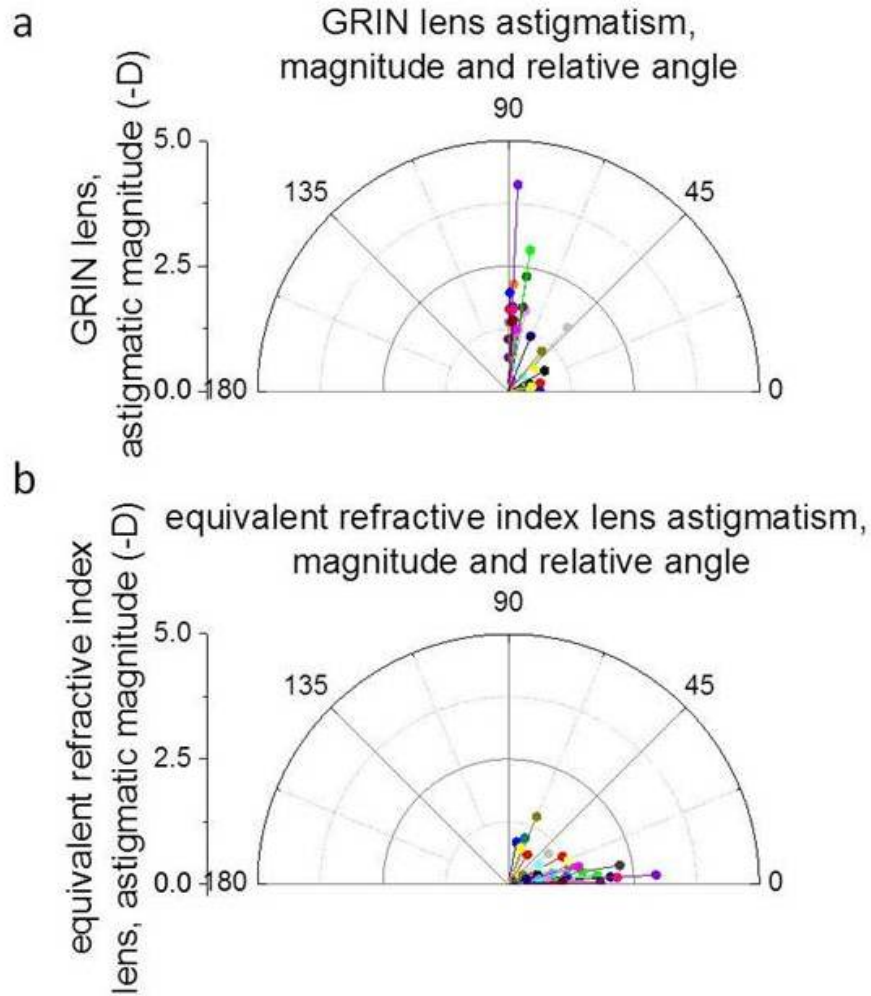


Figure 6.4: Figure 5: Power vector plot of a) the magnitude of astigmatism of the GRIN lens with the relative angle between GRIN lens astigmatic axis and aligned anterior astigmatic axis), and b) the magnitude of the lens with equivalent refractive index, with the relative angle between GRIN astigmatic axis and astigmatic axis of the equivalent refractive index lens.

the power vector graphs for total astigmatism for both, GRIN lenses (6.4a) and lenses with the corresponding homogeneous refractive index (6.4b). Figure 5a indicates the relative angle between the GRIN lens astigmatic axis and the aligned anterior surface astigmatic axis. The relative angle was <45 deg in 28.6 % of the lenses (average value 64,1 deg, with 54.3% of the lenses >80 deg). The relative angle between GRIN lens and the lens with a homogeneous refractive index (5b) was on average 22.9 deg, with 82.9% of the lenses <45 deg. Figure 6.5 shows the age-dependency of the relative astig-

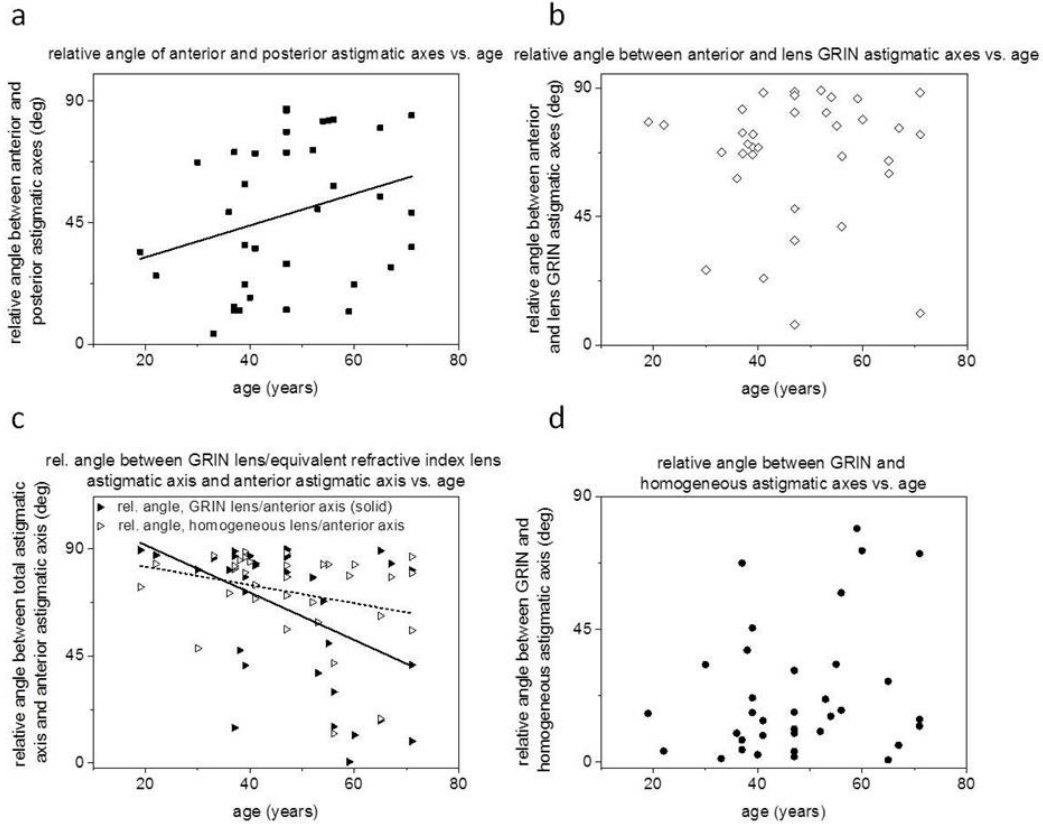


Figure 6.5: Relative astigmatic angle between (a) the anterior and posterior surface astigmatic axis as a function of age ($r=0.290$ $p=0.091$ slope= 0.586 deg/year); (b) the anterior astigmatic axis and the GRIN astigmatic axis as a function of age ($r=0.007$ $p=0.966$ slope= -0.012 deg/year); (c) anterior lens surface astigmatism and total lens astigmatism (with GRIN); $r=0.452$ $p=0.006$ slope= -0.997 deg/year and with homogeneous refractive index ($r=0.281$ $p=0.102$ slope= -0.384), and (d) GRIN astigmatic axis and equivalent refractive index lens astigmatic axis as a function of age ($r=0.393$ $p=0.169$ slope= -0.384 deg/year).

matic angle. Although there is a tendency for the relative astigmatic angle between the anterior and posterior lens to increase with age, this change does not reach statistical significance (Figure 4a; $r=0.290$, $p=0.091$, slope: 0.586). The meridional change of p_2 did not show any age dependency (figure 6.5b, $r=0.007$ $p=0.966$ slope= -0.012 , fit not shown). Interestingly, there is a difference between the total lens astigmatic axis (with GRIN) and that of the anterior lens axis (taken as a reference for each lens), and this difference decreases statistically significantly with age (Figure 6.5c; $r=0.452$ $p=0.006$ slope= -0.997). On average, the relative axis between the anterior lens and total lens astigmatism is close to 90 deg in the young lens, suggesting a compensatory role by the

GRIN and posterior surface, with age this relative angle decreases to values closer to 45 deg in the older lenses. Conversely, when the astigmatic axis with equivalent refractive index is similarly compared with the anterior lens, this dependency was not present (4c, dashed line, $r=0.281$ $p=0.102$ slope=-0.384), neither was any age dependency detected for the posterior lens surface (not shown). Figure 4d shows the relative angle between GRIN astigmatic axis and the astigmatic axis of the lens with the equivalent refractive index with age. As mentioned before, most of the lenses (82.9%) had a relative angle < 45 deg, throughout all ages. No age-dependency was found.

6.4 Discussion

In this study, we distinguish between lens astigmatism and surface astigmatism to find how the different lens characteristics (GRIN, individual surfaces) relate to each other. We found a significant decrease of the astigmatic magnitude of the anterior surface with age, from positive values towards zero (slope= -0.005 D/year; $r=0.397$, $p=0.018$), while the posterior surface and the GRIN of the lens did not change significantly with age. The total astigmatism of the lens decreased from negative values towards zero (GRIN: slope: 0.022 D/year $r=0.359$, $p=0.034$; Lens with equivalent refractive index: slope: 0.018D/year, $r=0.380$, $p=0.024$). These results suggest that the anterior lens surface has a larger impact to the total lens astigmatism, and that the influence of the GRIN to the lens astigmatism is limited. This behavior is also supported by the high correlation between the total lens astigmatism and the anterior lens astigmatism (GRIN lens: $p=3.9E-6$ $r=0.693$; equivalent refractive index lens: $p=4.1E-4$, $r=0.565$). Although our results regarding the orientation of the astigmatic axis lack of information on the actual orientation of the lens in vivo, we could study the relative angle between anterior and posterior surface axis of astigmatism, GRIN axis and anterior axis of astigmatism, GRIN axis and posterior axis of astigmatism, and between the lens total astigmatism axis with GRIN and the lens total astigmatism axis with a homogeneous refractive index. We did not find that the GRIN plays part in the astigmatism compensation with age, indeed, the relative angle between the GRIN lens astigmatic axis and the astigmatic axis of the equivalent refractive index lens was on average 22.9 deg, with 82.9% of the lenses <45 deg. In keeping with the predominant role of the anterior lens, we found a strong correlation between the anterior astigmatic axis and the GRIN astigmatic axis as a function of age ($r=0.007$ $p=0.966$ slope=-0.012 deg/year). The relative angle between anterior and posterior astigmatic axis had a tendency to increase, but showed no significant correlation ($r=0.290$ $p=0.091$ slope=0.586 deg/year).

Our results support the findings from Elawad[63] who found that direct astigmatism

came from the anterior lens surface. The magnitude and axis of the internal crystalline lens astigmatism has been a matter of controversy. Traditionally, Javals rule is assumed, implying a linear relationship between corneal and refractive astigmatism, with a constant offset of 0.5 D of against-the-rule astigmatism (arising from internal astigmatism). However, multiple modifications of this rule have been proposed, and significant deviations in the relationship between ocular and corneal astigmatism have been reported with age[108, 77, 78]. Corneal astigmatism is known to change its axis with age (from with-the-rule to against-the-rule[220][86]). Reports of the magnitude of corneal and refractive astigmatism differ across studies, with some works reporting no change of either, while others report significant changes of one or both. In those studies, internal astigmatism is computed indirectly from comparisons of corneal and ocular astigmatism, sometimes from different datasets, which may pose uncertainties.

The lack of knowledge of the real orientation to the cornea in our study, sets obvious limitations. Nevertheless, the study showed, that the crystalline lens' individual optics, like lens surfaces and the GRIN, can be studied individually to understand the purpose of every component in the crystalline lens as an optical instrument. In vivo studies are necessary to investigate these components with accommodation, but an ongoing challenge is the measurement of the GRIN in vivo from OCT data. Off-axis OCT seems to be a possible answer, as it was shown by de Castro et al. (A. de Castro,poster, ARVO2014).

7

Imaging crystalline lens microscopic structures of intact in vitro mammal lenses using confocal microscopy

This chapter is based on a paper in process "Imaging crystalline lens microscopic structures of intact in vitro mammal lenses using confocal microscopy". The coauthors of this study were Jorge Lamela, Sergio Ortiz, and Susana Marcos. The author of this thesis (i) was involved in the design of the microscope, (ii) built the microscope (in collaboration with Jorge Lamela and Sergio Ortiz), (iii) designed the experimental setup, (iv) did the sample preparation, (v) performed the experiments (in collaboration with Jorge Lamela and Sergio Ortiz), (vi) performed data processing, (v) analyzed the results, (vi) prepared the manuscript (in collaboration with Susana Marcos). As a result we could use a custom made confocal microscope to image lens epithelium, the lens suture, and lens fiber cells in different regions and depth of the crystalline lens in an in situ lens.

7.1 Introduction

In the last chapters we drew conclusions about the crystalline lens' optics based on the lens' shape and its reconstructed gradient index. This was done for in vitro lenses and

lenses in a stretching system to simulate the lens' natural accommodation in vivo. In the long run, these information have to be connected with structural aspects on the microscopic level. The ability to investigate the lens' shape, GRIN and microscopic structures with simulated accommodation (or in the living eye), would give greater insights into the interplay of optical properties and lens structures and would lead to a better and more comprehensive understanding of lens mechanisms. As a first step, we investigated the ability to image essential lens structures in an intact rabbit lens (still attached to the eye) by means of a custom built confocal microscope.

Confocal microscopy has been used as a way to observe the eye in an unstained, unfixed, living condition. It provides high-resolution, high-contrast images, and the capability to section the sample without contact. One of the main research fields of the confocal microscope in eye research is the in vivo investigation of the cornea[152, 271, 153, 31, 173, 184, 101, 185, 171, 144, 135, 134]. Therefore, literature on crystalline lens investigation, realized with confocal microscopy are scarce. Jeacocke et al.[102] and Xiao et al.[271] observed an isolated crystalline lens of a rabbit with a confocal microscopy. Xiao et al. could image lens fibers of a freshly excised lens about 200 μ m below the lens surface. Masters et al.[150] used confocal microscopy to image the fine structure of an in-situ rabbit crystalline lens from an ex-vivo rabbit eye through the full thickness of the cornea and aqueous humour. With this study, Masters could show the feasibility of confocal microscopy to achieve high-contrast images of transparent objects across 1.7 mm (crossing cornea and aqueous humour). Masters did a comparison study between confocal microscopy and SEM to further investigate the potential of confocal microscopy for observations of pathology in the anterior segment of the eye[154]. In vitro confocal microscopy showed high resolution images of the epithelium and superficial lens fibers. The compared SEM images fully supported the confocal images, and the study showed that confocal microscopy has the potential to become a useful tool for imaging the anterior segment of the eye.

Marcantonio et al.[148] used confocal microscopy to determine the level of cell coverage in human donor capsular bags with implants intraocular lenses, and concluded that they contain a large population of cells, even up to 13 years after surgery. Stachs et al. analyzed different cutting effects of femtosecond laser cuts with confocal microscopy[239]. He analyzed the cuts inside of crystalline lens tissue and found altered tissue scattering properties with intact lens fibers to concrete fiber separation, depending on femtosecond laser pulse energies. A study by Maddala et al.[143] also imaged lens fibers with a confocal microscope. The transparency of the lens depends on packing symmetry and membrane organization of the lens fibers. Concluding from his results, Maddala suggested a key role for the protein Periaxin in maturation, packing, and membrane organization of lens fiber cells. Reiss et al.[211] revealed a degenerative aging process

of porcine lens tissue, using Brillouin microscopy and confocal microscopy. Interesting approaches have been done using second harmonic generation (SHG) signals to investigate the cornea[93, 262]. Morishige[163] used two-photon confocal microscopy (with a variable-wavelength femtosecond lasers to produce SHG signals) to investigate the cornea of three mammalian species. He found that especially human corneas had a unique organizational pattern with sutural lamellae to provide important biomechanical support (not found in mouse or rabbit corneas). In this chapter, we demonstrated the potential of confocal light microscopy (CLM) in the anterior pole of the eye lenses by performing an in vitro study of in eye intact rabbit lenses. This method is suitable for quantifying the lens structures in the intact crystalline lens, holding promise for applications in vivo and for microscopic analysis of the lens under accommodative forces.

7.2 Methods

7.2.1 Lens sample preparation

Rabbit eyes were obtained from a local slaughterhouse and transported at a temperature around 4°C (figure 7.1A). The cornea was removed from the eye (figure 7.1B), and an iridectomy was done (figure 7.1C). The whole eye was then mounted on a ring for stabilization (figure 7.1D), and was positioned in a cuvette filled with water, whereby the water level was about 1 mm above the anterior surface of the eye (figure 7.1E). The cuvette was immediately placed under the microscopy objective (figure 7.1F). All eyes were imaged 2-24h post-mortem.

7.2.2 Confocal microscopy

Measurements were done with a custom made optical microscope, which can operate alternately or simultaneously as a confocal microscope or a multiphoton microscope. For this study, only the confocal part of the microscope was used. The confocal microscope is explained in detail in chapter 2. For these measurements, a 20 μm pinhole was used. The image resolution was set to 1000 by 1000 pixels. Each image took approximately 5 s. The objective was a Olympus water dipping objective (LUMPlan FL40x NA 0.8).

7.2.3 Imaging protocol

After the lens was positioned in the cuvette and under the microscope objective, the tip of the objective was dipped into the water and approximated to the crystalline lens,

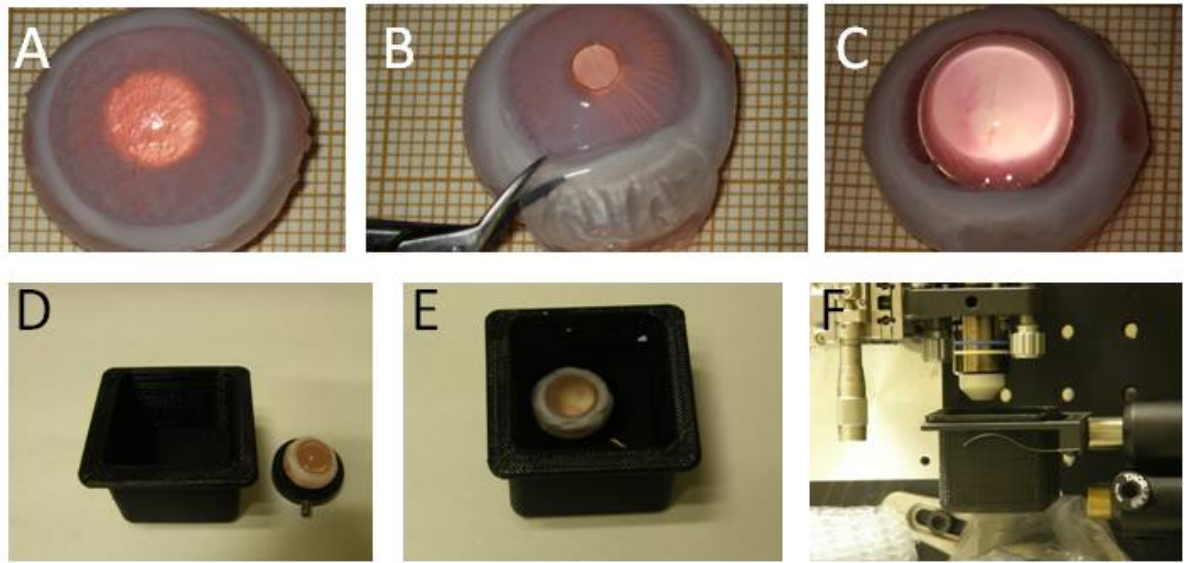


Figure 7.1: Lens sample preparation: A Fresh (3-48h post-mortem) rabbit eyes were obtained from the slaughterhouse B Cornea excision C Iridectomy D The eye is mounted on a ring for stabilization E The eye is positioned in a cuvette filled with water F The cuvette placed under the microscope objective

without touching it. The main areas of the lens which were images were the region around the apex of the lens (including the lens suture), the region in between the apex and the outer periphery, and the region at the very outer periphery (see figure 7.2).

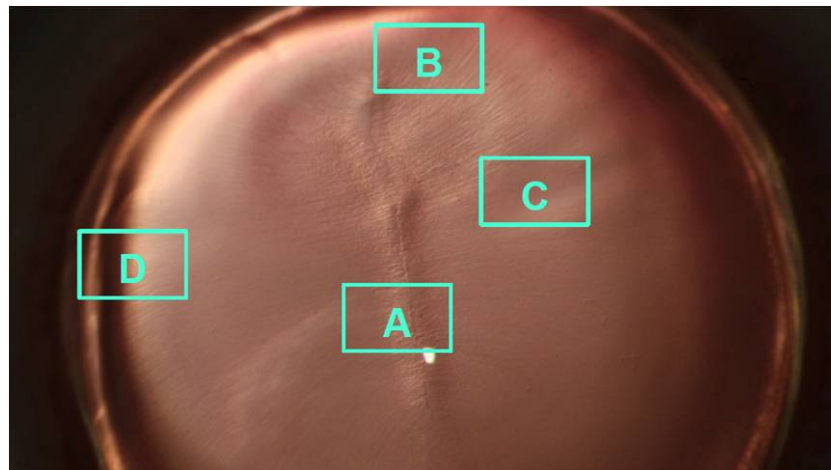


Figure 7.2: Regions images A region around the lens suture B, C Region in D outer periphery of the lens

In any of those regions, a z-scan was performed. The starting point was an area of 230

$\mu\text{m} \times 230 \mu\text{m}$ at the lens surface. The lens surface was defined as the lens capsule, which, under the confocal microscope, appears as a smooth surface. The z-scan consisted of imaging each area in step sizes in z-direction of 0.2 to 0.5 μm . The z-scan was performed as long as the image quality was considered reasonable. In general, the image depth is limited to the anterior lens cortex, because of the free working distance of the microscope objective and the low reflectivity of the ocular lens. After doing a z-scan in the central region of the lens, the cuvette was positioned to image another area of the lens. Because of the high curvature of the lens surface, the positioning usually included a manual inclination of the cuvette, so that the eventual region was positioned perpendicular to the laser beam.

7.2.4 Data analysis

The data was analyzed using image processing software ImageJ and Amira 5.3.3 (see chapter 2). ImageJ was used to remove all obvious scanning artifacts present in the images. Therefore, for every image, the Fourier transform was computed and the displayed power spectrum was edited to remove scanning artifacts. This was done by filtering or masking spots or lines on the frequency domain image (see figure 7.3). The inverse transform would then present an image which suppressed the selected frequencies and lead to an image with less artifacts. Background noise was filtered out using pre-programmed imageJ functions. All Images of one complete z-scan were then saved as a stack. A stack of z-scans was then loaded in Amira. The AlignSlices module was used to align the 2D slices to each other. That way, the resulting 3D image would not appear distorted. The stack was then visualized in 3D within the Amira editor, and color coded.

7.3 Results

The main results of this study are presented in figures 7.4 to 7.7. Figure 7.4(a) shows the lens fibers and the lens suture. The lens suture is the dark line, running perpendicularly to the well organized pattern of lens fibers. It is thought to be the link to two fiber endings, since the lens fibers are not long enough to stretch around the lens completely. Since the intracellular space between fibers cells are in the range of nm, we cannot really distinguish the single fiber cells. To approximate the width of the fiber cell, we used the dark region in between the fibers (which are indeed zones of low reflectivity, but not a visualization of intracellular space) to distinguish one fiber cell from another. With that definition, the fibers were counted over a predefined region of 50 μ , 100 μ , 150 μ and 200 μ and the length of the region was divided by fiber number (see Kuszak et al.[124]).

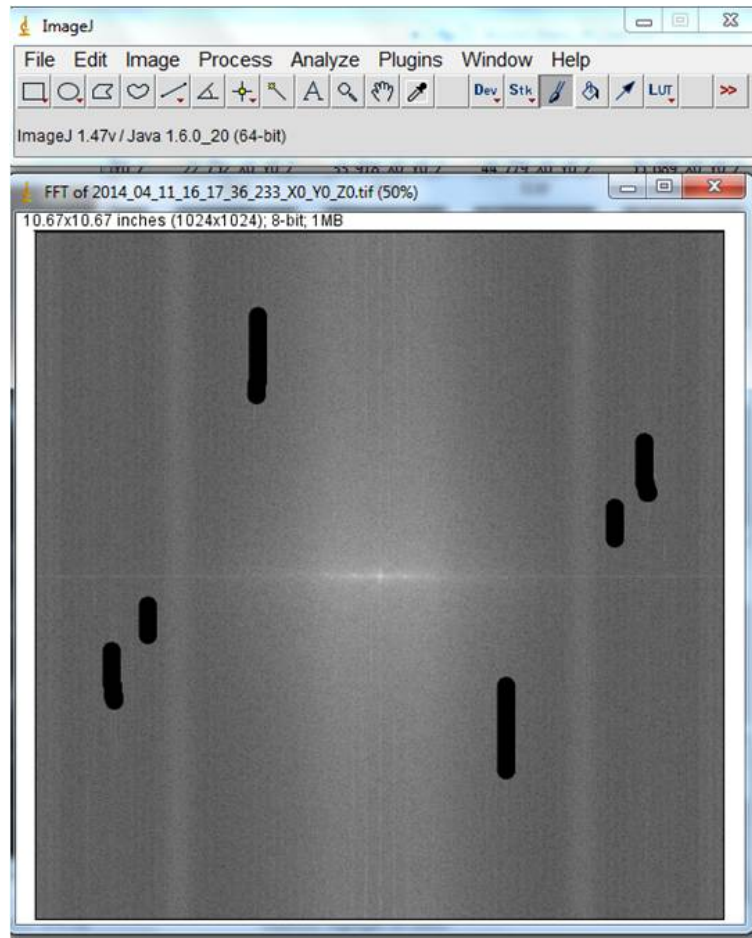


Figure 7.3: Edited power spectrum in imageJ (screenshot)

Hereinafter we give the fiber number per 50μ as a mean value.

In the region of the lens apix, the lens fibers were measured to a mean of 10 lens fibers per 50μ . Figure 7.4(b) shows a $350 \mu\text{m} \times 350 \mu\text{m} \times 25 \mu\text{m}$ 3D image of the same region of the line suture of the rabbit lens with the overlying epithelium layer visible (cut out region).

Figure 7.5 shows the lens fibers of region B (see figure 7.2), running diagonally from the lens suture. The lens fibers are on average wider than directly at the lens suture (9 lens fibers per $50\mu\text{m}$ in mean). The image is partly darkened because of the higher curvature of the lens in this region.

Images 7.6(a)-(c) were taken in between the periphery of the lens and the lens apex (region C in figure 7.2, $230 \mu\text{m} \times 230 \mu\text{m} \times 25 \mu\text{m}$).

Figure 7.6(a) (right, C1) shows the lens epithelium and (cut out) the beginning of

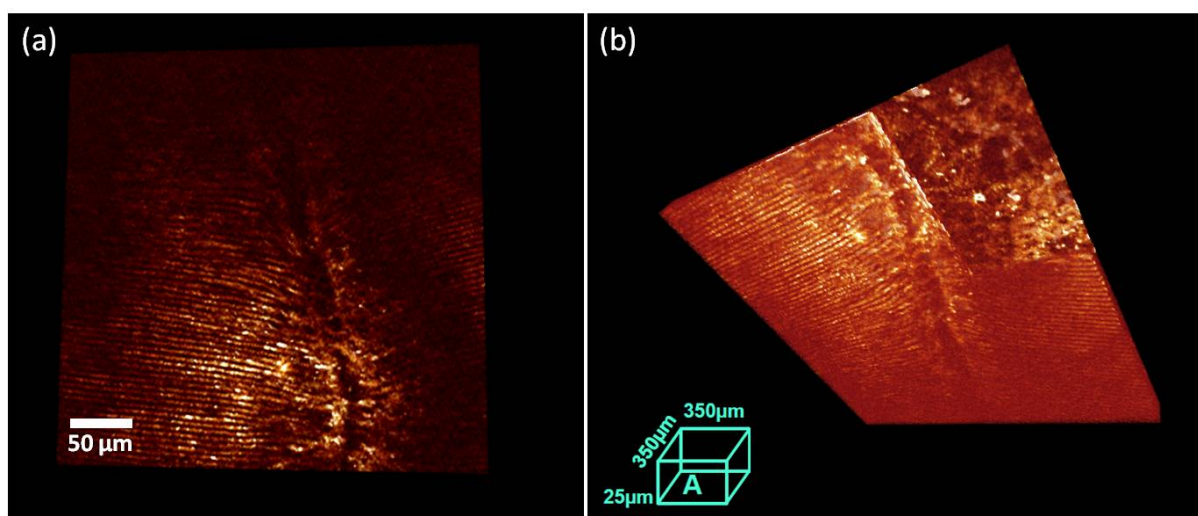


Figure 7.4: Image of lens suture taken in region A (see figure 7.2). (a) Lens suture and lens fibers (b) Lens suture and lens fibers with overlaying epithelium (cut out region)

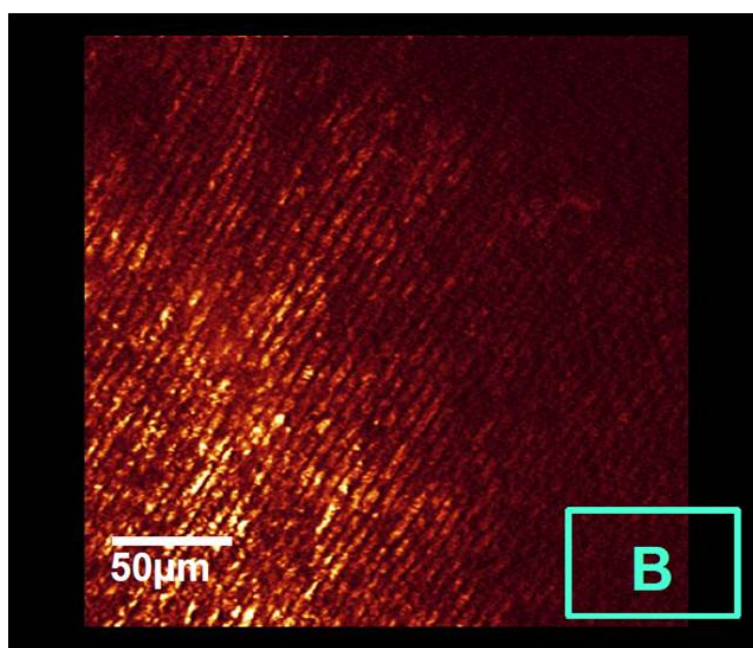


Figure 7.5: Lens fibers in region B (see figure 2). The lens fibers are on average wider than directly at the lens suture.

the subjacent layer of fiber cells. The lens epithelium can be easily identified by the roundish cell nuclei (*) spread over the layer, surrounded by cytoplasm. The size of the cell nuclei was measured between 10 and 25 μm . On the left side of figure 7.6(a), a more detailed image of the nuclei and surrounding cytoplasm is shown (110x110 μm region).

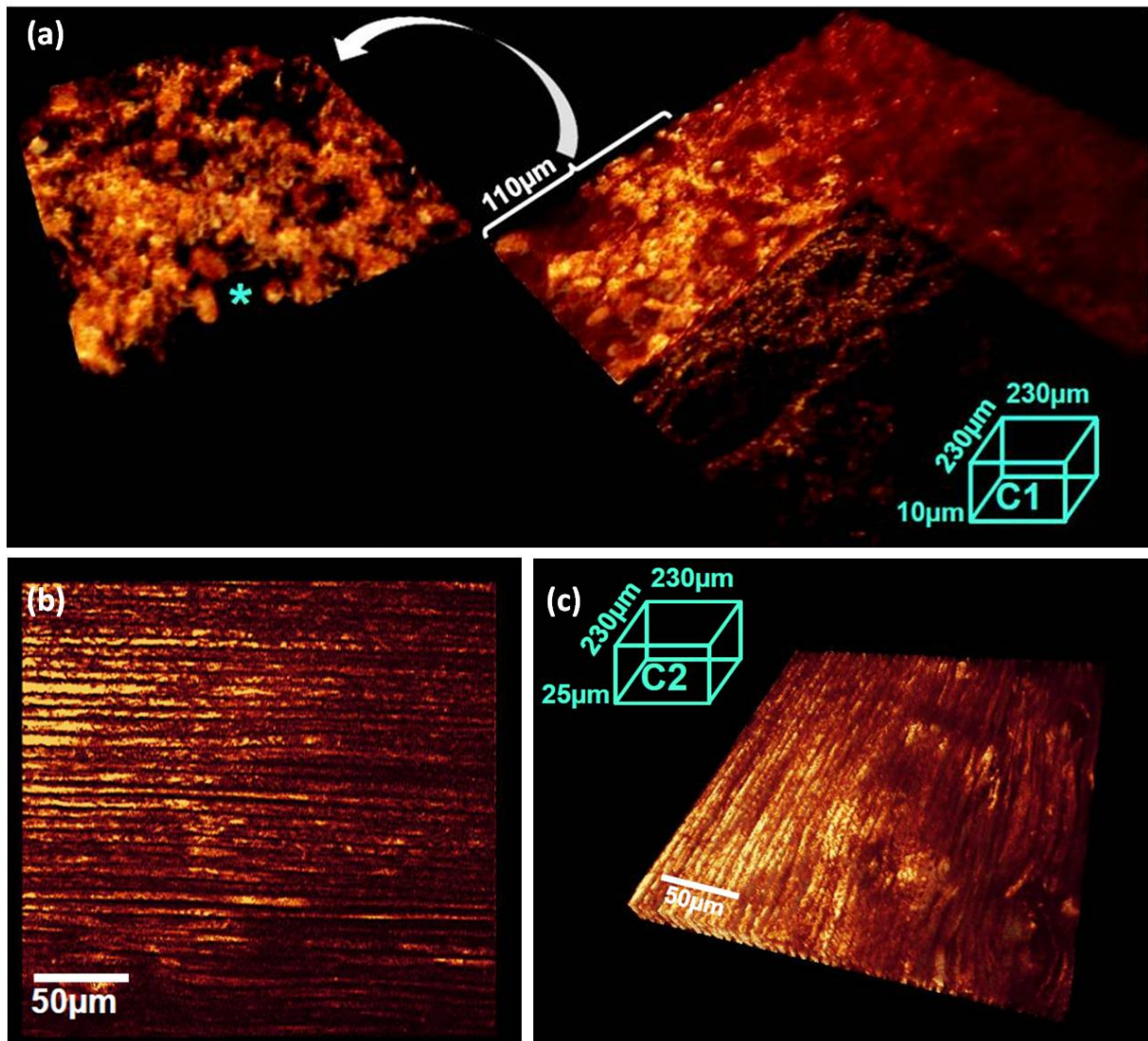


Figure 7.6: Z-scan through lens epithelium (a) with cell nuclei (*) and lens fibers (c). C1 and C2 are two parts of one z-scan, here shown separately for better visualization. (b) shows a 2D image of an image plane in C2.

As a continuation of the z-scan (approximately $10\ \mu\text{m}$ deeper towards the lens center), in figure 7.6 (c), a $25\ \mu\text{m}$ thick layer of fiber cells are shown. 8-9 fiber cells were counted within a $50\ \mu\text{m}$ region. In the lower right corner of figure 7.6(c), the lens fibers seem destroyed in their usually organized structure (also visible in the 2D view, lower left corner, figure 7.6(b)). The reason for these cavities is not clear. They might be similar to the findings of Masters et al.[151] who also found "large empty extracellular vacuolar elements of varying size in between the fibres" in regions close to the lens epithelium. It

might also come from heating effects and, consequently, bubble formation, as mentioned by Stachs et al.[239].

Figure 7.7 (a) to (c) show lens fibers at the outer periphery of the lens, hence, at a very curved region. Because of the high curvature of the lens in this region, epithelial cells and peripheral fibers are found in the same optical section (figure 7.7(a)). The lens fibers in this region are around $25\text{ }\mu\text{m}$ wide, hence, much wider than at the lens apex. Magnification of the lens fibers (figure 7.7(c)) show apparently tongue-and-groove so-called attachments (described by Dickson et al.[52] as "randomly distributed, elevated fingerprint-like patterns").

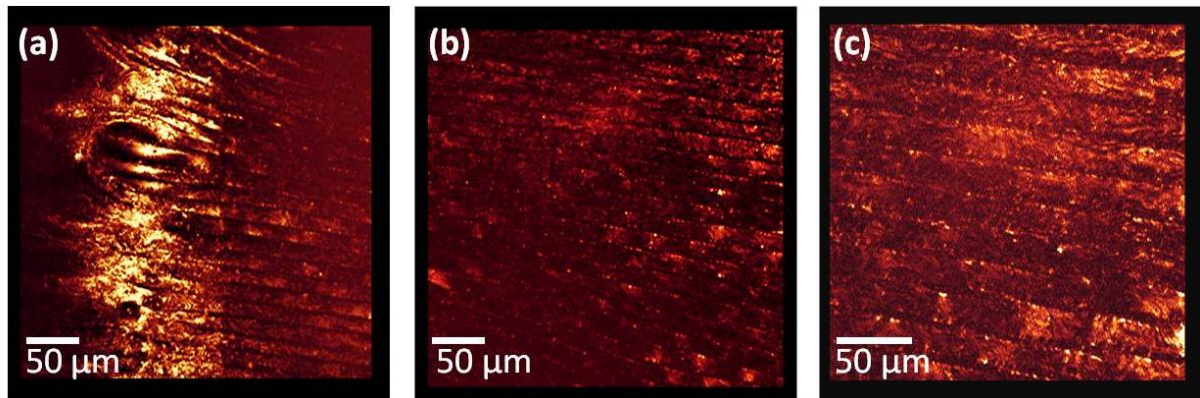


Figure 7.7: Lens fibers at the outer periphery of the lens (region D in figure 7.2). Because of the lens curvature, epithelial cells and peripheral fibers are found in the same optical section (figure 6 (a)). Between first (a) and second (b) image is a difference in z-direction of $5\text{ }\mu\text{m}$. (c) shows a magnification of the lens fibers.

7.4 Discussion

In the present study we proposed to use confocal microscopy for structural imaging of the intact crystalline lens in situ after corneal excision and iridectomy. We used a custom made confocal microscope to image lens epithelium, the lens suture, and lens fiber cells in different regions and depth of the crystalline lens. Since we were working with rabbit lenses, the lens suture was a vertical anterior line suture, and the lens fibers were the widest at the periphery of the lens, which is in accordance with the findings of Kuszak et al.[124], who found that the end of lens fibers taper to approximately $1/3$ - $1/2$ their width at the equator.

We were able to image the lens epithelium, with the lens nuclei prominently visible within the lens cytoplasm. The difference between the lens epithelium layer and the lens

fiber layer was well defined when imaging the lens in a region with great curvature, where epithelial cells and peripheral fibers cells could be found in the same optical section.

At the periphery the lens fibers showed small linking structures, like ball-and-socket interdigitations, tongue-and-groove attachments, and spikey posseses, which can be compared with images from scanning electron microscopy literature[52]. Usually the organization of lens fibers are investigated by means of scanning electron microscopy (for rabbit lenses for example Willekens et al.[267] and Kappelhof et al.[106]). SEM gives a great insight into lens structures up to nm range. Therefore, literature on lens fiber cells and their interlocking structure are usually based on SEM microscopy. Examples of different mammals are rats[266, 83, 65], rabbits [266, 84, 267], pigs[97], monkeys[266][52, 132, 268], and humans[132, 155, 253, 263, 44, 248]. The obvious downside of SEM is the destructive method it applies to obtain images. Confocal microscopy, on the other hand, can be used without even touching the lens, and this study was conducted to show the potential of confocal light microscopy in the anterior pole of the eye lenses, and to show that this method is suitable for quantifying the lens structures in the intact crystalline lens, holding promise for applications in vivo and for microscopic analysis of the lens under accommodative forces (as introduced by Glasser, ARVO2010). We found that confocal microscopy is well suited for possible investigations with simulated accommodation or even in vivo measurements to investigate the mechanics of lens fibers when accommodating.

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8

Activities during the Ph.D. Period 2010-14

8.1 Publications

ISI-WOK JCR data refers to the 2012 database

journal category: Ophthalmology (oph), Optics (opt)

(Authors are indicated by signature order)

1. [2014] Mengchan Sun, Judith Birkenfeld, Alberto de Castro, Sergio Ortiz, Susana Marcos, *OCT 3-D surface topography of isolated human crystalline lenses*
Article in Journal: BOE, vol. 5, pp. 3547-3561
(5year-IF: 3.179, 8/80 Opt. JCR 2012)
Publication date: 10/2014
2. [2014] Judith Birkenfeld, Alberto de Castro, Susana Marcos, *Contribution of shape and gradient refractive index to the spherical aberration of isolated human lenses.*
Article in Journal: IOVS, vol. 55, pp. 2599-2607
(5year-IF: 3.730, 6/59 Oph. JCR 2012)
Publication date: 04/2014
3. [2014] Pablo Prez-Merino, Judith Birkenfeld, Carlos Dorronsoro, Sergio Ortiz, Sonia Durn, Ignacio Jimnez-Alfaro, Susana Marcos, *Aberrometry in patients implanted with Accommodative Intraocular Lenses.*

Article in Journal: American Journal of Ophthalmology, vol. 157, pp. 1077-1089
(5year-IF: 4.292, 4/59 Oph. JCR 2012)

Publication date: 05/2014

4. [2014] Mengchan Sun, Alberto de Castro, Sergio Ortiz, Pablo Perez, Judith Birkenfeld, Susana Marcos, *Intraocular lens alignment from an en face OCT image Purkinje-like method.*

Article in Journal: Optical Engineering, vol. 53

(5year-IF: 0.783, 52/80 Opt. JCR 2012)

Publication date: 06/2014

5. [2014] Susana Marcos, Sergio Ortiz, Pablo Prez-Merino, Judith Birkenfeld, Sonia Durn, Ignacio Jimnez-Alfaro, *Three-dimensional evaluation of accommodating intraocular lens shift and alignment in vivo.*

Article in Journal: Ophthalmology, vol. 53, pp. 45-55

(5year-IF: 5.777, 2/59 Oph. JCR 2012)

Publication date: 01/2014

6. [2013] Alberto de Castro, Judith Birkenfeld, Bianca Maceo, Fabrice Manns, Esdras Arrieta, Jean-Marie Parel and Susana Marcos, *Influence of Shape and Gradient Refractive Index in the Accommodative Changes of Spherical Aberration in Nonhuman Primate Crystalline Lenses.*

Article in Journal: IOVS, vol. 54, pp. 6197-6207

(5year-IF: 3.730, 6/59 Oph. JCR 2012)

Publication date: 11/09/2013

7. [2013] Judith Birkenfeld, Alberto de Castro, Sergio Ortiz, Daniel Pascual, Susana Marcos, *Contribution of the gradient refractive index and shape to the crystalline lens spherical aberration and astigmatism.*

Article in Journal: Vision Research, vol. 86, pp. 27-34

(5year-IF: 2.345, 16/59 Oph. JCR 2012)

Publication date: 28 June 2013

8. [2013] Sergio Ortiz, Pablo Prez-Merino, Sonia Durn, Miriam Velasco-Ocana, Judith Birkenfeld, Alberto de Castro, Ignacio Jimnez-Alfaro, and Susana Marcos, *Full OCT anterior segment biometry: an application in cataract surgery*

Article in Journal: Biomedical Optics Express, vol. 4, pp. 387-396

(5year-IF: 3.179, 8/80 Opt. JCR 2012)

Publication date: 31 January 2013

9. [2011] S. Marcos, A. de Castro, E. Gamba, Judith Birkenfeld, S. Ortiz, P. Perez Marino, and C. Dorronsoro, *Ocular Imaging and crystalline lens Optical properties.* Published in Frontiers in Optics 2011/Laser Science XXVII, OSA Technical Digest

(Optical Society of America, 2011), paper FTuH4.
<http://www.opticsinfobase.org/abstract.cfm?URI=FiO-2011-FTuH4>

8.1.1 Works in Progress

1. Judith Birkenfeld, Alberto de Castro, Susana Marcos, *Astigmatism of the in vitro human lens: Surface and gradient refractive index contributions*
Expected publication date: 2014
2. Bianca Maceo, Alberto de Castro, Judith Birkenfeld, Esdras Arrieta, Jean-Marie A. Parel, Susana Marcos, Fabrice Manns, *Lens Spherical Aberrations in Cynomolgus Monkeys: Comparison of Laser Ray Tracing Measurements and Reconstructed GRIN Model Predictions*
Expected publication date: 2014
3. Alberto de Castro, Judith Birkenfeld, Bianca M. Maceo, Marco Ruggeri, Esdras A. Arrieta, Jean-Marie A. Parel, Fabrice Manns, Susana Marcos. *Crystalline lens gradient refractive index and posterior surface shape from multiple orientations OCT imaging: towards a reconstruction in vivo?*
Expected publication date: 2015
4. Judith Birkenfeld, Jorge Lamela, Sergio Ortiz, Susana Marcos *Imaging crystalline lens microscopic structures of intact in vitro mammal lenses using confocal microscopy.*
Expected publication date: 2015

8.2 Diffusion of the results

8.2.1 Summary

My work has been presented in 19 occasions, 6 given by me, 13 by my collaborators.

- 6 presentation in international conferences
 - 4 talks
 - 2 poster
- 13 presentation by my colleagues in international conferences
 - 8 talks
 - 5 posters

8.2.2 Talks and poster presented by Judith Birkenfeld

(Authors are indicated by signature order, *p* =poster, *t* =talk)

1. [2014 – *p*] ARVO 2014, Orlando, FL, USA, May 2014, poster on “Imaging crys-

talline lens microscopic structures of intact in vitro mammal lenses using confocal microscopy”, coauthors: Jorge Lamela, Sergio Ortiz, Susana Marcos

2. [2013 – *p*] ARVO 2013, Seattle, WA, USA, 8 May 2013, poster on “Contribution of shape and gradient index to the spherical aberration of donor human lenses”, coauthors: Alberto de Castro, and Susana Marcos
3. [2012 – *t*] Accommodation Club 2012, Miami, FL, USA, 11 May 2012, talk “Change in the Gradient Refractive Index of Cynomolgus Monkey Lenses with Simulated Accommodation”, coauthors: Alberto de Castro, Bianca Maceo, Esdras Arrieta, Fabrice Manns, Jean-Marie Parel, Susana Marcos
4. [2012 – *t*] ARVO 2012, Fort Lauderdale, FL, USA, 7 May 2012, talk “Contribution of Crystalline Lens Surface Shape and Gradient Index Distribution to Spherical Aberration”, coauthors: Alberto de Castro, Sergio Ortiz, Pablo Prez-Merino, Enrique Gamba, Daniel Pascual, Susana Marcos
5. [2012 – *t*] IONS 11, Paris, France, 22 February 2012, talk “Role of lens surface and gradient index for spherical aberration in mammal lenses”, coauthors: Alberto de Castro, Sergio Ortiz, Pablo Prez-Merino, Enrique Gamba, Daniel Pascual, and Susana Marcos
6. [2011 – *t*] ARVO 2011, Fort Lauderdale, FL, USA, 3 May 2011, talk “Three-dimensional reconstruction of the isolated human crystalline lens Gradient Index distribution”, coauthors: Alberto de Castro, Sergio Ortiz, Pablo Prez-Merino, Enrique Gamba, Susana Marcos

8.2.3 Presentation by Judith Birkenfeld’s collaborators

(Authors are indicated by signature order, *p* =poster, *t* =talk)

1. [2014 – *t*] ARVO 2014, Orlando, FL, USA, May 2014, talk “OCT-based reconstruction of the crystalline lens Gradient Refractive index: changes with age and accommodation”, authors: Susana Marcos et al.
2. [2014 – *p*] ARVO 2014, Orlando, FL, USA, May 2014, poster on “Crystalline lens gradient refractive index and posterior surface shape from multiple orientations OCT imaging: towards a reconstruction in vivo?”, authors: Alberto de Castro, Judith Birkenfeld, Bianca M. Maceo, Marco Ruggeri, Esdras A. Arrieta, Jean-Marie A. Parel, Fabrice Manns, Susana Marcos
3. [2014 – *p*] ARVO 2014, Orlando, FL, USA, May 2014, poster on “OCT 3-D surface topography of isolated human crystalline lenses”, authors: Mengchan Sun, Judith Birkenfeld, Alberto de Castro, Sergio Ortiz, Susana Marcos

4. [2014 – p] ARVO 2014, Orlando, FL, USA, May 2014, poster on “Lens Spherical Aberrations in Cynomolgus Monkeys: Comparison of Laser Ray Tracing Measurements and Reconstructed GRIN Model Predictions”, authors: Bianca M. Maceo, Alberto de Castro, Judith Birkenfeld, Esdras Arrieta, Jean-Marie A. Parel, Susana Marcos, Fabrice Manns
5. [2013 – t] International Myopia Conference, Asilomar, CA, USA, 19.-22. August 2013, talk “In vivo 3-D quantification of the crystalline lens based on Optical Coherence Tomography”, authors: Susana Marcos, Sergio Ortiz, Enrique Gamba, J.B., Mengchan Sun, Pablo Prez-Merino
6. [2013 – t] Bienal de Física 2013, Valencia, Spain, 15.-18. July 2013, talk “Microscopios multifotónica y de generación de segundo armónico: Aplicaciones al estudio del fenómeno de la acomodación del ojo y de nuevos tratamientos oculares”, authors: Jorge Lamela, J.B., Sergio Ortiz, Daniel Pascual, Susana Marcos
7. [2013 – t] Bienal de Física 2013, Valencia, Spain, 15.-18. July 2013, talk “Geometría y biometría 3D del segmento anterior del ojo con OCT”, authors: Sergio Ortiz, Pablo Prez-Merino, Judith Birkenfeld, Nicolás Alejandro, Sonia Duran, Susana Marcos
8. [2013 – t] ARVO 2013, Fort Lauderdale, FL, USA, 5 May 2013, talk “Three-dimensional biometry and alignment in eyes implanted with Accommodative IOLs as a function of accommodative demand”, authors: Susana Marcos, Sergio Ortiz, Pablo Perez-Merino, Miriam Velasco, Mengchan Sun, Judith Birkenfeld, Sonia Duran, Ignacio Jimenez-Alfaro
9. [2012 – t] X Reunión Nacional de Óptica, Zaragoza, Spain, 4-7 September 2012, talk “Análisis y cuantificación de imágenes 3D de OCT del segmento anterior del ojo: Técnicas y aplicaciones”, authors: Sergio Ortiz, Pablo Prez-Merino, Judith Birkenfeld, Nicolás Alejandro, Susana Marcos
10. [2012 – p] X Reunión Nacional de Óptica, Zaragoza, Spain, 4-7 September 2012, poster on “Líneas de Investigación del Laboratorio de Óptica Visual y Biofotónica (Instituto de Óptica, CSIC)”, authors: S. Marcos, S. Barbero, J. Birkenfeld, C. Dorronsoro, A. de Castro, D. Corts, E. Gamba, P. de Gracia, S. Kling, J. Lamela, D. Pascual, P. Prez-Merino, S. Ortiz, A. Radhakrishnan, L. Sawides, M. Sun, M. Velasco, M. Vias, N. Alejandro, S. Durn, I. Jimnez-Alfaro
11. [2012 – p] 6th Topical Meeting of Vision and Physiological Optics (EMVPO-2012), European Optical Society, Dublin, Ireland, 22 August 2012, talk “Quantifying performance of accommodative intraocular lenses with 3-D anterior segment OCT”, authors: Susana Marcos, Sergio Ortiz, Pablo Prez-Merino, Judith Birkenfeld, So-

nia Durn, Ignacio Jimnez-Alfaro

12. [2012 – t] 6th European Meeting on Visual And Physiological Optics (EMVPO) Dublin, Ireland, 20 August 2012, talk “Gradient refractive index reconstruction in accommodating non-human primate crystalline lenses”, authors: Alberto de Castro, Judith Birkenfeld, Bianca Maceo, Esdras Arrieta, Fabrice Manns, Jean-Marie Parel, Susana Marcos
13. [2011 – t] Frontiers in Optics, San Jose, CA, USA, 16-20 October 2011, talk “Ocular Imaging and crystalline lens Optical properties”, authors: Susana Marcos, Alberto de Castro, Enrique Gamba, Judith Birkenfeld, Sergio Ortiz, Pablo Perez Marino, and Carlos Dorronsoro

8.3 Visits and Stays in Research Institutions

1. [10/2013] Bascom Palmer Eye Institute, Miami, USA 2 weeks, from October 8 to October 25 of 2013. Topics: Gradient refractive index in monkey eyes / Off-axis OCT
2. [05/2012 – 07/2012] Bascom Palmer Eye Institute, Miami, USA 9.5 weeks, from May 5 to July 12 of 2012. Topics: Gradient refractive index in monkey eyes
3. [09/2011 – 12/2011] Bascom Palmer Eye Institute, Miami, USA 12 weeks, from September 16 to December 13 of 2011. Topics: Gradient refractive index in monkey eyes

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